

## Supplement to OPTMAN code, manual Version 10 (2008)

Efrem Sh. SOUKHOVITSKI\*, Satoshi CHIBA, Roberto CAPOTE†, Jose M. QUESADA‡,  
Satoshi KUNIEDA§ and Gennadij B. MOROGOVSKIJ\*

Advanced Science Research Center  
Japan Atomic Energy Agency  
Tokai, Naka, Ibaraki 319-1195, Japan

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A coupled-channel code to calculate nuclear cross sections, OPTMAN, has been improved after OPTMAN version 8 was published as JAERI-Data/code 2005-002. The most important addition is the inclusion of dispersive relationships between imaginary and real parts of the optical potential within a Lane consistent formalism. The obtained CCOM potentials allow for the description of nucleon induced reactions up to 200 MeV, including (p,n) reactions with excitation of isobaric analog states. Relativistic corrections consistent with those used in the ECIS06 code are also added.

Keywords: NUCLEAR REACTIONS  $\sigma_{tot}$ , direct (n,n), (n,n'), (p,p), (p,p'), (p,n'),  $\sigma_{p,nonel}$ , OPTMAN, Soft-rotator and Rigid-rotator Model, Collective level structure, Coupled-channels Method, Lane consistency, Coulomb correction, Isobaric analog states, ISTC B-1319, Theory, Numerical Algorithms, Inputs-outputs.

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Japan Atomic Energy Agency

\*Joint Institute of Energy and Nuclear Research, Sosny (Belarus)

†NAPC-Nuclear Data Section, International Atomic Energy Agency (Austria)

‡Departamento de Física, Atomica, Molecular y Nuclear, Universidad de Sevilla (Spain)

§ Nuclear Data Center, Japan Atomic Energy Agency

## OPTMAN コードマニュアル, Version 10 (2008)

Efrem Sh. SOUKHOVITSKI\*, 千葉 敏、Roberto CAPOTE†、Jose M. QUESADA‡、  
国枝 賢、Gennadij B. MOROGOVSKIJ\*

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チャンネル結合理論に基づいて原子核の断面積を計算するコード OPTMAN が改良されたのでマニュアルを更新する。重要な変更点はポテンシャルの虚数部と実数部を関係付ける分散関係をレーンモデルに基づいて導入した点である。これによって核子エネルギーで 200 MeV 程度までの断面積を計算することが可能である。また (p,n) アイソバリックアナログ状態への断面積の計算アルゴリズムと、ECIS06 と同様の相対論的補正を導入した。

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原子力科学研究所 (駐在): 〒 319-1195 茨城県那珂郡東海村白方白根 2-4

\* エネルギー原子核合同研究所 (ミンスク)

† 国際原子力機関

‡ セベリア大学

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# 1 Introduction

For more than twenty years, an original coupled-channels optical model code OPTMAN[1] has been developed at the Joint Institute of Energy and Nuclear Research to investigate nucleon-nucleus interaction mechanisms and as a basic tool for nuclear data evaluation for reactor design and other applications. Results of such activities for e.g.,  $^{235}\text{U}$ ,  $^{239}\text{Pu}$ ,  $^{236}\text{U}$ ,  $^{233}\text{U}$ ,  $^{238}\text{Pu}$  etc., were included in evaluated Nuclear Data Library BROND [2] of former Soviet Union. In addition to standard rigid rotator and harmonic vibrator coupling schemes encoded in widely-used JUPITER [3] and ECIS [4] codes, level-coupling schemes based on a non-axial soft-rotator model are included for the even-even nuclei in OPTMAN. This allows accounting of stretching of soft nuclei by rotations, which results in change of equilibrium deformations for excited collective states compared with that of the ground state. This is a critical point for reliable predictions [5, 6, 7] of the cross-sections based on the coupled-channels method.

OPTMAN had been developed over many years, and used for evaluation of reactor oriented nuclear data. So, originally it was written considering only neutrons as the projectile with possible upper incident energy of about 20MeV[5, 6, 7]. In 1995-1998, this code was successfully used as a theoretical base for nuclear data evaluation for minor actinides carried out in the framework of ISTC Project CIS-03-95, financial party of which was Japan. In 1997, the OPTMAN code was installed at Nuclear Data Center of Japan Atomic Energy Research Institute and an active collaboration started. After that time, many new options were added to the code following demands from a broad range of applications: power reactors, shielding design, radiotherapy, transmutations of nuclear wastes, and stellar nucleosynthesis.

Calculations with OPTMAN are now possible both for neutrons and protons as the projectile, and the upper incident nucleon energy is extended to at least 200 MeV [8]. The current version of soft-rotator model of OPTMAN takes into account the non-axial quadrupole, octupole and hexadecapole deformations, and  $\beta_2$ ,  $\beta_3$  and  $\gamma$ -vibrations with account of nuclear volume conservation. The OPTMAN is now able to analyze the collective level structure, E2, E3, E4  $\gamma$ -transition probabilities and reaction data in a self-consistent manner, which makes results of such analyses more reliable. We have found that this model was flexible enough so that OPTMAN can be applied not only to heavy rotational nuclei [9, 10], but can be also applied very successfully even to light nuclei, for example,  $^{12}\text{C}$ [11, 12] and  $^{28}\text{Si}$ [13], and medium vibrational nuclei such as  $^{52}\text{Cr}$  [14],  $^{56}\text{Fe}$  [15, 16] and  $^{58}\text{Ni}$  [17]. Later, energy dependence of the optical potential has been continuously improved guided by physical principles. Such features as the high-energy saturation behavior consistent with Dirac phenomenology, Elton and Madland relativistic corrections, and properties stemming from the nuclear matter theory had been taken into consideration. Therefore, OPTMAN already had capabilities applicable for analyses of nucleon interaction with light, medium and heavy nuclei for a wide energy

range, which will be crucially important to fulfill many nuclear data demands.

However, as the code, especially the mathematical algorithms were not described in detail before, and it were still a “black box” for most of the users, it were decided to undertake a complete modernization of the code. During 2001-2004 years currently available theoretical approaches and some new, more accurate advanced mathematical solutions and algorithms had been included in the code. They have resulted in a user-friendly code for complex coupled-channels optical model calculations. A complete OPTMAN code’s manual [1], giving details of the soft-rotator model, physical ideas and computation algorithms developed and incorporated into modernized code, description of input and output files, allowing easy running, understanding and using of OPTMAN outputs had been published in 2005. That manual covered the activity performed under the Project Agreement B-521 with the International Science and Technology Center (Moscow). The Financing Party for the Project was Japan.

After 2004 the main efforts in the development of the OPTMAN code were devoted to the improvement of the optical model formalism. We managed to incorporate dispersion relationships between imaginary and real parts of the optical potential [18], that allow to decrease significantly the number of optical parameters used, making them more physically grounded, that leads to the much better description of experimental data and accurate predictions of small differences in the total cross section along the given isotopic chain [19, 20].

We also extended the OPTMAN code allowing for calculations of the (p,n) cross sections with excitation of isobaric analog states and therefore, making possible to derive Lane consistent potentials including a new Coulomb correction [21]. Finally relativistic generalization used in the OPTMAN code had been made completely consistent with the one used in the latest versions of the ECIS [22] code.

All this modernization of OPTMAN code had been carried within activity performed under the Project Agreement B-1319 with the International Science and Technology Center (Moscow). The Financing Parties for the Project were Japan and European Union, job was requested by JAEA and IAEA.

## **2 New Options of OPTMAN Code**

New developments of the OPTMAN code follow. From now on we will refer to the previous code’s manual [1] as a *MANUAL*.

## 2.1 Dispersive Optical Potential Incorporated in OPTMAN Code

In this subsection we describe optical potential form used in OPTMAN when dispersive relations between real and imaginary parts of the potential are taken into account. For the non-dispersive potential formulae Eqs. (89)-(93) of the *MANUAL* can be used.

We consider that dispersive potential depending of energy can be written as

$$\begin{aligned}
V(r, R(\theta', \varphi'), E) = & - \left\{ V_{HF}(E) + \Delta V_R^{Coul}(E) \right\} \times f_{WS}(r, R_R(\theta', \varphi')) + \\
& - \left[ \Delta V_V(E) + \Delta V_V^{Coul}(E) + iW_V(E) \right] f_{WS}(r, R_V(\theta', \varphi')) \\
& - \left[ \Delta V_D(E) + \Delta V_D^{Coul}(E) + iW_D(E) \right] g_{WS}(r, R_D(\theta', \varphi')) \\
& + \left( \frac{\hbar}{m_\pi c} \right)^2 [V_{so}(E) + \Delta V_{so}(E) + iW_{so}(E)] \times \frac{1}{r} \frac{d}{dr} f_{WS}(r, R_{so}) \left( \hat{\sigma} \cdot \hat{L} \right) \quad (1) \\
& + V_{Coul}(r, R_c(\theta', \varphi'))
\end{aligned}$$

where the first term is the real smooth, so called Hartree-Fock (HF), volume potential. Successive complex-valued terms are the volume, surface and spin-orbit potentials, all containing the corresponding dispersive contributions  $\Delta V_V(E)$ ,  $\Delta V_D(E)$  and  $\Delta V_{so}(E)$ . It is known that the energy dependence of the depth  $V_{HF}(E)$  is due to the replacement of a microscopic nonlocal HF potential by a local equivalent. For a Gaussian non-locality,  $V_{HF}(E)$  is a linear function of  $E$  for large negative  $E$  and is an exponential for large positive  $E$ . Following Mahaux and Sartor [23], the energy dependence of the smooth HF part of the nuclear mean field is taken as that found by Lipperheide [24], accounting the isospin dependence [25]:  $V_{HF}(E) = V_R^{DISP} \left[ 1 + (-1)^{Z'+1} \frac{C_{viso}}{V_R^{DISP}} \frac{N-Z}{A} \right] \exp(-\lambda_R(E - E_F))$ . Apart from the standard prescription for the Coulomb correction already discussed as Eq. (94) in *MANUAL*, a new procedure has been now incorporated, which will be discussed in a later section. Similar Coulomb correction terms  $\Delta V_V^{Coul}(E)$  and  $\Delta V_D^{Coul}(E)$  are also calculated for volume  $\Delta V_V(E)$  and surface  $\Delta V_D(E)$  dispersive contributions to the real potential. The geometrical form factors are given as

$$\begin{aligned}
f_{WS}(r, R_i(\theta', \varphi')) & = \left\{ 1 + \exp \left[ \frac{r - R_i(\theta', \varphi')}{a_i} \right] \right\}^{-1}, \quad i = R, V, so \quad (2) \\
g_{WS}(r, R_D(\theta', \varphi')) & = -4a_D \frac{d}{dr} f(r, R_D(\theta', \varphi'))
\end{aligned}$$

where  $R_i(\theta', \varphi')$  denotes the deformed radii with the deformations considered, while spin-orbit potential is considered to be not deformed. The Coulomb potential  $V_{Coul}(r, R_c(\theta', \varphi'))$  was calculated using a multipole expansion of a charged ellipsoid with uniform charge density within the Coulomb radius  $R_c$  and zero outside as suggested by Bassel *et al.* [26]. The spherical term of the Coulomb potential was calculated by taking account of the diffuseness of the charge density distribution of the form  $f_c = [1 + \exp(r - R_c^0)/a_c]^{-1}$  [27].

In our formulation of the OMP in Eq.(1) the geometrical parameters of the Hartree-Fock potential  $R_R$  and  $a_R$  are in general different from the geometrical parameters  $R_V$ ,  $a_V$ ,  $R_D$ ,  $a_D$  of the volume and surface absorptive potentials; but the real and imaginary spin-orbit terms share the same  $R_{so}$  and  $a_{so}$  parameters. Therefore the volume dispersive contribution has different geometry (determined by  $R_V$  and  $a_V$ ) from the real smooth volume potential (determined by  $R_R$  and  $a_R$ ). As a result we have two separate volume contributions to the potential (as can be seen in the first and second line of Eq.(1)), effectively giving us more flexibility for the fitting of the experimental data. The present optical potential includes relativistic corrections as discussed by Elton [28] and explained in our recent paper [18].

It is useful to represent the variation of surface  $W_D(E)$  and volume absorption potential  $W_V(E)$  depth with energy in functional forms suitable for the dispersive optical model analysis. A commonly used energy dependence for the imaginary-surface term has been suggested by Delaroche *et al.* [29],

$$W_D(E) = A_D \frac{(E - E_F)^2}{(E - E_F)^2 + WID_D^2} \exp(-\lambda_D(E - E_F)). \quad (3)$$

An alternative energy dependence for the imaginary-surface term has been used by Charity *et al.* [30],

$$W_D(E) = A_D \left[ \frac{(E - E_F)^2}{(E - E_F)^2 + WID_D^2} - \frac{(E - E_F - E_{SHIFT})^2}{(E - E_F - E_{SHIFT})^2 + WID_D^2} \Theta(E - E_F - E_{SHIFT}) \right], \quad (4)$$

where  $\Theta(X)$  is Heaviside's step function, being non-zero if  $X > 0$ . One can see that contrary to Delaroche *et al.* [29] and *MANUAL* formulae (90) for which surface potential term had been decreasing exponentially, now as an option, it can decrease as the difference of two Breit-Wigner terms with the same height, but shifted by energy which could allow a better fit in some cases.

The isospin dependence of the surface and volume potential terms (the Lane term [25]) was considered in imaginary surface  $W_D(E)$  and volume  $W_V(E)$  potentials as follow,

$$A_{D,V} = W_{D,V}^{DISP} \left[ 1 + (-1)^{Z'+1} \frac{C_{wiso,wviso}}{W_{D,V}^{DISP}} \frac{N - Z}{A} \right]. \quad (5)$$

An energy dependence for the imaginary volume term has been suggested in studies of nuclear matter theory by Brown and Rho [31]:

$$W_V(E) = A_V \frac{(E - E_F)^2}{(E - E_F)^2 + (WID_V^{DISP})^2}. \quad (6)$$

The assumption that the imaginary potential  $W(E)$  is symmetric about  $E' = E_F$  is plausible for small values of  $|E' - E_F|$ . However, as was pointed out by Mahaux and Sartor [23] this approximate symmetry no longer holds for large values of  $|E' - E_F|$ . In fact the influence of the nonlocality of the imaginary part of the microscopic mean field will produce an increase of the empirical imaginary potential  $W(E')$  at large positive  $E'$  and approaches zero at large negative  $E'$  [32, 33]. The



DOM analysis of neutron scattering on  $^{27}\text{Al}$  [34] and  $^{232}\text{Th}$  [18] showed the importance of the non-local contribution to describe total cross-section  $\sigma_T$  data for energies above 100 MeV using a non-symmetric version of the volume absorptive potential for large positive and large negative energies. Following Mahaux and Sartor [23], we assume that the absorption strengths are only modified outside some fixed energy interval around the Fermi energy  $[E_F - E_a, E_F + E_a]$ . They used  $E_a = 60$  MeV, however this value is fairly arbitrary [23] and we will use it as a fitting parameter. Let us assume that the non-local imaginary potential to be used in the dispersive integral is denoted by  $\widetilde{W}_v(E)$ , then we can write [35]

$$\widetilde{W}_V(E) = W_V(E) - W_V(E) \frac{(E_F - E - E_a)^2}{(E_F - E - E_a)^2 + E_a^2}, \text{ for } E < E_F - E_a, \quad (7)$$

and

$$\widetilde{W}_V(E) = W_V(E) + \alpha \left[ \sqrt{E} + \frac{(E_F + E_a)^{3/2}}{2E} - \frac{3}{2} \sqrt{(E_F + E_a)} \right], \text{ for } E > E_F + E_a. \quad (8)$$

These functional forms are chosen in such a way that the function and its first derivative are continuous at  $E' = |E_F - E_a|$ . At large negative energies the volume absorption decreases and goes asymptotically to zero. On the contrary, at large positive energies nucleons sense the ‘‘hard core’’ repulsive region of the nucleon-nucleon interaction and  $\widetilde{W}_V(E)$  diverges like  $\alpha\sqrt{E}$ . Using a model of a dilute Fermi gas hard-sphere the coefficient  $\alpha$  was estimated to be equal to  $1.65 \text{ MeV}^{1/2}$  by Mahaux and Sartor[33], assuming that the Fermi momentum  $k_F$  is equal to  $1.36 \text{ fm}^{-1}$  and that the radius of the repulsive hard core is equal to  $0.4 \text{ fm}$ . However, as pointed out by Charity *et al* [30], the value of this quantity is not constrained theoretically and they obtained  $\alpha$  values ranging from 0 to  $2.2 \text{ MeV}^{1/2}$ . With some other assumptions  $\alpha$  can be as small as 0.3, so we decided to keep this parameter adjustable.

### 2.1.1 Dispersive corrections

In a dispersive relation treatment, the real potential strength consists of the Hartree-Fock term  $V_{HF}(\mathbf{r}, E)$  which varies slowly with energy, and a correction term  $\Delta V(\mathbf{r}, E)$  calculated using a dispersion relation. Under favorable conditions of analyticity in the complex  $E$ -plane the real part  $\Delta V$  can be constructed from the knowledge of the imaginary part  $W$  on the real axis through the dispersion relation

$$\Delta V(\mathbf{r}, E) = \frac{\mathcal{P}}{\pi} \int_{-\infty}^{\infty} \frac{W(\mathbf{r}, E')}{E' - E} dE', \quad (9)$$

where we have now explicitly indicated the radial and energy dependence of these quantities and  $\mathcal{P}$  means that the principal value of the integral should be taken. To simplify the problem, the geometry of the imaginary terms of the OMP are usually assumed to be energy-independent. They are expressed in terms of a Woods-Saxon function  $f_{WS}(r, R_i(\theta', \varphi'))$  or its derivative  $g_{WS}(r, R_i(\theta', \varphi'))$ .

In such case the radial functions factorize out of the integrals and the energy-dependence is completely accounted for by two overall multiplicative strengths  $\Delta V(E)$  and  $W(E)$ . Both of these factors contain, we note, volume and surface contributions. The dispersive treatment employed in this work to calculate dispersive contributions  $\Delta V_V(E)$ ,  $\Delta V_D(E)$  and  $\Delta V_{so}(E)$  was described in detail in ref. [18], so we refer interested readers to these contributions. Details of numerical and analytical methods incorporated in the OPTMAN code to calculate dispersive contributions are thoroughly described in original references [36, 37, 38]

## 2.2 Changes in Relativistic Generalization of Non-Relativistic Schrödinger Equation Consistent with One used in ECIS Code

The OPTMAN code already included relativistic generalization and in the present version we decided to make it consistent with what is coded in the ECIS [4] code. Let us define that the total relativistic energy of the projectile  $E_p^t$  is

$$E_p^t = E_p + m_p c^2,$$

so we can determine the relativistic factor  $\gamma$  for projectile

$$\gamma = \frac{E_p + m_p c^2}{m_p c^2},$$

and its momentum as

$$p_p^2 c^2 = E_p^{t2} - m_p^2 c^4 = (\gamma^2 - 1) m_p^2 c^4.$$

Assuming an interacting target at rest with mass  $M$ , the total energy of the interacting system is

$$E_{tot} = E_p + m_p c^2 + M c^2 = \gamma m_p c^2 + M c^2.$$

Now we can determine the velocity of the  $CM$  (center-of-mass) in the laboratory frame

$$V_{CM} = \frac{p_p c^2}{E_{tot}},$$

so that using relativistic factor  $\gamma$

$$V_{CM}^2 = \left( \frac{p_p c^2}{E_{tot}} \right)^2 = \frac{(\gamma^2 - 1) m_p^2 c^6}{(\gamma m_p c^2 + M c^2)^2} = \frac{(\gamma^2 - 1) m_p^2 c^2}{(\gamma m_p + M)^2}.$$

Using Lorentz transformations, we can easily find momentum  $p$  of interacting particle and target (they are opposite and equal in  $CM$ ), assuming that the momentum of the target in Lab.system vanishes

$$p = \frac{V_{CM}/c^2 M c^2}{\sqrt{1 - V_{CM}^2/c^2}} = \frac{V_{CM} M}{\sqrt{1 - V_{CM}^2/c^2}},$$

and

$$\begin{aligned}
p^2 &= \frac{V_{CM}^2 M^2}{1 - V_{CM}^2/c^2} = \frac{(\gamma^2 - 1)m_p^2 c^2}{(\gamma m_p + M)^2} M^2 / \left(1 - \frac{(\gamma^2 - 1)m_p^2 c^2}{(\gamma m_p + M)^2 c^2}\right) \\
&= \frac{(\gamma^2 - 1)m_p^2 c^2 M^2}{(2\gamma m_p M + M^2 + m_p^2)}.
\end{aligned}$$

Relativistic interacting particle wave number is  $k^2 = \left(\frac{p}{\hbar}\right)^2$ , so energies of interacting particle and target in the *CM* system are

$$\begin{aligned}
E_{p,CM}^2 &= p^2 c^2 + m_p^2 c^4 = m_p^2 c^4 \frac{\gamma^2 M^2 + 2\gamma m_p M + m_p^2}{2\gamma m_p M + M^2 + m_p^2}, \\
E_{M,CM}^2 &= p^2 c^2 + M^2 c^4 = M^2 c^4 \frac{\gamma^2 m_p^2 + 2\gamma m_p M + M^2}{2\gamma m_p M + M^2 + m_p^2}.
\end{aligned}$$

Above we considered  $p = p_p = p_M$  in the *CM*; in the OPTMAN code we use these energies instead of masses in reduced mass formulae:

$$\frac{E_{p,CM} E_{M,CM}}{E_{p,CM} + E_{M,CM}} = m_p M c^2 \frac{\sqrt{[\gamma^2 M^2 + 2\gamma m_p M + m_p^2][\gamma^2 m_p^2 + 2\gamma m_p M + M^2]}/[2\gamma m_p M + M^2 + m_p^2]}}{\left\{m_p \sqrt{\gamma^2 M^2 + 2\gamma m_p M + m_p^2} + M \sqrt{\gamma^2 m_p^2 + 2\gamma m_p M + M^2}\right\}}. \quad (10)$$

So we found necessary formulas for incident particle wave number and reduced mass of the particle + zero velocity target in terms of relativistic factor  $\gamma$  and masses in relativistic case. The corresponding energy of the system  $E_{CM}$  in the *CM* system is

$$\begin{aligned}
E_{CM} &= E_{p,CM} + E_{M,CM} \\
&= \frac{m_p c^2 [\gamma^2 M^2 + 2\gamma m_p M + m_p^2]^{1/2} + M c^2 [\gamma^2 m_p^2 + 2\gamma m_p M + M^2]^{1/2}}{[2\gamma m_p M + M^2 + m_p^2]^{1/2}}.
\end{aligned}$$

The kinetic energy  $E$  is equal to the *CM* energy minus the rest mass of the components:

$$\begin{aligned}
E &= E_{CM} - m_p c^2 - M c^2 \\
&= \left\{ m_p c^2 ([\gamma^2 M^2 + 2\gamma m_p M + m_p^2]^{1/2} - [2\gamma m_p M + M^2 + m_p^2]^{1/2}) \right. \\
&\quad \left. + M c^2 ([\gamma^2 m_p^2 + 2\gamma m_p M + M^2]^{1/2} - [2\gamma m_p M + M^2 + m_p^2]^{1/2}) \right\} / [2\gamma m_p M + M^2 + m_p^2]^{1/2}.
\end{aligned}$$

Following Elton's suggestions [28], we can multiply optical potential strengths (except for the spin-orbit and Coulomb terms) by a factor  $K(E)$ , as a relativistic optical potential generalization arising from the reduction of the Dirac equation to the corresponding Schroedinger equation for non-relativistic target and relativistic projectile. Elton suggests  $K(E)$  to be  $E/(M_p c^2)$ . Therefore, the factor grows without limit as the projectile kinetic energy  $E$  grows. We defined this factor as suggested by Madland [39], namely  $K(E) = 2E/(E + M_p c^2)$ , which saturates at 2 as incident energy grows, and allows easier fitting of experimental data. This factor is optional and controlled by the input key MEREL (see below).

### 2.3 Lane consistency of the OMP and excitation of the Isobaric Analog States (IAS)

An isospin-dependent coupled channel optical model potential, can be used to predict quasi-elastic (p,n) scattering to the isobaric analogue states (IAS) of the target nucleus; being such exercise the best test of the quality of the isovector part of the optical potential. It has been pointed out by Lane [25] that the optical model potential can be written in a charge-independent form. The extent to which we can state that a derived optical model potential is *Lane-consistent* can be established from the basic Lane equations [25]:

$$\begin{aligned} V_{pp} &= V_0 + \frac{N - Z}{4A} V_1, \\ V_{nn} &= V_0 - \frac{N - Z}{4A} V_1, \\ V_{pn} &= \frac{\sqrt{N - Z}}{2A} V_1, \end{aligned} \quad (11)$$

where  $V_0$  and  $V_1$  are the isoscalar and isovector components of the potential with the Coulomb interaction switched off, respectively. In such way one is able to calculate the charge-exchange channels in a (p,n) reaction (to the *elastic IAS* and excited states of the rotational band of the residual nucleus).

An accurate calculation of the nucleon scattering from deformed nuclei must include the coupling to the low-lying collective states. A very successful computational method to account for the importance of the multistep processes is the coupled-channel (CC) method using Tamura's formalism [3], which permits an exact solution of the Lane CC equations. The Coulomb displacement energy,  $\Delta_C$ , between the ground state and its corresponding IAS is well approximated by the empirical relation [40],  $\Delta_C = 1.444Z/A^{1/3} - 1.13$  MeV, with  $Z$  being equal the average charge of the target and residual nuclei in the reaction. As an example for actinide targets, the value of  $\Delta_C$  is about 20 MeV. The coupling form-factors for the charge-exchange calculations are defined as

$$\begin{aligned} &\langle \nu; 0_{IAS}^+ | V(\tau, \vec{r}) | \pi; 0_{gs}^+ \rangle \\ &= \langle \nu | \mathcal{T} | \pi \rangle \langle 0_{IAS}^+ | V_1^{diag}(r) | 0_{gs}^+ \rangle \\ &= \frac{\sqrt{(N - Z)}}{2A} \langle 0_{IAS}^+ | V_1^{diag}(r) | 0_{gs}^+ \rangle, \end{aligned} \quad (12)$$

for the "quasi-elastic"  $0_{gs}^+ \rightarrow 0_{IAS}^+$  excitation of the IAS, as a particular case of

$$\begin{aligned} &\langle \nu; I^+(\text{residual}) | V(\tau, \vec{r}) | \pi; I^+(\text{target}) \rangle \\ &= \langle \nu | \mathcal{T} | \pi \rangle \langle I^+(\text{residual}) | V_1(\vec{r}) | I^+(\text{target}) \rangle \\ &= \frac{\sqrt{(N - Z)}}{2A} \langle I^+(\text{residual}) | V_1(\vec{r}) | I^+(\text{target}) \rangle, \end{aligned} \quad (13)$$

for the coupling between analog states of both rotational bands, and

$$\begin{aligned}
& \langle \nu; I^{+'}(\text{residual}) | V(\tau, \vec{r}) | \pi; I^+(\text{target}) \rangle \\
& = \langle \nu | \mathcal{T} | \pi \rangle \langle I^{+'}(\text{residual}) | V_1^{\text{coupl}}(\vec{r}) | I^+(\text{target}) \rangle \\
& = \frac{\sqrt{(N-Z)}}{2A} \langle I^{+'}(\text{residual}) | V_1^{\text{coupl}}(\vec{r}) | I^+(\text{target}) \rangle, \tag{14}
\end{aligned}$$

for the coupling between  $I' \neq I$  states. In these expressions  $\mathcal{T}$  is the isospin operator,  $\nu$  and  $\pi$  represent the entrance and exit isospin states of projectile and ejectile, respectively, and  $V_1^{\text{diag}}(r)$  and  $V_1^{\text{coupl}}(\vec{r})$  are the usual spherical and deformed components of the isovector potential, as defined in the Tamura's *canonical* work [3]. These expression have been implemented into the OPTMAN code allowing to directly calculate the quasi-elastic scattering cross-section or to consider existing IAS scattering data during the optical potential fitting. In this way the isovector component of the potential is much better constrained.

By the moment calculations of IAS excitations in OPTMAN are possible for coupling built on rigid rotator wave functions (rotational model) only, but it will be also possible for soft rotator option of OPTMAN soon.

## 2.4 Lane Consistent Coulomb Correction Option

The Coulomb correction used for incident charged particles are well established. In the OPTMAN code, the Coulomb correction is applied for incident protons to account for the change of the interacting proton energy due to the Coulomb repulsion by nucleus. In the previous OPTMAN versions such corrections were assumed either as an energy-independent constant to be added to the real potential or as the value proportional to the derivative of the real potential:

$$\Delta V_R^{Coul}(E) = -C_{Coul} \frac{ZZ'}{A^{1/3}} \frac{d}{dE}(V_R(E)). \tag{15}$$

One can see, that such correction is just the first term of Taylor expansion of the proton potential accounting for the Coulomb repulsion, assuming that the "effective" interacting energy of the proton is  $E - C_{Coul} \frac{ZZ'}{A^{1/3}}$ . It should be noted that constant  $e^2$  in our definition is included in constant  $C_{Coul}$ . In fact, the term  $C_{Coul} \frac{ZZ'}{A^{1/3}}$  is an estimation of the kinetic energy loss of the incident proton in the interaction region due to the Coulomb repulsion[41]. Indeed, the optical potential at this "effective" energy becomes

$$V(E - C_{Coul} \frac{ZZ'}{A^{1/3}}) = V(E) - C_{Coul} \frac{ZZ'}{A^{1/3}} \frac{d}{dE}(V_R(E)) + \dots \tag{16}$$

Above formula (16) is a generalization of the previously used Coulomb correction, which considers such corrections in all orders. It has been included in the OPTMAN code by using the "effective"

energy  $E - C_{Coul} \frac{ZZ'}{A^{1/3}}$  for incident protons instead of the physical incident energy  $E$ . The constant  $C_{Coul}$  is an adjustable constant meant to account for the "effective" radius of interaction of proton in nucleus. It is expected to be near one. It should be noted that full Coulomb correction as defined by [16] is a pre-condition to the exact Lane consistency of the derived OMP.

### 3 Changes in OPTMAN Code Inputs Due to New Options

We describe changes in the input data of OPTMAN code following *MANUAL* notations [1]. Units of the data are MeV, fm, barn and amu (Carbon units). Data have card image form and we'll show below only those Records changed due to new options, the other Records will not be described.

#### 3.1 Description of the Input Records, that had been Changed

- Record 2 - FORMAT(16I2)

MEJOB, MEPOT, MEHAM, MEPRI, MESOL, MESH, MESH, MEHAO, MEAPP, MEVOL, MEREL, MECUL, MERZZ, MERRR, MEDIS, MERIP

Two new switches MEDIS and MERIP describing the models options had been added.

- MEJOB = 1 - optical calculations, = 2 - adjustment of optical potential parameters;
- MEPOT = 1 -rotational model potential, =2 - potential expanded by derivatives (must be used for soft-rotor potentials);
- MEHAM = 1 - nuclear Hamiltonian of rotational model, = 2 - not in use yet, = 3 - Davydov-Chaban model, = 4 - Davydov-Philipov model, = 5 - nuclear Hamiltonian with account of  $\gamma$  softness [6, 7] , =6,7 - not in use yet;
- MEPRI = 0 - short output, additional output as MEPRI is growing ;
- MESOL = 1 - code will choose which method of coupled channels system solution for a certain  $J^\pi$  should be used to reduce calculation time, = 2 - exact solution, =3 - solution using iterations, with spherical solutions as zero approximation, >3 - solution using iterations, with exact coupled channels solution with the number of coupled states equal MESOL as zero approximation, MESOL must be less than or equal to 20;
- MESH = 1 - rigid hexadecapole deformations are not taken into account, = 2 - with account of axial rigid hexadecapole deformations, = 3 - rigid hexadecapole deformations depending on  $\gamma$  (Ref.[42]), = 4 - rigid hexadecapole deformations in the most general case [42] ;

- MESH0 = 0 - nuclear shape without octupole deformations, = 1 - nuclear shape with axial octupole deformations, = 2 - nuclear shape with non-axial octupole deformations;
- MEHA0 = 0 - nucleus is rigid to octupole deformations, = 1 - nucleus is soft to non symmetric octupole, = 2 - nucleus is soft to symmetric octupole deformations scaled by  $\beta_2$ [5] , = 3 - nucleus is soft to symmetric octupole deformations not scaled by  $\beta_2$ ;
- MEAPP = 0 - solution with the potential dependency on level energy losses in the channel, =1 - quick solution without potential dependency on level energy losses, can be used when energies of levels are much less than particle incident energy;
- MEVOL = 0 - solution without account of volume conservation, =1 - account of volume conservation in uniform nuclear density approximation[43] , =2 - common case[44] , presenting nuclear density distribution by real potential form factor;
- MEREL = 0 - calculations using non-relativistic Schrödinger formalism, =1 -account of relativistic kinematics and potential dependence, =2 - account of relativistic kinematic only, =3 - account of relativistic kinematics and real potential dependence;
- MECUL = 0 - solution with Coulomb correction potential proportional to the derivative of real potential dependence (*MANUAL* Eq. (94)), = 1 - Coulomb correction is energy independent and applied to real potential, =2 - Lane consistent Coulomb correction with "effective" interacting proton energy equal  $E - C_{Coul} \frac{ZZ'e^2}{A^{1/3}}$  (Eq. (16)), applied to real and imaginary potentials, excluding SO potential, =3 - same change of "effective" interacting proton energy as for MECUL=2, is applied only to real potential;
- MERZZ = 0 - charge radius is considered constant, = 1 - charge radius is considered energy dependent (*MANUAL* Eq. (96));
- MERRR = 0 - real potential radius is considered constant, = 1 - charge radius is considered energy dependent (*MANUAL* Eq. (95));
- MEDIS = 0 - without account of dispersion relationships between real and imaginary potential terms, = 1 - with account of dispersion relationships between all real and imaginary potentials terms, =2 - with account of dispersion relationships between all real and imaginary potentials terms, excluding SO potential;
- MERIP = 0 - for CS calculation option (MEJOB = 1) potential parameters are read only at one energy and the potential is calculated for all other energies analytically, =1 - for CS calculation option (MEJOB = 1) code reads potential blocks for each requested energy, this option is used in the RIPL [45] interface. If MEJOB = 2 (adjustment) MERIP is not used.

- Record 8 - FORMAT(E12.7,4I2) for MEHAM = 1, FORMAT(E12.7,7I2) for MEHAM > 1

Characteristics of nuclear levels, NCA (I) -partition level characteristic is added.

(EL (I) , JO (I) , NPO (I) , KO (I) , NCA (I) , I=1,NUR) - Cards 8a, 8b, .. for MEHAM =1

(EL (I) , JO (I) , NPO (I) , NTU (I) , NNB (I) , NNG (I) , NNO (I) , NCA (I) , I=1,NUR) -  
Cards 8a, 8b, .. for MEHAM > 1

- EL (I) -energy of the  $I^{th}$  level;
- JO (I) - spin of the level multiplied by two;
- NPO (I) - parity of the level = +1 - for positive, = -1 - for negative;
- KO (I) -  $K$  of the levels band multiplied by two;
- NTU (I) - the number of rotational energy solution  $\tau$ ;
- NNB (I) - the number of  $\beta_2$  oscillation function solution  $n_{\beta_2}$ ;
- NNG (I) - the number of  $\gamma$  oscillation function solution  $n_{\gamma}$ ;
- NNO (I) - the number of  $\beta_3$  oscillation function solution  $n_{\beta_3}$ ;
- NCA (I) - key determining partitions: = 0 for levels of interacting nuclei, =1 - for isobaric analog levels of  $A(N - 1, Z + 1)$  excited in  $p + A(N, Z) \rightarrow n' + A(N - 1, Z + 1)$  reaction.

● Record 10 - FORMAT(6E12.7)

Optical potential parameters, for details see Section 6 of the MANUAL. New potential parameters are added to the card 10i and one more card 10j is added to the potential input card

VR0, VR1, VR2, VR3, VRLA, ALAVR - card 10a

- VR0 - real optical potential  $V_R$  constant term  $V_R^0$ ;
- VR1 - real optical potential  $V_R$  linear term  $V_R^1$ ;
- VR2 - real optical potential  $V_R$  square term  $V_R^2$ ;
- VR3 - real optical potential  $V_R$  cubic term  $V_R^3$ ;
- VRLA - real optical potential  $V_R$  exponential term  $V_R^{DISP}$ ;
- ALAVR - real optical potential  $V_R$  exponential constant  $\lambda_R$ .

WD0, WD1, WDA1, WDBW, WDWID, ALAWD - card 10b

- WD0 - imaginary surface potential  $W_D$  constant term  $W_D^0$ ;
- WD1 - imaginary surface potential  $W_D$  linear term  $W_D^1$ ;



- WDA1 - imaginary surface potential  $W_D$  linear term  $W_D^1$  for the projectile energies above  $E_{change}$  ;
- WDBW - imaginary surface potential  $W_D$  term  $W_D^{DISP}$  ;
- WDWID - imaginary surface potential  $W_D$  dispersion width  $WID_D$  ;
- ALAWD - imaginary surface potential  $W_D$  exponential constant  $\lambda_D$ .

WC0, WC1, WCA1, WCBW, WCWID, BNDC - card 10c

- WC0 - imaginary volume potential  $W_V$  constant term  $W_V^0$  ;
- WC1 - imaginary volume potential  $W_V$  linear term  $W_V^1$  ;
- WCA1 - imaginary volume potential  $W_V$  linear term  $W_V^1$  for the projectile energies above  $E_{change}$  ;
- WCBW - imaginary volume potential  $W_V$  term  $W_V^{DISP}$  ;
- WCWID - imaginary volume potential  $W_V$  dispersion width  $WID_V$  ;
- BNDC - boundary energy where linear potential slopes change  $E_{change}$ .

VS, ALASO, WSO, WS1, WSBW, WSWID - card 10d

- VS - real spin orbit potential  $V_{SO}^0$  ;
- ALASO - real spin orbit potential  $V_{SO}$  exponential constant  $\lambda_{so}$  ;
- WSO - imaginary spin orbit potential  $W_{SO}$  constant term  $W_{SO}^0$  ;
- WS1 - imaginary spin orbit potential  $W_{SO}$  linear term  $W_{SO}^1$  ;
- WSBW - imaginary spin orbit potential  $W_{SO}$  term  $W_{SO}^{DISP}$  ;
- WSWID - imaginary spin orbit potential  $W_{SO}$  dispersion width  $WID_{SO}$ .

RR, RRBWC, RRWID, PDIS, AR0, AR1 - card 10e

- RR - real potential radius  $R_R$  ;
- RRBWC - real potential radius  $R_R$  dispersion constant  $C_R$  ;
- RRWID - real potential radius  $R_R$  dispersion width  $WID_R$  ;
- PDIS - dispersion power constant  $S$  ;
- AR0 - real potential diffuseness  $a_R$  constant term  $a_R^0$  ;
- AR1 - real potential diffuseness  $a_R$  linear term  $a_R^1$ .

RD, AD0, AD1, RC, AC0, AC1 - card 10f

- RD - imaginary surface potential radius  $R_D$ ;
- AD0 - imaginary surface potential diffuseness  $a_D$  constant term  $a_D^0$ ;
- AD1 - imaginary surface potential diffuseness  $a_D$  linear term  $a_D^1$ ;
- RC - imaginary volume potential radius  $R_V$ ;
- AC0 - imaginary volume potential diffuseness  $a_V$  constant term  $a_V^0$ ;
- AC1 - imaginary volume potential diffuseness  $a_V$  linear term  $a_V^1$ .

RW, AW0, AW1, RS, AS0, AS1 - card 10g

- RW - imaginary Gaussian potential radius  $R_W$ ;
- AW0 - imaginary Gaussian potential diffuseness  $a_W$  constant term  $a_W^0$ ;
- AW1 - imaginary Gaussian potential diffuseness  $a_W$  linear term  $a_W^1$ ;
- RS - spin orbit potential radius  $R_{SO}$ ;
- AS0 - spin orbit potential diffuseness  $a_{SO}$  constant term  $a_{SO}^0$ ;
- AS1 - spin orbit potential diffuseness  $a_{SO}$  linear term  $a_{SO}^1$ .

RZ, RZBWC, RZWID, AZ, CCOUL, ALF - card 10h

- RZ - equivalent charged ellipsoid radius  $R_C$ ;
- RZBWC - equivalent charged ellipsoid radius  $R_C$  dispersion constant  $C_C$ ;
- RZWID - equivalent charged ellipsoid radius  $R_C$  dispersion width  $WID_C$ ;
- AZ - charged ellipsoid radius diffuseness  $a_Z$ ;
- CCOUL - Coulomb correction constant  $C_{Coul}$ ;
- ALF- mixture coefficient for imaginary volume and Gaussian potentials  $\alpha$ .

CISO, WCISO, WDISO, EA, WDSHI, WDWID2 - card 10i

- CISO - constant for real potential isospin term  $C_{viso}$ ;
- WCISO - constant for imaginary surface potential isospin term  $C_{wiso}$ ;
- WDISO - constant for imaginary volume potential isospin term  $C_{wviso}$ ;
- EA - energy constant  $E_a$  at which symmetry of volume absorption potential is broken;
- WDSHI - energy shift constant  $E_{SHIFT}$  for surface potential formulae;
- WDWID2 - width of the second Breit-Wigner component in the imaginary surface potential

ALFNEW, VRD - card 10j

- ALFNEW - non-locality constant  $\alpha$  in Eq.(7,8);
- VRD - constant not in use by the moment, will be used to create bottle-neck shape of the real potential at high incident energies

**This block of potential parameter cards should be repeated NST times if MERIP = 1, as than these parameters for each energy should be created by RIPL interface.**

**Below are the changes in input cards that are used when potential parameters adjustment is desired (MEJOB = 1).**

- Record 12 - FORMAT(6I2)

Input of flags that determine which parameters will be adjusted; if  $NPJ(I) = 1$  the chosen parameters will be adjusted. The number of the parameters to be adjusted must not exceed 20. The adjustment of optical model parameters WDSHI, WDWID2, ALFNEW and VRD is added in this code's version. Please note that the sequence of inputted flags follows the sequence of optical potential parameters for users convenience. Please also note that comparing with *MANUAL* version  $NPJ(51)$  was changed to  $NPJ(56)$ . We show below the sequence of all  $NPJ(I)$  with adjusted parameters.

( $NPJ(I)$ ,  $I=1,64$ ) cards 12a,...12k

- $NPJ(1)$  - flag for real optical potential  $V_R$  constant term adjustment VR0;
- $NPJ(2)$  - flag for real optical potential  $V_R$  linear term adjustment VR1;
- $NPJ(3)$  - flag for real optical potential  $V_R$  square term adjustment VR2;
- $NPJ(4)$  - flag for real optical potential  $V_R$  cubic term adjustment VR3;
- $NPJ(5)$  - flag for real optical potential  $V_R$  exponential term  $V_R^{DISP}$  adjustment VRLA;
- $NPJ(6)$  - flag for real optical potential  $V_R$  exponential exponential constant  $\lambda_R$  adjustment ALAVR;
- $NPJ(7)$  - flag for imaginary surface potential  $W_D$  constant term adjustment WD0;
- $NPJ(8)$  - flag for imaginary surface potential  $W_D$  linear term adjustment WD1;
- $NPJ(9)$  - flag for imaginary surface potential  $W_D$  linear term  $W_D^1$  for particle incident energies above  $BNDC=E_{change}$  adjustment WDA1;

- NPJ (10) - flag for imaginary surface potential  $W_D$  dispersion term  $W_D^{DISP}$  adjustment WDBW;
- NPJ (11) - flag for imaginary surface potential  $W_D$  dispersion width term  $WID_D$  adjustment WDWID;
- NPJ (12) - flag for imaginary surface potential  $W_D$  exponential constant  $\lambda_D$  adjustment ALAWD;
- NPJ (13) - flag for imaginary volume potential  $W_V$  constant term linear term  $W_V^0$  adjustment WC0;
- NPJ (14) - flag for imaginary volume potential  $W_V$  linear term  $W_V^1$  adjustment WC1;
- NPJ (15) - flag for imaginary volume potential  $W_V$  linear term  $W_V^1$  for particle incident energies above  $BNDC=E_{change}$  adjustment WCA1;
- NPJ (16) - flag for imaginary surface potential  $W_V$  dispersion term  $W_V^{DISP}$  adjustment WCBW;
- NPJ (17) - flag for imaginary surface potential  $W_V$  dispersion width term  $WID_V$  adjustment WCWID;
- NPJ (18) - flag for energy where linear potential slopes change  $E_{change}$  adjustment BNDC;
- NPJ (19) - flag for real spin-orbit potential  $V_{SO}$  adjustment VS;
- NPJ (20) - flag for real spin-orbit potential  $V_{SO}$  exponential constant  $\lambda_{so}$  adjustment ALASO;
- NPJ (21) - flag for imaginary spin-orbit potential  $W_{SO}$  constant term  $W_{SO}^0$  adjustment WSO;
- NPJ (22) - flag for imaginary spin-orbit potential  $W_{SO}$  linear term  $W_{SO}^1$  adjustment WS1;
- NPJ (23) - flag for imaginary spin-orbit potential  $W_{SO}$  dispersion term  $W_{SO}^{DISP}$  adjustment WSBW;
- NPJ (24) - flag for imaginary spin-orbit potential  $W_{SO}$  dispersion dispersion width  $WID_{SO}$  adjustment WSWID;
- NPJ (25) - flag for real potential radius  $R_R$  adjustment RR;
- NPJ (26) - flag for real potential radius  $R_R$  dispersion constant  $C_R$  adjustment RRBWC;
- NPJ (27) - flag for real potential radius  $R_R$  dispersion dispersion width  $WID_R$  adjustment RRWID;

- NPJ (28) - flag for dispersion power constant  $S$  adjustment PDIS;
- NPJ (29) - flag for real potential diffuseness  $a_R$  constant term adjustment AR0;
- NPJ (30) - flag for real potential diffuseness  $a_R$  linear term  $a_R^1$  adjustment AR1;
- NPJ (31) - flag for imaginary surface potential radius  $R_D$  adjustment RD;
- NPJ (32) - flag for imaginary surface potential diffuseness  $a_D$  constant term adjustment AD0;
- NPJ (33) - flag for imaginary surface potential diffuseness  $a_D$  linear term  $a_D^1$  adjustment AD1;
- NPJ (34) - flag for imaginary volume potential radius  $R_V$  adjustment RC;
- NPJ (35) - flag for imaginary volume potential diffuseness  $a_V$  constant term adjustment AC0;
- NPJ (36) - flag for imaginary volume potential diffuseness  $a_V$  linear term  $a_V^1$  adjustment AC1;
- NPJ (37) - flag for Gaussian potential radius  $R_W$  adjustment RW;
- NPJ (38) - flag for Gaussian potential radius diffuseness  $a_W$  constant term adjustment AW0;
- NPJ (39) - flag for Gaussian potential radius diffuseness  $a_W$  linear term  $a_W^1$  adjustment AW1;
- NPJ (40) - flag for spin-orbit potential radius  $R_{SO}$  adjustment RS;
- NPJ (41) - flag for spin-orbit potential diffuseness  $a_{SO}$  constant term adjustment AS0;
- NPJ (42) - flag for spin-orbit potential diffuseness  $a_{SO}$  linear term  $a_{SO}^1$  adjustment AS1;
- NPJ (43) - flag for equivalent charged ellipsoid radius  $R_C$  adjustment RZ;
- NPJ (44) - flag for equivalent charged ellipsoid radius  $R_C$  dispersion constant  $C_C$  adjustment RZBWC;
- NPJ (45) - flag for equivalent charged ellipsoid radius  $R_C$  dispersion width  $WID_C$  adjustment RZWID;
- NPJ (46) - flag for equivalent charged ellipsoid radius  $R_C$  diffuseness  $a_Z$  adjustment AZ;
- NPJ (47) - flag for Coulomb correction constant adjustment CCOUL;
- NPJ (48) - flag for mixture coefficient  $\alpha$  for imaginary volume and Gaussian potentials adjustment ALF;

- NPJ ( 49 ) - flag for real potential isospin constant  $C_{viso}$  adjustment CISO;
- NPJ ( 50 ) - flag for imaginary surface potential isospin constant  $C_{wiso}$  adjustment WCISO;
- NPJ ( 51 ) - flag for imaginary volume potential isospin constant  $C_{wviso}$  adjustment WDISO;
- NPJ ( 52 ) - flag for energy constant  $E_a$  at which symmetry of volume absorption potential is broken adjustment EA;
- NPJ ( 53 ) - flag for energy shift constant  $E_{SHIFT}$  for surface potential formulae adjustment WDSHI;
- NPJ ( 54 ) - flag for width of the second BW component in imaginary shape potential adjustment WDWID2;
- NPJ ( 55 ) - flag for constant  $\alpha$  in Eq.(7,8) adjustment ALFNEW;
- NPJ ( 56 ) - flag for constant that later, will be used to create bottle shape of the real potential at high incident energy adjustment VRD;
- NPJ ( 57 ) - flag for equilibrium deformation  $\beta_{20}$  for non-axial Hamiltonian model adjustment BET0;
- NPJ ( 58 ) - flag for equilibrium deformation  $\beta_{30}$  for non-axial Hamiltonian adjustment BET3;
- NPJ ( 59 ) - flag for rigid deformation  $\beta_4$  for non-axial Hamiltonian model adjustment BET4;
- NPJ ( 60 ) - flag for rotational model axial rigid deformation  $\beta_{20}$  adjustment BET ( 2 ) ;
- NPJ ( 61 ) - flag for rotational model axial rigid deformation  $\beta_{40}$  adjustment BET ( 4 ) ;
- NPJ ( 62 ) - flag for rotational model axial rigid deformation  $\beta_{60}$  adjustment BET ( 6 ) ;
- NPJ ( 63 ) - flag for nuclear softness  $\mu_{\beta_3} = \mu_\epsilon$  adjustment AMUO;
- NPJ ( 64 ) - flag for nuclear softness  $\mu_{\gamma_0}$  adjustment AMG0.

**Next block of cards will be read for each energy point EE ( I ) , that means NST blocks.**

**This is the last card in input.**

### 3.2 Example of Input File

We discuss a test input for a rigid rotator case with potential adjustment. Differences between adjustment and cross section calculations options and between rigid and soft rotator options of the code are the same as in the previous version of the code and are well described in *MANUAL*.

### 3.2.1 Input for the Rigid-rotator Model

Here one can find a description of an actual input file that was used for adjustment of  $^{181}\text{Ta}$  optical model parameters, which is attached below.

- Card 1 - text information about code running.
- Card 2 - defines that a run with adjustment of optical potential parameters (MEJOB=2) for rigid-rotator CC-model (MEPOT=1) is given; the rotational model nuclear Hamiltonian will be used to determine coupling (MEHAM=1); output will be the shortest possible (MEPRI=0); code will use accurate ten channels coupled solution as zero approximation for iteration method for coupled equations system solution (MESOL=10); switches MESSHA, MESSHO, MEHAO are set to zero, but will be ignored as MEPOT=1; energy dependence of optical potential in different channels caused by energy losses due to levels excitation will be ignored in this run (MEAPP=0); MEVOL is set zero, but will be ignored as MEPOT=1; relativistic generalization will be taken into account by relativistic kinematics, but potential would not be multiplied by  $K(E)$  factor (MEREL=2); Coulomb correction will be taken by using "effective energy decrease due to proton repulsion by nuclei for a real potential only (MECUL=3); Coulomb and real potential radii will be considered energy independent (MERZZ= MERRR=0); with account of dispersion relationships between all real and imaginary potentials terms (MEDIS=1); MERIP=1 key is ignored as (MEJOB=2).
- Card 3 - defines that eight levels are coupled (NUR=8); potential will be adjusted using experimental data at one energy point simultaneously (NST=1); the highest order  $L$  for radius deformation expansion is chosen to be equal 6 (NPD=6); resulting optical potential will be expanded up to  $Y_{80}$  spherical harmonic (LAS=8); MTET is set zero, but will be ignored as MEJOB=2, waves with angular momentum greater than 50 will be ignored (LLMA=50); coupling of no more than 100 equations will be taken into account, other ignored (NCMA=100); no more than 100  $J^\pi$  states ordered by increasing  $J$  will be taken into consideration (NSMA=100); coupled equations will be ordered one by one, starting with the first level, then the second and so on (KODMA=0).
- Card 4 - one energy, as it was stated in Card 3 (NST=1), in which adjustment of optical potential to experimental data will be organized in the code.
- Card 5 - one MCHAE (I) flag (NST=1 in Card 3), in our case equal 1 (experimental data for incident proton).
- Cards 6-13 - excitation energies and quantum numbers of the height levels (NUR=8 in Card 3) that should be coupled first four levels are from  $^{181}\text{Ta}$  GS rotational band (NCA (1-4) = 0) and

four levels from fifth to eighth are isobaric analog states of  $^{181}\text{W}$  ( $\text{NCA}(5-8)=1$ ).

- Card 14 - quantities of incident particle and interacting nucleus.
- Cards 15-24 - starting set of optical potential parameters, even in case some values of the potential are equal to zero their radii and diffuseness must be set non-zero.
- Card 25 -  $\beta_2$ ,  $\beta_4$  and  $\beta_6$  as defined by  $\text{NPD}=6$  on Card 3.
- Cards 26-36 -  $\text{NPJ}(\text{I})$  values; in our case only  $\text{NPJ}(49)$ ,  $\text{NPJ}(55)$  and  $\text{NPJ}(61)$  are equal 1:  $\text{CISO}$ ,  $a$  constants and  $\beta_{40}$  deformation are to be adjusted in this run.

Next cards 37-55 apply to the experimental data used for adjustment at the first energy point.

- Card 37 - information about experimental data at the first experimental energy point ( $E_n=27$  MeV), that will be used in optical parameters adjustment; total cross section is not used in adjustment at this energy point ( $\text{NT}(1)=0$ ); reaction cross section is not used in adjustment at this energy ( $\text{NR}(1)=0$ ); no experimental integral scattering cross section with excitation of a group of levels will be used for adjustment ( $\text{NGN}(1)=0$ ), while one group of angular distributions ( $\text{NGD}(1)=1$ ); strength functions  $S_0$  and  $S_1$  will not be adjusted at this energy ( $\text{NSF1}(1)=\text{NSF2}(1)=0$ ).
- Card 38 - total and reaction cross sections and their errors at first energy; as  $\text{NT}(1)$  and  $\text{NR}(1)$  are equal to zero so this card is blank.
- Card 39 - information about experimental angular distributions: initial and final excited levels of the group and number of experimental angles ( $\text{NID}(1, \text{K})=5$ ,  $\text{NFD}(1, \text{K})=8$ ,  $\text{MTD}(1, \text{K})=32$ ,  $\text{K}=1$ ); in our case  $\text{NID}$  and  $\text{NFD}$  do not coincide, it means that experimental angular distributions present excitation of a group of levels from fifth to eighth (isobaric analog states levels).
- Cards 40-55 - experimental angular distributions as angle, data, error. .
- Card 56 - three values of absolute accuracies for adjusted parameters as defined by  $\text{NPJ}(\text{I})$  in Cards 26-36.
- Card 57 - zero (initial  $\chi^2$  is unknown).

The actual rigid-rotator input file for adjustment of  $^{181}\text{Ta}$  optical model parameters using experimental angular distribution of emitted neutrons in the  $p + ^{181}\text{Ta} \rightarrow n' + ^{181}\text{W}$  reaction

---



```

1  Ta-181  rigid -dispersive
2  02010100100000000000020300000101
3  008001006008000050100100000    NUR NST NPD LAS MTET LLMA NCMA NSMA KODMA
4  0.270000+02
5  01
6  0.0000000-0007+10700    LEVELS OF Ta-181
7  0.1362660-0009+10700    EL(I),JO(I),NPO(I),KO(I),NCA(I)
8  0.30162  -0011+10700
9  0.49519  -0013+10700
10 1.7630000+0107+10701    IOS LEVELS OF W-181
11 1.7796266+0109+10701    EL(I),JO(I),NPO(I),KO(I),NCA(I)
12 1.7931620+0111+10701
13 1.8125190+0113+10701
14 1.00866520 .5000000-00 .1810000+03 .7300000+02-.6819980+01-.6517450+01
15 -.0000000+02 .0000000+00 .0000000+00 .1080000+01 .5030000+02 .9130000-02
16 .0000000+01-.0000000-00 .0000000+01 .2213500+02 .1469000+02 .0961000-01
17 .0000000+01-.0000000-00 .0000000+01 .1076000+02 .8384000+02 .0000000+01
18 .6040000+01 .0050000000 .0000000+01-.0000000-00 -3.10 0 160.
19 .1239450+01 .00 .1000000+03 2.00 .6295000-00 .0000000-03
20 .1240000+01 .5085000-00 .0000000-02 .1077100+01 .6870000+00 .0000000-02
21 .1000000+01 .1000000+01 .0000000-00 .1129300+01 .5900000-00-.0000000-02
22 .1062600+01 .00 .1000000+01 .640 1.3300 .1000000+01
23 .2100000+02 .2940000+02 .0000000+03 052.0 500000. .0
24 1.080000+00 .0
25 0.251000+00-0.07900E-00-0.50000E-01
26 000000000000
27 000000000000
28 000000000000
29 000000000000
30 000000000000
31 000000000000
32 000000000000
33 000000000000
34 010000000000
35 010000000000
36 01000000
37 000000010000    p,n data from here !!!!!!!
38 .0000000+01 .0000000-01 .0000000-00 .0000000-00    27MeV, No total
39 050832
40 2.50000E+00 1.46000E-03 3.65000E-04 7.90000E+00 1.82000E-03 4.55000E-04
41 1.61000E+01 2.10000E-03 5.25000E-04 2.44000E+01 2.30000E-03 5.75000E-04
42 3.26000E+01 2.40000E-03 6.00000E-04 3.83000E+01 1.59000E-03 3.97500E-04
43 4.66000E+01 1.03000E-03 2.57500E-04 5.50000E+01 8.90000E-04 2.22500E-04
44 6.24000E+01 7.40000E-04 1.85000E-04 7.16000E+01 6.10000E-04 1.52500E-04
45 8.57000E+01 2.30000E-04 5.75000E-05 9.94000E+01 2.30000E-04 5.75000E-05
46 1.16100E+02 1.30000E-04 3.25000E-05 1.29900E+02 1.09000E-04 2.72500E-05

```

```

47  1.46400E+02  1.40000E-04  3.50000E-05  1.60200E+02  1.17000E-04  2.92500E-05
48  8.60000E+00  2.01000E-03  2.61300E-04  1.66000E+01  1.54000E-03  1.54000E-04
49  1.67000E+01  1.30000E-03  5.20000E-04  2.41000E+01  2.54000E-03  2.03200E-04
50  3.16000E+01  2.31000E-03  1.84800E-04  3.21000E+01  2.09000E-03  2.09000E-04
51  4.55000E+01  9.50000E-04  1.23500E-04  5.35000E+01  9.60000E-04  1.34400E-04
52  5.41000E+01  6.90000E-04  1.10400E-04  6.10000E+01  8.90000E-04  8.90000E-05
53  6.90000E+01  5.20000E-04  8.84000E-05  7.01000E+01  5.90000E-04  8.85000E-05
54  8.24000E+01  4.20000E-04  1.09200E-04  9.95000E+01  3.80000E-04  9.12000E-05
55  1.13400E+02  2.80000E-04  7.00000E-05  1.28900E+02  1.90000E-04  7.60000E-05
56  .2000000-03  .2000000-03  .100000-03
57  0.0

```

---

### 3.3 Running the Code OPTMAN and Description of Output Files

For describing the test of the OPTMAN code we took the same soft-rotator  $^{12}\text{C}$  non-dispersive potential input, that has been used in *MANUAL* calculations (see pages 58-59 of *MANUAL*). To allow running with the new version of the code we needed to add keys MEDIS=0 and MERIP=0 in the second line of input and add several new potential parameters to the last card (WDISO=0.0 - imaginary volume isovector term equal zero, EA=52 - ignored in case of non-dispersive calculations, WDSHI=500 - allowing to use exponential decrease of surface imaginary potential, and ignoring the alternative formulation for it, WDWID2=0 - not used) of old potential and add one more card in potential description (ALFNEW=1.65, VRD=0.0), which both have not been used by the code, as it is not dispersive case, but should be inputted. As one runs the code it will first ask the name of the input file and, after it will be read, the code asks to give the name for the output file, in which all the calculated information will be stored. While running, eleven more files are created by the code: CR-SECT, TRANSME, GNASH, ANG-DIST, ANG-POL, OPTMAN.CS, OPTMAN.ANG, OPTMAN.ICS, OPTMAN.TL, OPTMAN.LEG and ANGL-DIS.YW. First five have been described in *MANUAL*, and here we just repeat this description for users convenience. The file CR-SECT contains one line for each calculated energy. Line starts with a value of energy, and then follow total cross sections, reaction cross sections and level excitation cross sections ordered by their excitation energies (in case calculations are made for protons total and elastic cross-sections being infinity include contribution of nuclear amplitude only). The file TRANSME contains nucleon transmission coefficients for each calculated energy to be used for the code STAPRE [46]. For each calculated energy the first line indicates incident energy and  $l_{\text{max}}$  - maximum angular momentum for which nucleon transmission coefficients are available for this energy, and then lines with  $l_{\text{max}} + 1$  transmissions follow ordered by growing angular momentum starting with  $T_{l=0}$  by six in a line. As the calculated transmission coefficients are intended to be used for the code STAPRE [46] they are aver-

aged over  $J$  dependence and depend only on angular momentum  $l$ . Many zero value higher angular momentum transmissions are outputted by code, it is convenient for STAPRE code calculations. The file GNASH contains nucleon transmission coefficients for each calculated energy to be used for the code GNASH [47]. For each energy the first line indicates the energy and  $l_{\max}$  - maximum angular momentum for which nucleon transmission coefficients are available for this energy, and according with the GNASH format lines with transmissions: first  $T_{l=0,j=1/2}$ , than pares  $T_{l,j=l-1/2}$ ,  $T_{l,j=1/2}$  follow ordered by growing angular momentum up to  $l_{\max}$  by six in a line. The file ANGL-DIST contains calculated angular distributions for all excited levels for energies described by Record 5 at angles distinguished by Record 7 of input file. The file ANGLE-DIST consists of NST block: the first card of the block indicated incident energy and type of the particle, than follow MTET lines with first angle and than angular distributions for excited levels in order following the order of levels input. The file ANGL-POL gives polynomial expansion for angular distributions determined by *MANUAL* Eqs. (87,88). This file consists of NST blocks. Each block starts with line describing incident energy and numbers of coefficients in expansions  $LKK=NL$  and  $LCOUL=ML$  (see *MANUAL* Eqs. (87,88). Second line of the block gives total cross sections, reaction cross sections and level excitation cross sections ordered by their excitation energies, than for each excited level follow  $NL$  of  $a_L(E)$  coefficients six in line and for elastically scattered protons  $ML$  pares of  $Reb_L(E)$  and  $Imb_L(E)$  coefficients. When one runs optical parameter adjustment files CR-SECT, TRANSME, GNASH, ANGL-DIST and ANGL-POL are organized but are blank. When one runs optical calculations code's option with MTET=0 input files ANGL-DIST, ANGL-POL are also blank. New files OPTMAN.CS, OPTMAN.ANG, OPTMAN.ICS, OPTMAN.TL and OPTMAN.LEG are the files create to be used by the EMPIRE nuclear data evaluation system [48], which are compatible with ECIS output files and are described in ECIS manual. The file ANGLDIS-YW is similar to the previously created ANGL-DIST and is being used for some specific purposes. Below we present these files for a run in which optical calculations of  $^{12}\text{C}$  cross sections are organized, as well as the appropriate input file. User can compare input file demonstrated here with the one in *MANUAL* for the same case to see the necessary changes. Small changes in calculated values of outputted files, compared with the *MANUAL* ones are due to the new relativistic generalization applied here, which is now completely consistent with the one used in the ECIS code.

#### INPUT file

---

```

1 C-12   OCTUPOLE CALCULATION bet3=ksi*bet2
2 01020500010402020000010001010000
3 0.42825E+01 0.27943E+01 0.11860E+00 0.32107E+00 0.16400E+00 0.12600E+00
4 0.14694E+00 0.21419E-03 0.52045E+00 0.24200E+00 0.21005E-01 0.83263E+00
5 0.35000E+00 0.22130E+00 0.00000E+00 0.17221E+01 0.00000E+00
6 005002004004040050100100001
```

```

7 .2090000+02 .3970000+02
8 0001
9 .1000000+01 .5000000+01 .1000000+02 .1500000+02 .1800000+02 .2100000+02
10 .2400000+02 .2700000+02 .3000000+02 .3300000+02 .3600000+02 .3900000+02
11 .4200000+02 .4500000+02 .4800000+02 .5100000+02 .5400000+02 .5600000+02
12 .5900000+02 .6200000+02 .6500000+02 .7000000+02 .7500000+02 .8000000+02
13 .8500000+02 .9000000+02 .9500000+02 .1000000+03 .1050000+03 .1100000+03
14 .1150000+03 .1200000+03 .1250000+03 .1300000+03 .1350000+03 .1400000+03
15 .1500000+03 .1600000+03 .1700000+03 .1800000+03
16 .0000000-0000+10100000000
17 .4440000+0104+10100000000
18 .9640000+0106-10100000100
19 .7650000+0100+10101000000
20 .1780000+0204+10200000000
21 1.00866520 .5000000-00 .1200000+02 .6000000+01-.1183500+02-.8950000+01
22 -.3094000+02-.0000000-00 .0002381+00 .3000000-06 .8630000+02 .4532000-02
23 .0000000+01-.0000000-00 .0000000+01 .9530000+01 .1533000+02 .1755000-01
24 .0000000+01-.0000000-00 .0000000+01 .1909000+02 .9429000+02 .0000000+01
25 .7840000+01 .0030000000 .0000000+01-.0000000-00 -3.100 160.
26 .1150000+01 .10 .5000000+02 2. .5190000-00 .2600000-02
27 .1072600+01 .4460000-00 .0000000-02 .1039700+01 .4570000+00 .0000000-02
28 .1000000+01 .1000000+01 .0000000-00 .1086300+01 .6050000-00-.0000000-02
29 .1145300+01 .00 .1000000+01 .680 1.1000 .1000000+01
30 .1300000+02 .3000000+02 .0000000+03 52.00 500. .0
31 1.65 .0

```

---

## OUTPUT file

---

```

1 C-12 OCTUPOLE CALCULATION bet3=ksi*bet2
2 INTERACTION OF PARTICLE, HAVING SPIN = 0.50
3 WITH NUCLEI A= 12.0000000
4 COUPLED CHANNELS METHOD
5 RELATIVISTIC KINEMATICS AND POTENTIAL DEPENDENCE
6 OPTICAL POTENTIAL WITHOUT DISPERSIVE RELATIONSHIPS
7 COULOMB CORRECTION PROPORTIONAL REAL POTENTIAL DER-VE
8 CHARGE RADIUS IS ENERGY DEPENDENT
9 REAL RADIUS IS ENERGY DEPENDENT
10 WITH AC. NONAXIAL GEXADECAPOLE DEFORMATIONS
11
12 HAMILTONIAN 5PA0 POTENTIAL EXPANDED BY BET0
13 WITH AC. NONAXIAL OCTUPOLE SOFT DEFORMATIONS
14
15 NUMBER OF COUPLED LEVELS= 5 NPD = 4
16 NUMBER OF TERMS IN POTENTIAL EXPANSION= 4
17
18
19 ENERGY LEVEL'S SPIN*2 NTU NNB NNG NNO NPO
20

```



79 JU1= 3 JU2= 3 NNT= 4 FOV(JU1, JU2, NNT) = 0.7459344D+03 ANU1= 0.7969259D-02 ANU2= 0.7969259D-02  
80 JU1= 3 JU2= 4 NNT= 1 FOV(JU1, JU2, NNT) = 0.3176693D+01 ANU1= 0.7969259D-02 ANU2= 0.2446976D+01  
81 JU1= 3 JU2= 4 NNT= 2 FOV(JU1, JU2, NNT) = 0.2092894D+02 ANU1= 0.7969259D-02 ANU2= 0.2446976D+01  
82 JU1= 3 JU2= 4 NNT= 3 FOV(JU1, JU2, NNT) = 0.1356776D+03 ANU1= 0.7969259D-02 ANU2= 0.2446976D+01  
83 JU1= 3 JU2= 4 NNT= 4 FOV(JU1, JU2, NNT) = 0.8938593D+03 ANU1= 0.7969259D-02 ANU2= 0.2446976D+01  
84 JU1= 3 JU2= 5 NNT= 1 FOV(JU1, JU2, NNT) = 0.4816125D+01 ANU1= 0.7969259D-02 ANU2= 0.2785297D-02  
85 JU1= 3 JU2= 5 NNT= 2 FOV(JU1, JU2, NNT) = 0.2551115D+02 ANU1= 0.7969259D-02 ANU2= 0.2785297D-02  
86 JU1= 3 JU2= 5 NNT= 3 FOV(JU1, JU2, NNT) = 0.1443104D+03 ANU1= 0.7969259D-02 ANU2= 0.2785297D-02  
87 JU1= 3 JU2= 5 NNT= 4 FOV(JU1, JU2, NNT) = 0.8626245D+03 ANU1= 0.7969259D-02 ANU2= 0.2785297D-02  
88 JU1= 4 JU2= 4 NNT= 1 FOV(JU1, JU2, NNT) = 0.4940767D+01 ANU1= 0.2446976D+01 ANU2= 0.2446976D+01  
89 JU1= 4 JU2= 4 NNT= 2 FOV(JU1, JU2, NNT) = 0.2992855D+02 ANU1= 0.2446976D+01 ANU2= 0.2446976D+01  
90 JU1= 4 JU2= 4 NNT= 3 FOV(JU1, JU2, NNT) = 0.1981755D+03 ANU1= 0.2446976D+01 ANU2= 0.2446976D+01  
91 JU1= 4 JU2= 4 NNT= 4 FOV(JU1, JU2, NNT) = 0.1380222D+04 ANU1= 0.2446976D+01 ANU2= 0.2446976D+01  
92 JU1= 4 JU2= 5 NNT= 1 FOV(JU1, JU2, NNT) = 0.3838257D+01 ANU1= 0.2446976D+01 ANU2= 0.2785297D-02  
93 JU1= 4 JU2= 5 NNT= 2 FOV(JU1, JU2, NNT) = 0.2517325D+02 ANU1= 0.2446976D+01 ANU2= 0.2785297D-02  
94 JU1= 4 JU2= 5 NNT= 3 FOV(JU1, JU2, NNT) = 0.1655250D+03 ANU1= 0.2446976D+01 ANU2= 0.2785297D-02  
95 JU1= 4 JU2= 5 NNT= 4 FOV(JU1, JU2, NNT) = 0.1113186D+04 ANU1= 0.2446976D+01 ANU2= 0.2785297D-02  
96 JU1= 5 JU2= 5 NNT= 1 FOV(JU1, JU2, NNT) = 0.5108385D+01 ANU1= 0.2785297D-02 ANU2= 0.2785297D-02  
97 JU1= 5 JU2= 5 NNT= 2 FOV(JU1, JU2, NNT) = 0.2815616D+02 ANU1= 0.2785297D-02 ANU2= 0.2785297D-02  
98 JU1= 5 JU2= 5 NNT= 3 FOV(JU1, JU2, NNT) = 0.1650190D+03 ANU1= 0.2785297D-02 ANU2= 0.2785297D-02  
99 JU1= 5 JU2= 5 NNT= 4 FOV(JU1, JU2, NNT) = 0.1018772D+04 ANU1= 0.2785297D-02 ANU2= 0.2785297D-02

100

POTENTIAL PARAMETERS V(R)

101

102  
103 VR0=-30.940 VR1= 0.0000 VR2= 0.0002381 RR= 1.1500 AR0= 0.5190 AR1= 0.0026  
104 WD0= 0.0000 WD1= 0.0000 VR3= 0.0000003 RD= 1.0726 AD0= 0.4460 AD1= 0.0000  
105 WC0= 0.0000 WC1= 0.0000 RC= 1.0397 AC0= 0.4570 AC1= 0.0000  
106 RW= 1.0000 AW0= 1.0000 AW1= 0.0000  
107 VSO= 7.8400 RS= 1.0863 AS0= 0.6050 AS1= 0.0000  
108 ALF= 1.0000 ANEU= 1.0087 RZ= 1.1453 AZ0= 0.6800  
109 BNDC= 0.00 WDA1= 0.0000 WCA1= 0.0000 CCOUL= 1.1000 CISO= 13.000 WCISO= 30.000  
110 WS0= 0.0000 WS1= 0.0000 VRLA=86.3000 ALAVR= 0.00453 WCBW=19.0900 WCWID=94.2900  
111 WDBW= 9.5300 WDWID=15.3300 ALAWD= 0.0175 EFERMN=-11.835 EFERMP= -8.950 ALASO= 0.0030  
112 PDIS= 2.0000 WSBW=-3.1000 WSWID= 160.00 RRBWC= 0.1000 RRWID= 50.00 RZBWC= 0.0000  
113 RZWID= 1.0000 EA= 52.0000 WDISO= 0.000 WDSHI= 500.00 WDWID2= 0.00 ALFNEW= 1.650  
114 VRD= 0.000

115

NUCLEUS CHARGE = 6.0000

116

117

118

SPHERICAL VOLUME INTEGRALS OF REAL POTENTIAL F-FACTORS AND DERIVATIVES:

119 VIRO= 8.625 VIR1= 19.696 VIR2= 16.714 VIR3= 5.472 CBET0= 1.000

SPHERICAL VOLUME INTEGRALS OF IMAGINARY (WC) F-FACTORS AND DERIVATIVES:

120 WICO= 6.132 WIC1= 15.117 WIC2= 13.501 WIC3= 4.471 CBETC= 0.893

SPHERICAL VOLUME INTEGRALS OF IMAGINARY (WD) F-FACTORS AND DERIVATIVES:

121 WID0= 11.922 WID1= 21.532 WID2= 10.714 WID3= 0.080 CBETD= 0.498

SPHERICAL VOLUME INTEGRALS OF IMAGINARY (WW) F-FACTORS AND DERIVATIVES:

122 WIW0= 0.000 WIW1= 0.000 WIW2= 0.000 WIW3= 0.000 CBETW= 0.000

123

POTENTIAL VALUES FOR INCIDENT ENERGY= 20.900000 MeV:

124

VR= 43.727 VRDC(DWv)= 0.000 Dws= 0.000 Wv= 2.053

125

Ws= 4.400 VSO= 7.107 DVSO= 0.000 WSO= -0.125

126

VHF= 0.000 VCUL= 0.000 VCULC= 0.000 RELPOT= 1.011

127

128

129

130

131

132

133

134

135

136

ORB. MOMENT

TRANSITIONS

SR

SI

137					
138	0	0.5737846069	-0.2536359445	-0.4370581955	
139	1	0.5374276509	-0.4786559860	-0.3214959692	
140	2	0.6352215737	-0.3613472888	-0.0780649826	
141	3	0.5547320159	0.2268952102	0.3438143496	
142	4	0.1286411602	0.8842311291	0.2237765843	
143	5	0.0161530565	0.9873378258	0.0504127033	
144	6	0.0023943416	0.9983776698	0.0114224488	
145	7	0.0002735766	0.9998513790	0.0024095287	
146	8	0.0000196752	0.4705782607	0.0001341173	
147	9	0.0000000000	0.0000000000	0.0000000000	
148	10	0.0000000000	0.0000000000	0.0000000000	
149	11	0.0000000000	0.0000000000	0.0000000000	
150	12	0.0000000000	0.0000000000	0.0000000000	
151	13	0.0000000000	0.0000000000	0.0000000000	
152	14	0.0000000000	0.0000000000	0.0000000000	
153	15	0.0000000000	0.0000000000	0.0000000000	
154	16	0.0000000000	0.0000000000	0.0000000000	
155	17	0.0000000000	0.0000000000	0.0000000000	
156	18	0.0000000000	0.0000000000	0.0000000000	
157	19	0.0000000000	0.0000000000	0.0000000000	
158	20	0.0000000000	0.0000000000	0.0000000000	
159	21	0.0000000000	0.0000000000	0.0000000000	

160

161 ANGULAR DISTRIBUTIONS OF SCATTERED PARTICLES

162

163	0.100E+01	0.109E+01	0.165E-01	0.240E-02	0.204E-02	0.126E-03
164	0.500E+01	0.106E+01	0.162E-01	0.241E-02	0.194E-02	0.125E-03
165	0.100E+02	0.967E+00	0.154E-01	0.241E-02	0.164E-02	0.124E-03
166	0.150E+02	0.833E+00	0.142E-01	0.241E-02	0.123E-02	0.121E-03
167	0.180E+02	0.740E+00	0.135E-01	0.242E-02	0.959E-03	0.119E-03
168	0.210E+02	0.643E+00	0.129E-01	0.242E-02	0.703E-03	0.117E-03
169	0.240E+02	0.547E+00	0.123E-01	0.242E-02	0.476E-03	0.114E-03
170	0.270E+02	0.454E+00	0.119E-01	0.243E-02	0.292E-03	0.111E-03
171	0.300E+02	0.368E+00	0.115E-01	0.244E-02	0.159E-03	0.108E-03
172	0.330E+02	0.292E+00	0.112E-01	0.246E-02	0.822E-04	0.105E-03
173	0.360E+02	0.226E+00	0.110E-01	0.248E-02	0.581E-04	0.101E-03
174	0.390E+02	0.171E+00	0.108E-01	0.251E-02	0.813E-04	0.975E-04
175	0.420E+02	0.127E+00	0.105E-01	0.254E-02	0.143E-03	0.939E-04
176	0.450E+02	0.925E-01	0.102E-01	0.259E-02	0.231E-03	0.902E-04
177	0.480E+02	0.666E-01	0.976E-02	0.263E-02	0.335E-03	0.866E-04
178	0.510E+02	0.481E-01	0.925E-02	0.267E-02	0.443E-03	0.831E-04
179	0.540E+02	0.355E-01	0.866E-02	0.271E-02	0.545E-03	0.797E-04
180	0.560E+02	0.297E-01	0.823E-02	0.274E-02	0.605E-03	0.775E-04
181	0.590E+02	0.240E-01	0.754E-02	0.277E-02	0.680E-03	0.745E-04
182	0.620E+02	0.209E-01	0.683E-02	0.279E-02	0.732E-03	0.717E-04
183	0.650E+02	0.194E-01	0.611E-02	0.279E-02	0.758E-03	0.692E-04
184	0.700E+02	0.188E-01	0.500E-02	0.277E-02	0.747E-03	0.658E-04
185	0.750E+02	0.186E-01	0.407E-02	0.271E-02	0.675E-03	0.636E-04
186	0.800E+02	0.178E-01	0.336E-02	0.262E-02	0.563E-03	0.625E-04
187	0.850E+02	0.160E-01	0.289E-02	0.250E-02	0.433E-03	0.626E-04
188	0.900E+02	0.137E-01	0.263E-02	0.236E-02	0.308E-03	0.639E-04
189	0.950E+02	0.110E-01	0.253E-02	0.222E-02	0.207E-03	0.664E-04
190	0.100E+03	0.839E-02	0.252E-02	0.207E-02	0.142E-03	0.698E-04
191	0.105E+03	0.617E-02	0.255E-02	0.193E-02	0.118E-03	0.740E-04
192	0.110E+03	0.447E-02	0.260E-02	0.179E-02	0.131E-03	0.789E-04
193	0.115E+03	0.330E-02	0.261E-02	0.166E-02	0.171E-03	0.843E-04
194	0.120E+03	0.260E-02	0.259E-02	0.153E-02	0.227E-03	0.900E-04

195 0.125E+03 0.224E-02 0.252E-02 0.140E-02 0.284E-03 0.957E-04  
196 0.130E+03 0.212E-02 0.242E-02 0.127E-02 0.333E-03 0.101E-03  
197 0.135E+03 0.213E-02 0.227E-02 0.114E-02 0.368E-03 0.107E-03  
198 0.140E+03 0.221E-02 0.210E-02 0.101E-02 0.388E-03 0.112E-03  
199 0.150E+03 0.250E-02 0.171E-02 0.777E-03 0.399E-03 0.122E-03  
200 0.160E+03 0.292E-02 0.135E-02 0.586E-03 0.410E-03 0.129E-03  
201 0.170E+03 0.337E-02 0.109E-02 0.461E-03 0.433E-03 0.133E-03  
202 0.180E+03 0.356E-02 0.100E-02 0.418E-03 0.447E-03 0.134E-03  
203  
204 NEUTRON ENERGY = 20.900000  
205 TOTAL CR-SECT.= 1.388248  
206 REACTION CR-SECT. = 0.385556  
207 TOTAL DIRECT CR-SECT. (ELASTIC + DIR.LEV EXCIT.) = 1.002692  
208  
209 NMAX CR-SECT. OF LEVEL EXCITATION  
210 1 0.909358  
211 2 0.061491  
212 3 0.025713  
213 4 0.005039  
214 5 0.001092  
215  
216 STRENGTH FUNCTIONS  
217 SF0= 0.1997542E-04 SF1= 0.2036385E-04 SF2= 0.2953541E-04  
218  
219 SPHERICAL VOLUME INTEGRALS OF REAL POTENTIAL F-FACTORS AND DERIVATIVES:  
220 VIR0= 8.683 VIR1= 19.106 VIR2= 15.798 VIR3= 5.132 CBET0= 1.000  
221 SPHERICAL VOLUME INTEGRALS OF IMAGINARY (WC) F-FACTORS AND DERIVATIVES:  
222 WIC0= 6.132 WIC1= 15.117 WIC2= 13.501 WIC3= 4.471 CBETC= 0.893  
223 SPHERICAL VOLUME INTEGRALS OF IMAGINARY (WD) F-FACTORS AND DERIVATIVES:  
224 WID0= 11.922 WID1= 21.532 WID2= 10.714 WID3= 0.080 CBETD= 0.498  
225 SPHERICAL VOLUME INTEGRALS OF IMAGINARY (WW) F-FACTORS AND DERIVATIVES:  
226 WIW0= 0.000 WIW1= 0.000 WIW2= 0.000 WIW3= 0.000 CBETW= 0.000  
227  
228 POTENTIAL VALUES FOR INCIDENT ENERGY= 39.700000 MeV:  
229  
230 VR= 39.714 VRDC(DWv)= 0.000 DWs= 0.000 Wv= 4.014  
231 Ws= 3.691 VSO= 6.775 DVSO= 0.000 WSO= -0.262  
232 VHF= 0.000 VCUL= 0.831 VCULC= 0.000 RELPOT= 1.021  
233  
234  
235  
236 ORB. MOMENT TRANSITIONS SR SI  
237  
238 0 0.5906007098 -0.4897811171 -0.2914561701  
239 1 0.5900982415 -0.5634083289 0.0091432975  
240 2 0.6110135297 -0.4348204717 0.2344173454  
241 3 0.5819856530 -0.0588710831 0.4037901075  
242 4 0.3566833447 0.5014775551 0.4680959806  
243 5 0.1085472180 0.8851663631 0.2544105027  
244 6 0.0249843609 0.9783258493 0.0940408379  
245 7 0.0056572024 0.9958405169 0.0326444816  
246 8 0.0011868145 0.9992743424 0.0110129421  
247 9 0.0002459203 0.9998661185 0.0036090923  
248 10 0.0000515024 0.9999729892 0.0011223856  
249 11 0.0000106250 0.9999944683 0.0003191821  
250 12 0.0000020932 0.9999989229 0.0000811183  
251 13 0.0000004219 0.9999997863 0.0000183648  
252 14 0.0000000778 0.9999999609 0.0000037438



253	15	0.0000000128	0.9999999936	0.0000006981
254	16	0.0000000019	0.4848484839	0.0000000425
255	17	0.0000000000	0.0000000000	0.0000000000
256	18	0.0000000000	0.0000000000	0.0000000000
257	19	0.0000000000	0.0000000000	0.0000000000
258	20	0.0000000000	0.0000000000	0.0000000000
259	21	0.0000000000	0.0000000000	0.0000000000
260	22	0.0000000000	0.0000000000	0.0000000000
261	23	0.0000000000	0.0000000000	0.0000000000
262	24	0.0000000000	0.0000000000	0.0000000000
263	25	0.0000000000	0.0000000000	0.0000000000

264

265

ANGULAR DISTRIBUTIONS OF SCATTERED PARTICLES

266

267	0.100E+01	0.597E+04	0.851E-02	0.984E-03	0.319E-02	0.482E-03
268	0.500E+01	0.549E+01	0.830E-02	0.102E-02	0.290E-02	0.479E-03
269	0.100E+02	0.781E+00	0.784E-02	0.111E-02	0.214E-02	0.471E-03
270	0.150E+02	0.746E+00	0.762E-02	0.126E-02	0.124E-02	0.462E-03
271	0.180E+02	0.664E+00	0.776E-02	0.138E-02	0.781E-03	0.457E-03
272	0.210E+02	0.554E+00	0.813E-02	0.153E-02	0.439E-03	0.453E-03
273	0.240E+02	0.438E+00	0.867E-02	0.172E-02	0.235E-03	0.448E-03
274	0.270E+02	0.330E+00	0.926E-02	0.195E-02	0.160E-03	0.444E-03
275	0.300E+02	0.238E+00	0.973E-02	0.220E-02	0.188E-03	0.438E-03
276	0.330E+02	0.165E+00	0.997E-02	0.248E-02	0.279E-03	0.429E-03
277	0.360E+02	0.110E+00	0.987E-02	0.274E-02	0.395E-03	0.417E-03
278	0.390E+02	0.709E-01	0.941E-02	0.298E-02	0.502E-03	0.401E-03
279	0.420E+02	0.446E-01	0.862E-02	0.315E-02	0.578E-03	0.381E-03
280	0.450E+02	0.281E-01	0.759E-02	0.325E-02	0.610E-03	0.357E-03
281	0.480E+02	0.185E-01	0.644E-02	0.325E-02	0.599E-03	0.330E-03
282	0.510E+02	0.133E-01	0.529E-02	0.317E-02	0.550E-03	0.300E-03
283	0.540E+02	0.108E-01	0.422E-02	0.300E-02	0.475E-03	0.270E-03
284	0.560E+02	0.999E-02	0.360E-02	0.284E-02	0.418E-03	0.250E-03
285	0.590E+02	0.932E-02	0.280E-02	0.257E-02	0.329E-03	0.221E-03
286	0.620E+02	0.886E-02	0.218E-02	0.226E-02	0.246E-03	0.194E-03
287	0.650E+02	0.833E-02	0.174E-02	0.195E-02	0.176E-03	0.169E-03
288	0.700E+02	0.703E-02	0.129E-02	0.144E-02	0.990E-04	0.135E-03
289	0.750E+02	0.540E-02	0.109E-02	0.102E-02	0.692E-04	0.110E-03
290	0.800E+02	0.377E-02	0.101E-02	0.702E-03	0.731E-04	0.926E-04
291	0.850E+02	0.244E-02	0.940E-03	0.489E-03	0.917E-04	0.818E-04
292	0.900E+02	0.150E-02	0.849E-03	0.359E-03	0.109E-03	0.751E-04
293	0.950E+02	0.943E-03	0.723E-03	0.287E-03	0.114E-03	0.707E-04
294	0.100E+03	0.661E-03	0.579E-03	0.249E-03	0.105E-03	0.672E-04
295	0.105E+03	0.543E-03	0.437E-03	0.225E-03	0.852E-04	0.637E-04
296	0.110E+03	0.498E-03	0.314E-03	0.204E-03	0.609E-04	0.597E-04
297	0.115E+03	0.467E-03	0.221E-03	0.181E-03	0.386E-04	0.553E-04
298	0.120E+03	0.422E-03	0.158E-03	0.155E-03	0.228E-04	0.503E-04
299	0.125E+03	0.359E-03	0.124E-03	0.126E-03	0.156E-04	0.451E-04
300	0.130E+03	0.283E-03	0.112E-03	0.994E-04	0.165E-04	0.399E-04
301	0.135E+03	0.207E-03	0.116E-03	0.759E-04	0.232E-04	0.349E-04
302	0.140E+03	0.141E-03	0.127E-03	0.578E-04	0.332E-04	0.305E-04
303	0.150E+03	0.728E-04	0.151E-03	0.388E-04	0.563E-04	0.241E-04
304	0.160E+03	0.973E-04	0.156E-03	0.360E-04	0.798E-04	0.218E-04
305	0.170E+03	0.166E-03	0.146E-03	0.386E-04	0.100E-03	0.222E-04
306	0.180E+03	0.202E-03	0.139E-03	0.399E-04	0.109E-03	0.228E-04

307

308 PROTON ENERGY = 39.700000

309 TOTAL CR-SECT.= 1.021217

310 REACTION CR-SECT. = 0.272251

```

311 REACTION CR-SECT. incl. LEVELS = 0.317651
312
313 NMAX CR-SECT. OF LEVEL EXCITATION
314 1 0.703566
315 2 0.028443
316 3 0.012605
317 4 0.002521
318 5 0.001832
319
320 STRENGTH FUNCTIONS
321 SF0= 0.1491828E-04 SF1= 0.1559499E-04 SF2= 0.1787253E-04

```

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### CR-SECT file

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```

1 2.09000E+01 1.38825E+00 3.85556E-01 9.09358E-01 6.14911E-02 2.57128E-02 5.03850E-03 1.09154E-03
2 3.97000E+01 1.02122E+00 2.72251E-01 7.03566E-01 2.84426E-02 1.26046E-02 2.52092E-03 1.83165E-03

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### TRANSME file

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```

1 0.209000E+02 21
2 0.57378461 0.53742765 0.63522157 0.55473202 0.12864116 0.01615306
3 0.00239434 0.00027358 0.00001968 0.00000000 0.00000000 0.00000000
4 0.00000000 0.00000000 0.00000000 0.00000000 0.00000000 0.00000000
5 0.00000000 0.00000000 0.00000000 0.00000000
6 0.397000E+02 25
7 0.59060071 0.59009824 0.61101353 0.58198565 0.35668334 0.10854722
8 0.02498436 0.00565720 0.00118681 0.00024592 0.00005150 0.00001062
9 0.00000209 0.00000042 0.00000008 0.00000001 0.00000000 0.00000000
10 0.00000000 0.00000000 0.00000000 0.00000000 0.00000000 0.00000000
11 0.00000000 0.00000000

```

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### GNASH file

---

```

1 0.209000E+02 50
2 0.5738E+00 0.5698E+00 0.7133E+00 0.5212E+00 0.5831E+00 0.4032E+00
3 0.8737E-01 0.6684E+00 0.1617E+00 0.1447E-01 0.2287E-02 0.1755E-01
4 0.2486E-02 0.2974E-03 0.4181E-04 0.2527E-03 0.0000E+00 0.0000E+00
5 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00
6 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00
7 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00
8 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00
9 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00

```

10	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
11	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
12	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
13	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
14	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
15	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
16	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
17	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
18	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
19	0.397000E+02	50				
20	0.5906E+00	0.6115E+00	0.6722E+00	0.5794E+00	0.5702E+00	0.5582E+00
21	0.2766E+00	0.5999E+00	0.4207E+00	0.9128E-01	0.2504E-01	0.1229E+00
22	0.2494E-01	0.6318E-02	0.1476E-02	0.5079E-02	0.9298E-03	0.3456E-03
23	0.8219E-04	0.1562E-03	0.2361E-04	0.1927E-04	0.4287E-05	0.2698E-05
24	0.6780E-07	0.8762E-06	0.1612E-06	-0.3181E-16	0.8511E-16	0.2641E-07
25	0.3847E-08	0.2618E-16	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
26	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
27	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
28	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
29	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
30	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
31	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
32	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
33	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
34	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
35	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00
36	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00	0.0000E+00

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### ANG-DIST file

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1						
2						
3	NEUTRON ENERGY = 20.900000					
4	0.100E+01	0.109E+01	0.165E-01	0.240E-02	0.204E-02	0.126E-03
5	0.500E+01	0.106E+01	0.162E-01	0.241E-02	0.194E-02	0.125E-03
6	0.100E+02	0.967E+00	0.154E-01	0.241E-02	0.164E-02	0.124E-03
7	0.150E+02	0.833E+00	0.142E-01	0.241E-02	0.123E-02	0.121E-03
8	0.180E+02	0.740E+00	0.135E-01	0.242E-02	0.959E-03	0.119E-03
9	0.210E+02	0.643E+00	0.129E-01	0.242E-02	0.703E-03	0.117E-03
10	0.240E+02	0.547E+00	0.123E-01	0.242E-02	0.476E-03	0.114E-03
11	0.270E+02	0.454E+00	0.119E-01	0.243E-02	0.292E-03	0.111E-03
12	0.300E+02	0.368E+00	0.115E-01	0.244E-02	0.159E-03	0.108E-03
13	0.330E+02	0.292E+00	0.112E-01	0.246E-02	0.822E-04	0.105E-03
14	0.360E+02	0.226E+00	0.110E-01	0.248E-02	0.581E-04	0.101E-03

15	0.390E+02	0.171E+00	0.108E-01	0.251E-02	0.813E-04	0.975E-04
16	0.420E+02	0.127E+00	0.105E-01	0.254E-02	0.143E-03	0.939E-04
17	0.450E+02	0.925E-01	0.102E-01	0.259E-02	0.231E-03	0.902E-04
18	0.480E+02	0.666E-01	0.976E-02	0.263E-02	0.335E-03	0.866E-04
19	0.510E+02	0.481E-01	0.925E-02	0.267E-02	0.443E-03	0.831E-04
20	0.540E+02	0.355E-01	0.866E-02	0.271E-02	0.545E-03	0.797E-04
21	0.560E+02	0.297E-01	0.823E-02	0.274E-02	0.605E-03	0.775E-04
22	0.590E+02	0.240E-01	0.754E-02	0.277E-02	0.680E-03	0.745E-04
23	0.620E+02	0.209E-01	0.683E-02	0.279E-02	0.732E-03	0.717E-04
24	0.650E+02	0.194E-01	0.611E-02	0.279E-02	0.758E-03	0.692E-04
25	0.700E+02	0.188E-01	0.500E-02	0.277E-02	0.747E-03	0.658E-04
26	0.750E+02	0.186E-01	0.407E-02	0.271E-02	0.675E-03	0.636E-04
27	0.800E+02	0.178E-01	0.336E-02	0.262E-02	0.563E-03	0.625E-04
28	0.850E+02	0.160E-01	0.289E-02	0.250E-02	0.433E-03	0.626E-04
29	0.900E+02	0.137E-01	0.263E-02	0.236E-02	0.308E-03	0.639E-04
30	0.950E+02	0.110E-01	0.253E-02	0.222E-02	0.207E-03	0.664E-04
31	0.100E+03	0.839E-02	0.252E-02	0.207E-02	0.142E-03	0.698E-04
32	0.105E+03	0.617E-02	0.255E-02	0.193E-02	0.118E-03	0.740E-04
33	0.110E+03	0.447E-02	0.260E-02	0.179E-02	0.131E-03	0.789E-04
34	0.115E+03	0.330E-02	0.261E-02	0.166E-02	0.171E-03	0.843E-04
35	0.120E+03	0.260E-02	0.259E-02	0.153E-02	0.227E-03	0.900E-04
36	0.125E+03	0.224E-02	0.252E-02	0.140E-02	0.284E-03	0.957E-04
37	0.130E+03	0.212E-02	0.242E-02	0.127E-02	0.333E-03	0.101E-03
38	0.135E+03	0.213E-02	0.227E-02	0.114E-02	0.368E-03	0.107E-03
39	0.140E+03	0.221E-02	0.210E-02	0.101E-02	0.388E-03	0.112E-03
40	0.150E+03	0.250E-02	0.171E-02	0.777E-03	0.399E-03	0.122E-03
41	0.160E+03	0.292E-02	0.135E-02	0.586E-03	0.410E-03	0.129E-03
42	0.170E+03	0.337E-02	0.109E-02	0.461E-03	0.433E-03	0.133E-03
43	0.180E+03	0.356E-02	0.100E-02	0.418E-03	0.447E-03	0.134E-03
44						
45						
46						
47	PROTON ENERGY = 39.700000					
48	0.100E+01	0.597E+04	0.851E-02	0.984E-03	0.319E-02	0.482E-03
49	0.500E+01	0.549E+01	0.830E-02	0.102E-02	0.290E-02	0.479E-03
50	0.100E+02	0.781E+00	0.784E-02	0.111E-02	0.214E-02	0.471E-03
51	0.150E+02	0.746E+00	0.762E-02	0.126E-02	0.124E-02	0.462E-03
52	0.180E+02	0.664E+00	0.776E-02	0.138E-02	0.781E-03	0.457E-03
53	0.210E+02	0.554E+00	0.813E-02	0.153E-02	0.439E-03	0.453E-03
54	0.240E+02	0.438E+00	0.867E-02	0.172E-02	0.235E-03	0.448E-03
55	0.270E+02	0.330E+00	0.926E-02	0.195E-02	0.160E-03	0.444E-03
56	0.300E+02	0.238E+00	0.973E-02	0.220E-02	0.188E-03	0.438E-03
57	0.330E+02	0.165E+00	0.997E-02	0.248E-02	0.279E-03	0.429E-03
58	0.360E+02	0.110E+00	0.987E-02	0.274E-02	0.395E-03	0.417E-03
59	0.390E+02	0.709E-01	0.941E-02	0.298E-02	0.502E-03	0.401E-03
60	0.420E+02	0.446E-01	0.862E-02	0.315E-02	0.578E-03	0.381E-03

61	0.450E+02	0.281E-01	0.759E-02	0.325E-02	0.610E-03	0.357E-03
62	0.480E+02	0.185E-01	0.644E-02	0.325E-02	0.599E-03	0.330E-03
63	0.510E+02	0.133E-01	0.529E-02	0.317E-02	0.550E-03	0.300E-03
64	0.540E+02	0.108E-01	0.422E-02	0.300E-02	0.475E-03	0.270E-03
65	0.560E+02	0.999E-02	0.360E-02	0.284E-02	0.418E-03	0.250E-03
66	0.590E+02	0.932E-02	0.280E-02	0.257E-02	0.329E-03	0.221E-03
67	0.620E+02	0.886E-02	0.218E-02	0.226E-02	0.246E-03	0.194E-03
68	0.650E+02	0.833E-02	0.174E-02	0.195E-02	0.176E-03	0.169E-03
69	0.700E+02	0.703E-02	0.129E-02	0.144E-02	0.990E-04	0.135E-03
70	0.750E+02	0.540E-02	0.109E-02	0.102E-02	0.692E-04	0.110E-03
71	0.800E+02	0.377E-02	0.101E-02	0.702E-03	0.731E-04	0.926E-04
72	0.850E+02	0.244E-02	0.940E-03	0.489E-03	0.917E-04	0.818E-04
73	0.900E+02	0.150E-02	0.849E-03	0.359E-03	0.109E-03	0.751E-04
74	0.950E+02	0.943E-03	0.723E-03	0.287E-03	0.114E-03	0.707E-04
75	0.100E+03	0.661E-03	0.579E-03	0.249E-03	0.105E-03	0.672E-04
76	0.105E+03	0.543E-03	0.437E-03	0.225E-03	0.852E-04	0.637E-04
77	0.110E+03	0.498E-03	0.314E-03	0.204E-03	0.609E-04	0.597E-04
78	0.115E+03	0.467E-03	0.221E-03	0.181E-03	0.386E-04	0.553E-04
79	0.120E+03	0.422E-03	0.158E-03	0.155E-03	0.228E-04	0.503E-04
80	0.125E+03	0.359E-03	0.124E-03	0.126E-03	0.156E-04	0.451E-04
81	0.130E+03	0.283E-03	0.112E-03	0.994E-04	0.165E-04	0.399E-04
82	0.135E+03	0.207E-03	0.116E-03	0.759E-04	0.232E-04	0.349E-04
83	0.140E+03	0.141E-03	0.127E-03	0.578E-04	0.332E-04	0.305E-04
84	0.150E+03	0.728E-04	0.151E-03	0.388E-04	0.563E-04	0.241E-04
85	0.160E+03	0.973E-04	0.156E-03	0.360E-04	0.798E-04	0.218E-04
86	0.170E+03	0.166E-03	0.146E-03	0.386E-04	0.100E-03	0.222E-04
87	0.180E+03	0.202E-03	0.139E-03	0.399E-04	0.109E-03	0.228E-04

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## ANG-POL file

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1  EN= 20.900 LKK= 16 LCOUL= 0
2  2.09000E+01 1.38825E+00 3.85556E-01 9.09358E-01 6.14911E-02 2.57128E-02 5.03850E-03 1.09154E-03
3
4  LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
5  ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
6  0.7236444E-01 0.5856492E-01 0.4359487E-01 0.3064064E-01 0.2004334E-01 0.1102915E-01
7  0.4910092E-02 0.1812372E-02 0.6086285E-03 0.1845791E-03 0.5168079E-04 0.1260924E-04
8  0.2571940E-05 0.4288946E-06 0.6065312E-07 0.5987911E-08
9
10 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
11 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
12 0.4893306E-02 0.1859738E-02 0.7033964E-03 0.1801537E-03 -0.1117703E-03 -0.1458413E-04
13 0.6033732E-04 0.5668852E-04 0.2931499E-04 0.1127401E-04 0.3805886E-05 0.1069167E-05
14 0.2616600E-06 0.4785041E-07 0.7018580E-08 0.6662893E-09
15
16 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
17 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE

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18 0.2046161E-02 0.3529176E-03 -0.1346601E-03 -0.2652838E-04 -0.1485108E-08 0.1308584E-04
19 0.4757761E-05 -0.8550424E-06 -0.1120968E-05 -0.4991774E-06 -0.1993302E-06 -0.5435987E-07
20 -0.1188290E-07 -0.2006546E-08 -0.2632848E-09 -0.2161875E-10
21
22 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
23 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
24 0.4009511E-03 0.5381315E-04 0.2113663E-04 -0.2369844E-04 0.2347598E-04 0.5741010E-04
25 0.3419885E-04 0.9867403E-05 0.4557760E-05 0.1267224E-05 0.3536885E-06 0.8368380E-07
26 0.1724356E-07 0.2461749E-08 0.2916871E-09 0.2406994E-10
27
28 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
29 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
30 0.8686176E-04 -0.3504770E-05 0.8932224E-05 0.9562821E-06 -0.1666646E-06 -0.5083541E-07
31 0.4903044E-08 0.6371851E-09 0.1080125E-09 0.1020842E-10 0.9459343E-12 0.7785877E-13
32 0.4900660E-14 0.2347191E-15 0.1004641E-16 0.3073881E-18
33
34 EN= 39.700 LKK= 32 LCOUL= 17
35 3.97000E+01 1.02122E+00 2.72251E-01 7.03566E-01 2.84426E-02 1.26046E-02 2.52092E-03 1.83165E-03
36
37 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
38 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
39 0.5598798E-01 0.5074457E-01 0.4284439E-01 0.3433034E-01 0.2611217E-01 0.1844517E-01
40 0.1185086E-01 0.6848544E-02 0.3594070E-02 0.1744406E-02 0.7982288E-03 0.348327E-03
41 0.1458440E-03 0.5866221E-04 0.2260938E-04 0.8320790E-05 0.2919245E-05 0.9773950E-06
42 0.3130308E-06 0.9608660E-07 0.2825816E-07 0.7945747E-08 0.2133602E-08 0.5455996E-09
43 0.1320229E-09 0.2997540E-10 0.6316940E-11 0.1219742E-11 0.2124699E-12 0.3259838E-13
44 0.4178419E-14 0.3762704E-15
45
46 LEGENDRE COEFFICIENTS FOR SCATTERED PROTONS
47 ANGULAR DISTRIBUTIONS - COULOMB AMPLITUDE
48 -0.5329607E-05 0.2724234E-04 -0.8424545E-05 0.2731975E-04 -0.7764765E-05 0.2542598E-04
49 -0.4043458E-05 0.2032444E-04 0.1452678E-05 0.1242014E-04 0.2210252E-05 0.4600765E-05
50 0.9878467E-06 0.1462332E-05 0.3641806E-06 0.4790582E-06 0.1257265E-06 0.1578749E-06
51 0.4086036E-07 0.5188394E-07 0.1241434E-07 0.1635135E-07 0.3430490E-08 0.4723129E-08
52 0.8457065E-09 0.1218796E-08 0.1853447E-09 0.2800696E-09 0.3659081E-10 0.5786558E-10
53 0.6614075E-11 0.1091870E-10 0.3810165E-12 0.6769627E-12
54
55 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
56 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
57 0.2263392E-02 0.1460025E-02 0.6737818E-03 0.1618302E-03 -0.8769747E-04 -0.1819126E-03
58 -0.1410000E-03 -0.4008538E-04 0.2540215E-04 0.4211670E-04 0.3300129E-04 0.1901077E-04
59 0.8405557E-05 0.2446360E-05 0.1313958E-07 -0.5207376E-06 -0.3885188E-06 -0.1830124E-06
60 -0.6194802E-07 -0.1350774E-07 -0.2310001E-10 0.1757144E-08 0.1083614E-08 0.4511116E-09
61 0.1504731E-09 0.4225555E-10 0.1016230E-10 0.2100226E-11 0.3707529E-12 0.5486393E-13
62 0.6476804E-14 0.5108695E-15
63
64 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
65 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
66 0.1003039E-02 0.5300818E-03 0.8737568E-04 -0.1194362E-03 -0.1343310E-03 -0.5849834E-04
67 0.8217083E-05 0.2319584E-04 0.1335915E-04 0.3594827E-05 -0.7882556E-06 -0.1578205E-05
68 -0.1149896E-05 -0.6060559E-06 -0.2559891E-06 -0.8898399E-07 -0.2556876E-07 -0.5977355E-08
69 -0.1062793E-08 -0.1048700E-09 0.1461403E-10 0.1094271E-10 0.3410178E-11 0.7540120E-12
70 0.1231391E-12 0.1279214E-13 -0.1053389E-15 -0.4179632E-15 -0.1150545E-15 -0.2050942E-16
71 -0.2642750E-17 -0.2141415E-18
72
73 LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
74 ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
75 0.2006087E-03 0.1086488E-03 0.5642774E-04 0.2164884E-04 0.1608478E-04 0.1187074E-04

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76 0.2341582E-04 0.3228098E-04 0.2758898E-04 0.1711844E-04 0.9261094E-05 0.4449792E-05
77 0.1969907E-05 0.8147737E-06 0.3161590E-06 0.1148564E-06 0.3897304E-07 0.1235672E-07
78 0.3672258E-08 0.1027123E-08 0.2709612E-09 0.6742394E-10 0.1580305E-10 0.3483977E-11
79 0.7183827E-12 0.1373959E-12 0.2410329E-13 0.3821297E-14 0.5370573E-15 0.6500135E-16
80 0.6395858E-17 0.4295177E-18
81
82          LEGENDRE COEFFICIENTS FOR SCATTERED NUCLEONS
83          ANGULAR DISTRIBUTIONS - NUCLEAR AMPLITUDE
84 0.1457584E-03 0.7087144E-04 0.2677219E-04 0.6858640E-05 -0.3044164E-05 -0.4037173E-05
85 -0.1106796E-05 0.3141728E-06 0.6018783E-06 0.3531678E-06 0.1648211E-06 0.6759515E-07
86 0.2670557E-07 0.1020991E-07 0.3666891E-08 0.1211379E-08 0.3684126E-09 0.1045432E-09
87 0.2806741E-10 0.7169551E-11 0.1737758E-11 0.3969828E-12 0.8483983E-13 0.1682128E-13
88 0.3065229E-14 0.5084441E-15 0.7599398E-16 0.1011545E-16 0.1180852E-17 0.1179210E-18
89 0.9545002E-20 0.5223029E-21

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### OPTMAN.CS file

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```

1 <CROSS-S.>      1.01  20.90000      12.00   0   3
2 0.13882485E+04
3 0.10026924E+04
4 0.38555610E+03
5 <CROSS-S.>      1.01  39.70000      12.00   0   1
6 0.74896559E+03

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### OPTMAN.ANG file

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```

1 <ANG.DIS.>      1.01  20.90000      12.00   0   5
2 1 0.0+ 0 40
3 1.0 0.10873E+04
4 5.0 0.10568E+04
5 10.0 0.96686E+03
6 15.0 0.83318E+03
7 18.0 0.74017E+03
8 21.0 0.64317E+03
9 24.0 0.54650E+03
10 27.0 0.45392E+03
11 30.0 0.36842E+03
12 33.0 0.29213E+03
13 36.0 0.22629E+03
14 39.0 0.17132E+03
15 42.0 0.12698E+03
16 45.0 0.92463E+02
17 48.0 0.66620E+02
18 51.0 0.48101E+02
19 54.0 0.35497E+02
20 56.0 0.29703E+02
21 59.0 0.23991E+02

```

22	62.0	0.20870E+02
23	65.0	0.19412E+02
24	70.0	0.18784E+02
25	75.0	0.18587E+02
26	80.0	0.17758E+02
27	85.0	0.16044E+02
28	90.0	0.13652E+02
29	95.0	0.10972E+02
30	100.0	0.83859E+01
31	105.0	0.61725E+01
32	110.0	0.44727E+01
33	115.0	0.33040E+01
34	120.0	0.25978E+01
35	125.0	0.22415E+01
36	130.0	0.21171E+01
37	135.0	0.21279E+01
38	140.0	0.22110E+01
39	150.0	0.24982E+01
40	160.0	0.29218E+01
41	170.0	0.33657E+01
42	180.0	0.35638E+01
43	2	2.0+ 0 40
44	1.0	0.16531E+02
45	5.0	0.16226E+02
46	10.0	0.15369E+02
47	15.0	0.14221E+02
48	18.0	0.13515E+02
49	21.0	0.12865E+02
50	24.0	0.12305E+02
51	27.0	0.11852E+02
52	30.0	0.11499E+02
53	33.0	0.11224E+02
54	36.0	0.10992E+02
55	39.0	0.10763E+02
56	42.0	0.10499E+02
57	45.0	0.10170E+02
58	48.0	0.97564E+01
59	51.0	0.92524E+01
60	54.0	0.86637E+01
61	56.0	0.82320E+01
62	59.0	0.75425E+01
63	62.0	0.68268E+01
64	65.0	0.61133E+01
65	70.0	0.50006E+01
66	75.0	0.40652E+01
67	80.0	0.33600E+01



68	85.0	0.28917E+01
69	90.0	0.26305E+01
70	95.0	0.25251E+01
71	100.0	0.25175E+01
72	105.0	0.25547E+01
73	110.0	0.25955E+01
74	115.0	0.26122E+01
75	120.0	0.25902E+01
76	125.0	0.25243E+01
77	130.0	0.24164E+01
78	135.0	0.22723E+01
79	140.0	0.21004E+01
80	150.0	0.17148E+01
81	160.0	0.13509E+01
82	170.0	0.10949E+01
83	180.0	0.10035E+01
84	3	3.0- 0 40
85	1.0	0.24046E+01
86	5.0	0.24061E+01
87	10.0	0.24098E+01
88	15.0	0.24138E+01
89	18.0	0.24161E+01
90	21.0	0.24187E+01
91	24.0	0.24227E+01
92	27.0	0.24292E+01
93	30.0	0.24399E+01
94	33.0	0.24561E+01
95	36.0	0.24789E+01
96	39.0	0.25086E+01
97	42.0	0.25446E+01
98	45.0	0.25858E+01
99	48.0	0.26298E+01
100	51.0	0.26739E+01
101	54.0	0.27150E+01
102	56.0	0.27391E+01
103	59.0	0.27682E+01
104	62.0	0.27865E+01
105	65.0	0.27920E+01
106	70.0	0.27702E+01
107	75.0	0.27096E+01
108	80.0	0.26152E+01
109	85.0	0.24952E+01
110	90.0	0.23592E+01
111	95.0	0.22159E+01
112	100.0	0.20718E+01
113	105.0	0.19307E+01

114	110.0	0.17938E+01
115	115.0	0.16605E+01
116	120.0	0.15293E+01
117	125.0	0.13989E+01
118	130.0	0.12687E+01
119	135.0	0.11394E+01
120	140.0	0.10126E+01
121	150.0	0.77731E+00
122	160.0	0.58629E+00
123	170.0	0.46131E+00
124	180.0	0.41777E+00
125	4	0.0+ 0 40
126	1.0	0.20446E+01
127	5.0	0.19409E+01
128	10.0	0.16431E+01
129	15.0	0.12267E+01
130	18.0	0.95881E+00
131	21.0	0.70252E+00
132	24.0	0.47560E+00
133	27.0	0.29172E+00
134	30.0	0.15948E+00
135	33.0	0.82195E-01
136	36.0	0.58132E-01
137	39.0	0.81325E-01
138	42.0	0.14266E+00
139	45.0	0.23113E+00
140	48.0	0.33500E+00
141	51.0	0.44296E+00
142	54.0	0.54492E+00
143	56.0	0.60541E+00
144	59.0	0.68016E+00
145	62.0	0.73189E+00
146	65.0	0.75839E+00
147	70.0	0.74689E+00
148	75.0	0.67542E+00
149	80.0	0.56299E+00
150	85.0	0.43283E+00
151	90.0	0.30786E+00
152	95.0	0.20691E+00
153	100.0	0.14214E+00
154	105.0	0.11784E+00
155	110.0	0.13065E+00
156	115.0	0.17117E+00
157	120.0	0.22663E+00
158	125.0	0.28401E+00
159	130.0	0.33296E+00

160	135.0	0.36769E+00
161	140.0	0.38759E+00
162	150.0	0.39950E+00
163	160.0	0.41004E+00
164	170.0	0.43332E+00
165	180.0	0.44671E+00
166	5	2.0+ 0 40
167	1.0	0.12570E+00
168	5.0	0.12518E+00
169	10.0	0.12358E+00
170	15.0	0.12096E+00
171	18.0	0.11893E+00
172	21.0	0.11658E+00
173	24.0	0.11394E+00
174	27.0	0.11104E+00
175	30.0	0.10790E+00
176	33.0	0.10458E+00
177	36.0	0.10111E+00
178	39.0	0.97522E-01
179	42.0	0.93878E-01
180	45.0	0.90221E-01
181	48.0	0.86601E-01
182	51.0	0.83068E-01
183	54.0	0.79675E-01
184	56.0	0.77514E-01
185	59.0	0.74460E-01
186	62.0	0.71672E-01
187	65.0	0.69192E-01
188	70.0	0.65841E-01
189	75.0	0.63582E-01
190	80.0	0.62502E-01
191	85.0	0.62633E-01
192	90.0	0.63949E-01
193	95.0	0.66374E-01
194	100.0	0.69786E-01
195	105.0	0.74030E-01
196	110.0	0.78932E-01
197	115.0	0.84305E-01
198	120.0	0.89967E-01
199	125.0	0.95744E-01
200	130.0	0.10148E+00
201	135.0	0.10703E+00
202	140.0	0.11229E+00
203	150.0	0.12153E+00
204	160.0	0.12856E+00
205	170.0	0.13296E+00

206	180.0	0.13446E+00				
207	<ANG.DIS.>	1.01	39.70000	12.00	0	5
208	1	0.0+	0	40		
209	1.0	0.59705E+07				
210	5.0	0.54908E+04				
211	10.0	0.78120E+03				
212	15.0	0.74650E+03				
213	18.0	0.66405E+03				
214	21.0	0.55361E+03				
215	24.0	0.43761E+03				
216	27.0	0.33008E+03				
217	30.0	0.23837E+03				
218	33.0	0.16516E+03				
219	36.0	0.11005E+03				
220	39.0	0.70878E+02				
221	42.0	0.44614E+02				
222	45.0	0.28105E+02				
223	48.0	0.18467E+02				
224	51.0	0.13313E+02				
225	54.0	0.10823E+02				
226	56.0	0.99897E+01				
227	59.0	0.93205E+01				
228	62.0	0.88627E+01				
229	65.0	0.83253E+01				
230	70.0	0.70323E+01				
231	75.0	0.53966E+01				
232	80.0	0.37728E+01				
233	85.0	0.24379E+01				
234	90.0	0.15032E+01				
235	95.0	0.94338E+00				
236	100.0	0.66079E+00				
237	105.0	0.54305E+00				
238	110.0	0.49806E+00				
239	115.0	0.46703E+00				
240	120.0	0.42236E+00				
241	125.0	0.35869E+00				
242	130.0	0.28291E+00				
243	135.0	0.20655E+00				
244	140.0	0.14094E+00				
245	150.0	0.72779E-01				
246	160.0	0.97280E-01				
247	170.0	0.16606E+00				
248	180.0	0.20160E+00				
249	2	2.0+	0	40		
250	1.0	0.85118E+01				
251	5.0	0.82978E+01				

252	10.0	0.78382E+01
253	15.0	0.76158E+01
254	18.0	0.77609E+01
255	21.0	0.81348E+01
256	24.0	0.86737E+01
257	27.0	0.92565E+01
258	30.0	0.97348E+01
259	33.0	0.99707E+01
260	36.0	0.98705E+01
261	39.0	0.94060E+01
262	42.0	0.86154E+01
263	45.0	0.75887E+01
264	48.0	0.64415E+01
265	51.0	0.52889E+01
266	54.0	0.42246E+01
267	56.0	0.35959E+01
268	59.0	0.27992E+01
269	62.0	0.21840E+01
270	65.0	0.17374E+01
271	70.0	0.12931E+01
272	75.0	0.10940E+01
273	80.0	0.10070E+01
274	85.0	0.94036E+00
275	90.0	0.84853E+00
276	95.0	0.72345E+00
277	100.0	0.57928E+00
278	105.0	0.43707E+00
279	110.0	0.31431E+00
280	115.0	0.22067E+00
281	120.0	0.15830E+00
282	125.0	0.12415E+00
283	130.0	0.11227E+00
284	135.0	0.11555E+00
285	140.0	0.12691E+00
286	150.0	0.15087E+00
287	160.0	0.15631E+00
288	170.0	0.14610E+00
289	180.0	0.13924E+00
290	3	3.0- 0 40
291	1.0	0.98393E+00
292	5.0	0.10159E+01
293	10.0	0.11100E+01
294	15.0	0.12598E+01
295	18.0	0.13812E+01
296	21.0	0.15333E+01
297	24.0	0.17217E+01

298	27.0	0.19473E+01
299	30.0	0.22037E+01
300	33.0	0.24760E+01
301	36.0	0.27419E+01
302	39.0	0.29751E+01
303	42.0	0.31500E+01
304	45.0	0.32461E+01
305	48.0	0.32518E+01
306	51.0	0.31654E+01
307	54.0	0.29951E+01
308	56.0	0.28425E+01
309	59.0	0.25699E+01
310	62.0	0.22644E+01
311	65.0	0.19476E+01
312	70.0	0.14437E+01
313	75.0	0.10199E+01
314	80.0	0.70236E+00
315	85.0	0.48860E+00
316	90.0	0.35894E+00
317	95.0	0.28723E+00
318	100.0	0.24889E+00
319	105.0	0.22519E+00
320	110.0	0.20445E+00
321	115.0	0.18125E+00
322	120.0	0.15468E+00
323	125.0	0.12648E+00
324	130.0	0.99373E-01
325	135.0	0.75921E-01
326	140.0	0.57815E-01
327	150.0	0.38754E-01
328	160.0	0.35982E-01
329	170.0	0.38557E-01
330	180.0	0.39943E-01
331	4	0.0+ 0 40
332	1.0	0.31887E+01
333	5.0	0.29017E+01
334	10.0	0.21400E+01
335	15.0	0.12427E+01
336	18.0	0.78140E+00
337	21.0	0.43925E+00
338	24.0	0.23481E+00
339	27.0	0.16019E+00
340	30.0	0.18773E+00
341	33.0	0.27923E+00
342	36.0	0.39525E+00
343	39.0	0.50245E+00

344	42.0	0.57788E+00
345	45.0	0.61022E+00
346	48.0	0.59858E+00
347	51.0	0.54985E+00
348	54.0	0.47546E+00
349	56.0	0.41799E+00
350	59.0	0.32911E+00
351	62.0	0.24610E+00
352	65.0	0.17635E+00
353	70.0	0.99013E-01
354	75.0	0.69250E-01
355	80.0	0.73137E-01
356	85.0	0.91696E-01
357	90.0	0.10850E+00
358	95.0	0.11373E+00
359	100.0	0.10489E+00
360	105.0	0.85209E-01
361	110.0	0.60931E-01
362	115.0	0.38552E-01
363	120.0	0.22783E-01
364	125.0	0.15610E-01
365	130.0	0.16468E-01
366	135.0	0.23191E-01
367	140.0	0.33230E-01
368	150.0	0.56308E-01
369	160.0	0.79751E-01
370	170.0	0.10046E+00
371	180.0	0.10935E+00
372	5	2.0+ 0 40
373	1.0	0.48170E+00
374	5.0	0.47873E+00
375	10.0	0.47094E+00
376	15.0	0.46183E+00
377	18.0	0.45690E+00
378	21.0	0.45257E+00
379	24.0	0.44846E+00
380	27.0	0.44382E+00
381	30.0	0.43765E+00
382	33.0	0.42893E+00
383	36.0	0.41680E+00
384	39.0	0.40071E+00
385	42.0	0.38054E+00
386	45.0	0.35660E+00
387	48.0	0.32957E+00
388	51.0	0.30043E+00
389	54.0	0.27033E+00

390	56.0	0.25029E+00
391	59.0	0.22110E+00
392	62.0	0.19378E+00
393	65.0	0.16906E+00
394	70.0	0.13481E+00
395	75.0	0.10962E+00
396	80.0	0.92565E-01
397	85.0	0.81773E-01
398	90.0	0.75127E-01
399	95.0	0.70740E-01
400	100.0	0.67214E-01
401	105.0	0.63685E-01
402	110.0	0.59735E-01
403	115.0	0.55258E-01
404	120.0	0.50334E-01
405	125.0	0.45144E-01
406	130.0	0.39923E-01
407	135.0	0.34941E-01
408	140.0	0.30481E-01
409	150.0	0.24109E-01
410	160.0	0.21772E-01
411	170.0	0.22206E-01
412	180.0	0.22803E-01

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OPTMAN.ICS file

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1	<INE.C.S.>	1.01	20.90000	12.00	0	4
2	0.61491105E+02					
3	0.25712818E+02					
4	0.50385012E+01					
5	0.10915372E+01					
6	<INE.C.S.>	1.01	39.70000	12.00	0	4
7	0.28442630E+02					
8	0.12604566E+02					
9	0.25209230E+01					
10	0.18316537E+01					

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OPTMAN.TL file

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1	<TLJ	>	1.01	20.90000	12.00	0	16
2	0.5	+	1				
3	1	0	0.5	5.7378461D-01			



4	1.5	+	1						
5	1	2	1.5	7.1334591D-01					
6	2.5	+	1						
7	1	2	2.5	5.8313868D-01					
8	3.5	+	1						
9	1	4	3.5	8.7373406D-02					
10	4.5	+	1						
11	1	4	4.5	1.6165536D-01					
12	5.5	+	1						
13	1	6	5.5	2.2870480D-03					
14	6.5	+	1						
15	1	6	6.5	2.4863075D-03					
16	7.5	+	1						
17	1	8	7.5	4.1809846D-05					
18	0.5	-	1						
19	1	1	0.5	5.6979255D-01					
20	1.5	-	1						
21	1	1	1.5	5.2124520D-01					
22	2.5	-	1						
23	1	3	2.5	4.0323268D-01					
24	3.5	-	1						
25	1	3	3.5	6.6835652D-01					
26	4.5	-	1						
27	1	5	4.5	1.4472244D-02					
28	5.5	-	1						
29	1	5	5.5	1.7553733D-02					
30	6.5	-	1						
31	1	7	6.5	2.9739545D-04					
32	7.5	-	1						
33	1	7	7.5	2.5273519D-04					
34	<TLJ	>	1.01	39.70000	12.00	0	31		
35	0.5	+	1						
36	1	0	0.5	5.9060071D-01					
37	1.5	+	1						
38	1	2	1.5	6.7224853D-01					
39	2.5	+	1						
40	1	2	2.5	5.7019019D-01					
41	3.5	+	1						
42	1	4	3.5	2.7661760D-01					
43	4.5	+	1						
44	1	4	4.5	4.2073594D-01					
45	5.5	+	1						
46	1	6	5.5	2.5037786D-02					
47	6.5	+	1						
48	1	6	6.5	2.4938568D-02					
49	7.5	+	1						

50	1	8	7.5	1.4760093D-03
51	8.5	+	1	
52	1	8	8.5	9.2975241D-04
53	9.5	+	1	
54	1	10	9.5	8.2186264D-05
55	10.5	+	1	
56	1	10	10.5	2.3607990D-05
57	11.5	+	1	
58	1	12	11.5	4.2874804D-06
59	12.5	+	1	
60	1	12	12.5	6.7800821D-08
61	13.5	+	1	
62	1	14	13.5	1.6115521D-07
63	14.5	+	1	
64	1	14	14.5	8.5105837D-17
65	15.5	+	1	
66	1	16	15.5	3.8466296D-09
67	0.5	-	1	
68	1	1	0.5	6.1145569D-01
69	1.5	-	1	
70	1	1	1.5	5.7941952D-01
71	2.5	-	1	
72	1	3	2.5	5.5815178D-01
73	3.5	-	1	
74	1	3	3.5	5.9986105D-01
75	4.5	-	1	
76	1	5	4.5	9.1284953D-02
77	5.5	-	1	
78	1	5	5.5	1.2293244D-01
79	6.5	-	1	
80	1	7	6.5	6.3177864D-03
81	7.5	-	1	
82	1	7	7.5	5.0791914D-03
83	8.5	-	1	
84	1	9	8.5	3.4558767D-04
85	9.5	-	1	
86	1	9	9.5	1.5621959D-04
87	10.5	-	1	
88	1	11	10.5	1.9272893D-05
89	11.5	-	1	
90	1	11	11.5	2.6977274D-06
91	12.5	-	1	
92	1	13	12.5	8.7616063D-07
93	14.5	-	1	
94	1	15	14.5	2.6414880D-08
95	15.5	-	1	

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OPTMAN.LEG

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1	<LEGENBRE>		1.01	20.90000	12.00	0	5
2	1	16	COUPLED LEVEL, NUMBER OF VALUES				
3	1	0	7.2364437268D-02				
4	1	1	5.8564916552D-02				
5	1	2	4.3594871585D-02				
6	1	3	3.0640641116D-02				
7	1	4	2.0043336199D-02				
8	1	5	1.1029154493D-02				
9	1	6	4.9100920395D-03				
10	1	7	1.8123723722D-03				
11	1	8	6.0862853319D-04				
12	1	9	1.8457905391D-04				
13	1	10	5.1680789259D-05				
14	1	11	1.2609243978D-05				
15	1	12	2.5719399628D-06				
16	1	13	4.2889464056D-07				
17	1	14	6.0653118851D-08				
18	1	15	5.9879111272D-09				
19	2	16	COUPLED LEVEL, NUMBER OF VALUES				
20	2	0	4.8933061531D-03				
21	2	1	1.8597380079D-03				
22	2	2	7.0339639725D-04				
23	2	3	1.8015373390D-04				
24	2	4	-1.1177027850D-04				
25	2	5	-1.4584126582D-05				
26	2	6	6.0337318908D-05				
27	2	7	5.6688520974D-05				
28	2	8	2.9314991472D-05				
29	2	9	1.1274013955D-05				
30	2	10	3.8058863883D-06				
31	2	11	1.0691670409D-06				
32	2	12	2.6166002281D-07				
33	2	13	4.7850413911D-08				
34	2	14	7.0185797767D-09				
35	2	15	6.6628931248D-10				
36	3	16	COUPLED LEVEL, NUMBER OF VALUES				
37	3	0	2.0461608574D-03				
38	3	1	3.5291760177D-04				
39	3	2	-1.3466006335D-04				
40	3	3	-2.6528381269D-05				

41	3	4	-1.4851076140D-09
42	3	5	1.3085843653D-05
43	3	6	4.7577614262D-06
44	3	7	-8.5504239255D-07
45	3	8	-1.1209677508D-06
46	3	9	-4.9917740300D-07
47	3	10	-1.9933022862D-07
48	3	11	-5.4359873705D-08
49	3	12	-1.1882900960D-08
50	3	13	-2.0065458118D-09
51	3	14	-2.6328478998D-10
52	3	15	-2.1618754055D-11
53	4	16	COUPLED LEVEL, NUMBER OF VALUES
54	4	0	4.0095114597D-04
55	4	1	5.3813153943D-05
56	4	2	2.1136626613D-05
57	4	3	-2.3698437268D-05
58	4	4	2.3475977527D-05
59	4	5	5.7410102724D-05
60	4	6	3.4198853087D-05
61	4	7	9.8674028411D-06
62	4	8	4.5577598936D-06
63	4	9	1.2672243782D-06
64	4	10	3.5368848489D-07
65	4	11	8.3683796220D-08
66	4	12	1.7243557721D-08
67	4	13	2.4617489546D-09
68	4	14	2.9168709366D-10
69	4	15	2.4069942819D-11
70	5	16	COUPLED LEVEL, NUMBER OF VALUES
71	5	0	8.6861758548D-05
72	5	1	-3.5047701379D-06
73	5	2	8.9322244359D-06
74	5	3	9.5628207382D-07
75	5	4	-1.6666460404D-07
76	5	5	-5.0835412517D-08
77	5	6	4.9030444034D-09
78	5	7	6.3718509893D-10
79	5	8	1.0801246490D-10
80	5	9	1.0208419009D-11
81	5	10	9.4593426970D-13
82	5	11	7.7858767077D-14
83	5	12	4.9006596556D-15
84	5	13	2.3471913763D-16
85	5	14	1.0046406941D-17
86	5	15	3.0738809080D-19

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ANGL-DIS.YW

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1	20.900000	40	5						
2	0.100000E+01	0.108726E+01	0.165310E-01	0.240460E-02	0.204458E-02	0.125697E-03			
3	0.500000E+01	0.105681E+01	0.162261E-01	0.240608E-02	0.194089E-02	0.125180E-03			
4	0.100000E+02	0.966861E+00	0.153692E-01	0.240978E-02	0.164307E-02	0.123577E-03			
5	0.150000E+02	0.833177E+00	0.142210E-01	0.241381E-02	0.122667E-02	0.120957E-03			
6	0.180000E+02	0.740173E+00	0.135154E-01	0.241609E-02	0.958808E-03	0.118927E-03			
7	0.210000E+02	0.643174E+00	0.128647E-01	0.241874E-02	0.702524E-03	0.116579E-03			
8	0.240000E+02	0.546501E+00	0.123050E-01	0.242268E-02	0.475604E-03	0.113939E-03			
9	0.270000E+02	0.453916E+00	0.118516E-01	0.242923E-02	0.291715E-03	0.111036E-03			
10	0.300000E+02	0.368416E+00	0.114990E-01	0.243992E-02	0.159481E-03	0.107905E-03			
11	0.330000E+02	0.292125E+00	0.112240E-01	0.245614E-02	0.821951E-04	0.104581E-03			
12	0.360000E+02	0.226286E+00	0.109918E-01	0.247890E-02	0.581318E-04	0.101105E-03			
13	0.390000E+02	0.171325E+00	0.107627E-01	0.250856E-02	0.813251E-04	0.975224E-04			
14	0.420000E+02	0.126981E+00	0.104988E-01	0.254464E-02	0.142664E-03	0.938780E-04			
15	0.450000E+02	0.924628E-01	0.101698E-01	0.258578E-02	0.231126E-03	0.902207E-04			
16	0.480000E+02	0.666197E-01	0.975644E-02	0.262980E-02	0.334999E-03	0.866006E-04			
17	0.510000E+02	0.481013E-01	0.925239E-02	0.267392E-02	0.442956E-03	0.830684E-04			
18	0.540000E+02	0.354969E-01	0.866369E-02	0.271499E-02	0.544917E-03	0.796747E-04			
19	0.560000E+02	0.297028E-01	0.823196E-02	0.273911E-02	0.605408E-03	0.775138E-04			
20	0.590000E+02	0.239909E-01	0.754245E-02	0.276818E-02	0.680158E-03	0.744602E-04			
21	0.620000E+02	0.208705E-01	0.682676E-02	0.278646E-02	0.731895E-03	0.716723E-04			
22	0.650000E+02	0.194125E-01	0.611327E-02	0.279200E-02	0.758385E-03	0.691921E-04			
23	0.700000E+02	0.187836E-01	0.500058E-02	0.277016E-02	0.746889E-03	0.658406E-04			
24	0.750000E+02	0.185870E-01	0.406518E-02	0.270960E-02	0.675424E-03	0.635819E-04			
25	0.800000E+02	0.177584E-01	0.336000E-02	0.261518E-02	0.562990E-03	0.625025E-04			
26	0.850000E+02	0.160444E-01	0.289170E-02	0.249522E-02	0.432834E-03	0.626332E-04			
27	0.900000E+02	0.136525E-01	0.263050E-02	0.235923E-02	0.307862E-03	0.639493E-04			
28	0.950000E+02	0.109719E-01	0.252508E-02	0.221593E-02	0.206906E-03	0.663737E-04			
29	0.100000E+03	0.838593E-02	0.251749E-02	0.207183E-02	0.142135E-03	0.697856E-04			
30	0.105000E+03	0.617254E-02	0.255474E-02	0.193073E-02	0.117836E-03	0.740303E-04			
31	0.110000E+03	0.447266E-02	0.259549E-02	0.179382E-02	0.130645E-03	0.789319E-04			
32	0.115000E+03	0.330403E-02	0.261224E-02	0.166050E-02	0.171167E-03	0.843053E-04			
33	0.120000E+03	0.259780E-02	0.259017E-02	0.152930E-02	0.226627E-03	0.899671E-04			
34	0.125000E+03	0.224149E-02	0.252431E-02	0.139887E-02	0.284012E-03	0.957440E-04			
35	0.130000E+03	0.211711E-02	0.241637E-02	0.126868E-02	0.332957E-03	0.101479E-03			
36	0.135000E+03	0.212789E-02	0.227225E-02	0.113936E-02	0.367691E-03	0.107034E-03			
37	0.140000E+03	0.221098E-02	0.210037E-02	0.101260E-02	0.387588E-03	0.112291E-03			
38	0.150000E+03	0.249824E-02	0.171485E-02	0.777314E-03	0.399496E-03	0.121526E-03			
39	0.160000E+03	0.292177E-02	0.135090E-02	0.586289E-03	0.410036E-03	0.128560E-03			
40	0.170000E+03	0.336567E-02	0.109490E-02	0.461315E-03	0.433318E-03	0.132961E-03			
41	0.180000E+03	0.356375E-02	0.100346E-02	0.417768E-03	0.446710E-03	0.134458E-03			

42 39.700000 40 5  
43 0.100000E+01 0.597054E+04 0.851176E-02 0.983933E-03 0.318872E-02 0.481695E-03  
44 0.500000E+01 0.549081E+01 0.829776E-02 0.101589E-02 0.290170E-02 0.478727E-03  
45 0.100000E+02 0.781204E+00 0.783821E-02 0.111004E-02 0.214000E-02 0.470945E-03  
46 0.150000E+02 0.746499E+00 0.761579E-02 0.125976E-02 0.124272E-02 0.461832E-03  
47 0.180000E+02 0.664046E+00 0.776093E-02 0.138116E-02 0.781403E-03 0.456904E-03  
48 0.210000E+02 0.553607E+00 0.813484E-02 0.153335E-02 0.439254E-03 0.452568E-03  
49 0.240000E+02 0.437610E+00 0.867365E-02 0.172171E-02 0.234810E-03 0.448458E-03  
50 0.270000E+02 0.330076E+00 0.925646E-02 0.194727E-02 0.160195E-03 0.443816E-03  
51 0.300000E+02 0.238366E+00 0.973484E-02 0.220367E-02 0.187733E-03 0.437651E-03  
52 0.330000E+02 0.165156E+00 0.997070E-02 0.247596E-02 0.279227E-03 0.428934E-03  
53 0.360000E+02 0.110055E+00 0.987053E-02 0.274185E-02 0.395248E-03 0.416801E-03  
54 0.390000E+02 0.708776E-01 0.940597E-02 0.297507E-02 0.502453E-03 0.400711E-03  
55 0.420000E+02 0.446141E-01 0.861541E-02 0.314997E-02 0.577883E-03 0.380540E-03  
56 0.450000E+02 0.281048E-01 0.758869E-02 0.324614E-02 0.610218E-03 0.356595E-03  
57 0.480000E+02 0.184668E-01 0.644155E-02 0.325181E-02 0.598580E-03 0.329568E-03  
58 0.510000E+02 0.133130E-01 0.528894E-02 0.316541E-02 0.549855E-03 0.300434E-03  
59 0.540000E+02 0.108229E-01 0.422457E-02 0.299508E-02 0.475461E-03 0.270327E-03  
60 0.560000E+02 0.998969E-02 0.359589E-02 0.284248E-02 0.417993E-03 0.250291E-03  
61 0.590000E+02 0.932049E-02 0.279916E-02 0.256987E-02 0.329108E-03 0.221098E-03  
62 0.620000E+02 0.886265E-02 0.218402E-02 0.226440E-02 0.246103E-03 0.193776E-03  
63 0.650000E+02 0.832526E-02 0.173744E-02 0.194758E-02 0.176346E-03 0.169057E-03  
64 0.700000E+02 0.703230E-02 0.129310E-02 0.144370E-02 0.990134E-04 0.134806E-03  
65 0.750000E+02 0.539663E-02 0.109402E-02 0.101992E-02 0.692496E-04 0.109622E-03  
66 0.800000E+02 0.377282E-02 0.100701E-02 0.702360E-03 0.731366E-04 0.925653E-04  
67 0.850000E+02 0.243786E-02 0.940358E-03 0.488603E-03 0.916964E-04 0.817733E-04  
68 0.900000E+02 0.150318E-02 0.848527E-03 0.358935E-03 0.108501E-03 0.751267E-04  
69 0.950000E+02 0.943375E-03 0.723450E-03 0.287228E-03 0.113731E-03 0.707398E-04  
70 0.100000E+03 0.660792E-03 0.579280E-03 0.248893E-03 0.104891E-03 0.672143E-04  
71 0.105000E+03 0.543048E-03 0.437072E-03 0.225193E-03 0.852093E-04 0.636854E-04  
72 0.110000E+03 0.498064E-03 0.314313E-03 0.204450E-03 0.609305E-04 0.597352E-04  
73 0.115000E+03 0.467031E-03 0.220667E-03 0.181247E-03 0.385517E-04 0.552579E-04  
74 0.120000E+03 0.422365E-03 0.158301E-03 0.154683E-03 0.227827E-04 0.503338E-04  
75 0.125000E+03 0.358689E-03 0.124152E-03 0.126481E-03 0.156098E-04 0.451435E-04  
76 0.130000E+03 0.282912E-03 0.112269E-03 0.993730E-04 0.164675E-04 0.399229E-04  
77 0.135000E+03 0.206545E-03 0.115547E-03 0.759210E-04 0.231910E-04 0.349411E-04  
78 0.140000E+03 0.140944E-03 0.126914E-03 0.578147E-04 0.332297E-04 0.304812E-04  
79 0.150000E+03 0.727793E-04 0.150870E-03 0.387541E-04 0.563084E-04 0.241091E-04  
80 0.160000E+03 0.972802E-04 0.156307E-03 0.359815E-04 0.797514E-04 0.217716E-04  
81 0.170000E+03 0.166059E-03 0.146102E-03 0.385568E-04 0.100461E-03 0.222061E-04  
82 0.180000E+03 0.201605E-03 0.139242E-03 0.399432E-04 0.109350E-03 0.228033E-04

### 3.4 Other Minor Changes

We had made some additional minor changes for users convenience. If the potential adjustment option of the OPTMAN code is used then the user is able to see the individual  $\chi^2$  value for each experimental data point used in adjustment, and calculated values compared with experimental ones with appropriate errors (In the previous version only global  $\chi^2$  for each run was outputted).

## 4 Conclusions

Following the ISTC B-1319 project's Working Plan we extended OPTMAN code's options by including the possibility of accounting dispersive relationships between imaginary and real parts of optical potential used. We also developed theoretical approaches and computational algorithms allowing predictions of isobaric analog states excitations in  $p + A(N, Z) \rightarrow n' + A(N - 1, Z + 1)$  reactions to test the Lane consistency of the potential and improve isovector potential fits. A new Lane consistent approach to define the Coulomb correction of proton optical potential has been proposed. Incorporation of relativistic generalization algorithm identical to one used in the ECIS code, now allows to get the same computational results for the two codes with identical inputs. New output files now allow to use them as inputs for the EMPIRE evaluation system.

Physical and mathematical ideas of the work have been developed in close cooperation among the Collaborating Parties, which are Japan and European Union and with support from the International Atomic Energy Agency (IAEA) Nuclear Data Section. The present version of modernized OPTMAN code was employed to derive regional nuclear optical potentials applicable up to 200 MeV incident nucleon energies for actinides ( $90 \leq Z \leq 98$ ), Ta – W – Hf nuclei and also a number of local potentials for nuclei as  $^{55}\text{Mn}$ ,  $^{103}\text{Ru}$ ,  $^{197}\text{Au}$ , etc, these derived high-quality dispersive optical potentials are already being used for high energy file evaluation in a cooperative work with our Collaborating Parties and the IAEA Nuclear Data Section. All these potentials are included in the latest version of the RIPL (Reference Input Parameter Library) library, which is an international database maintained by the IAEA.

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