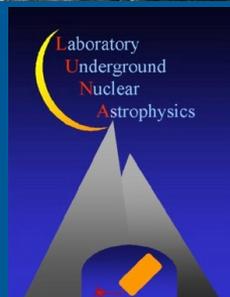


Nuclear fusion reaction measurements at LUNA

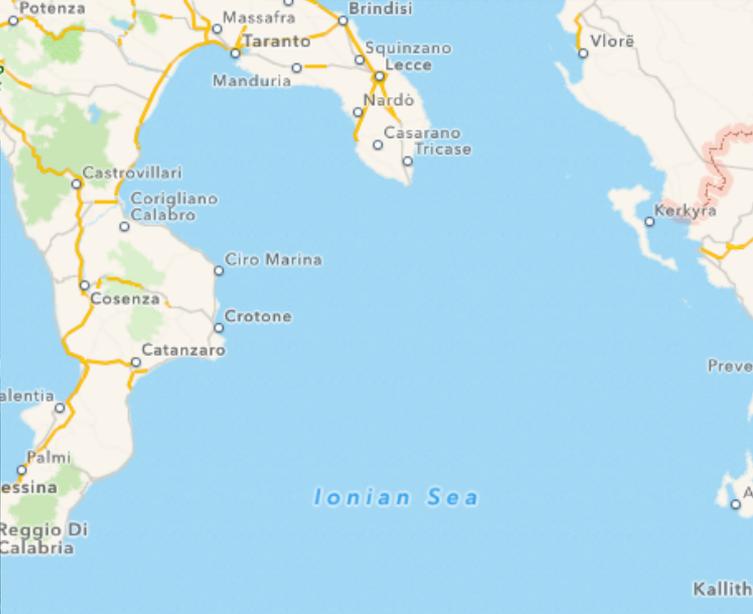
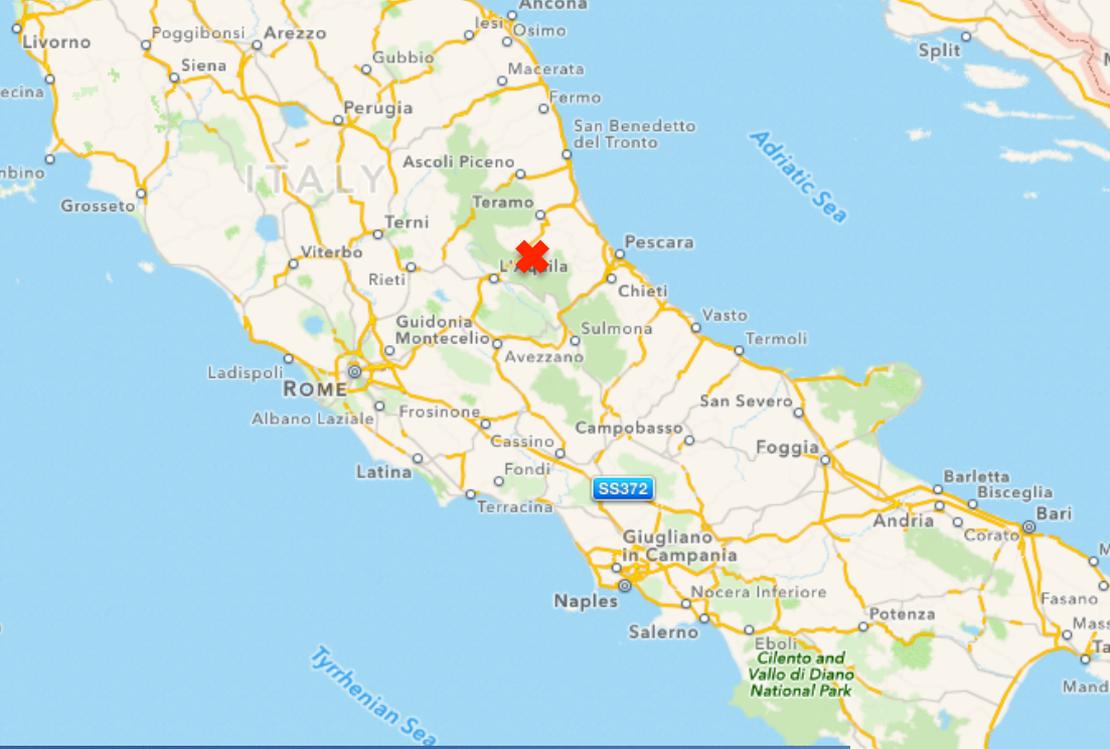


Marcell P. Takács



hzdr

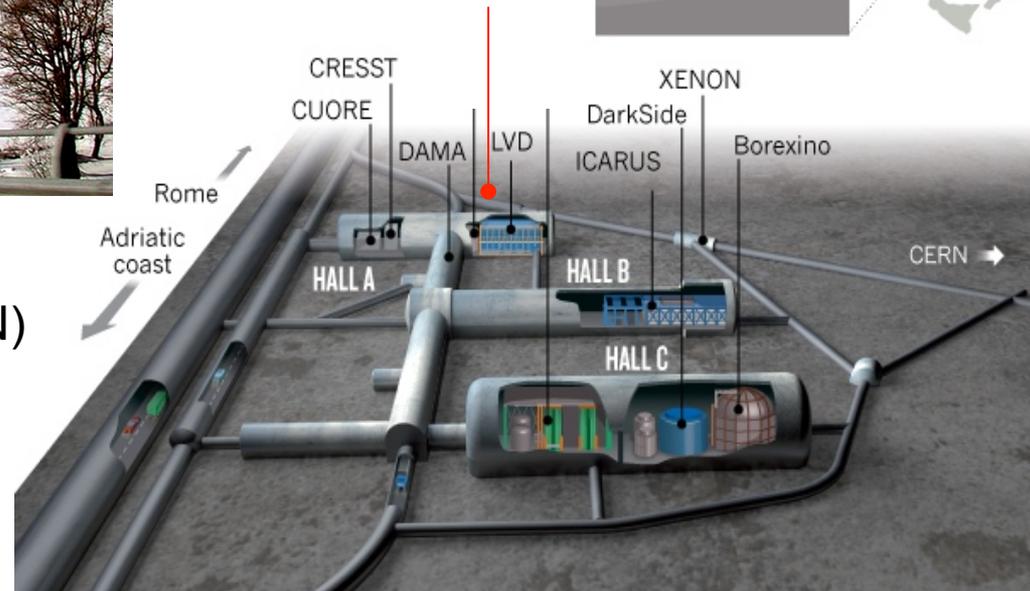
HELMHOLTZ
ZENTRUM DRESDEN
ROSSENDORF



Laboratory for **U**nderground **N**uclear **A**strophysics



Luna



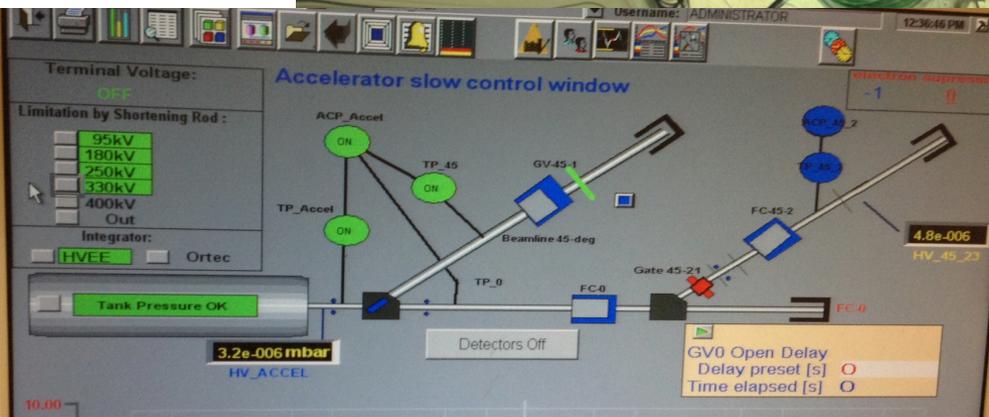
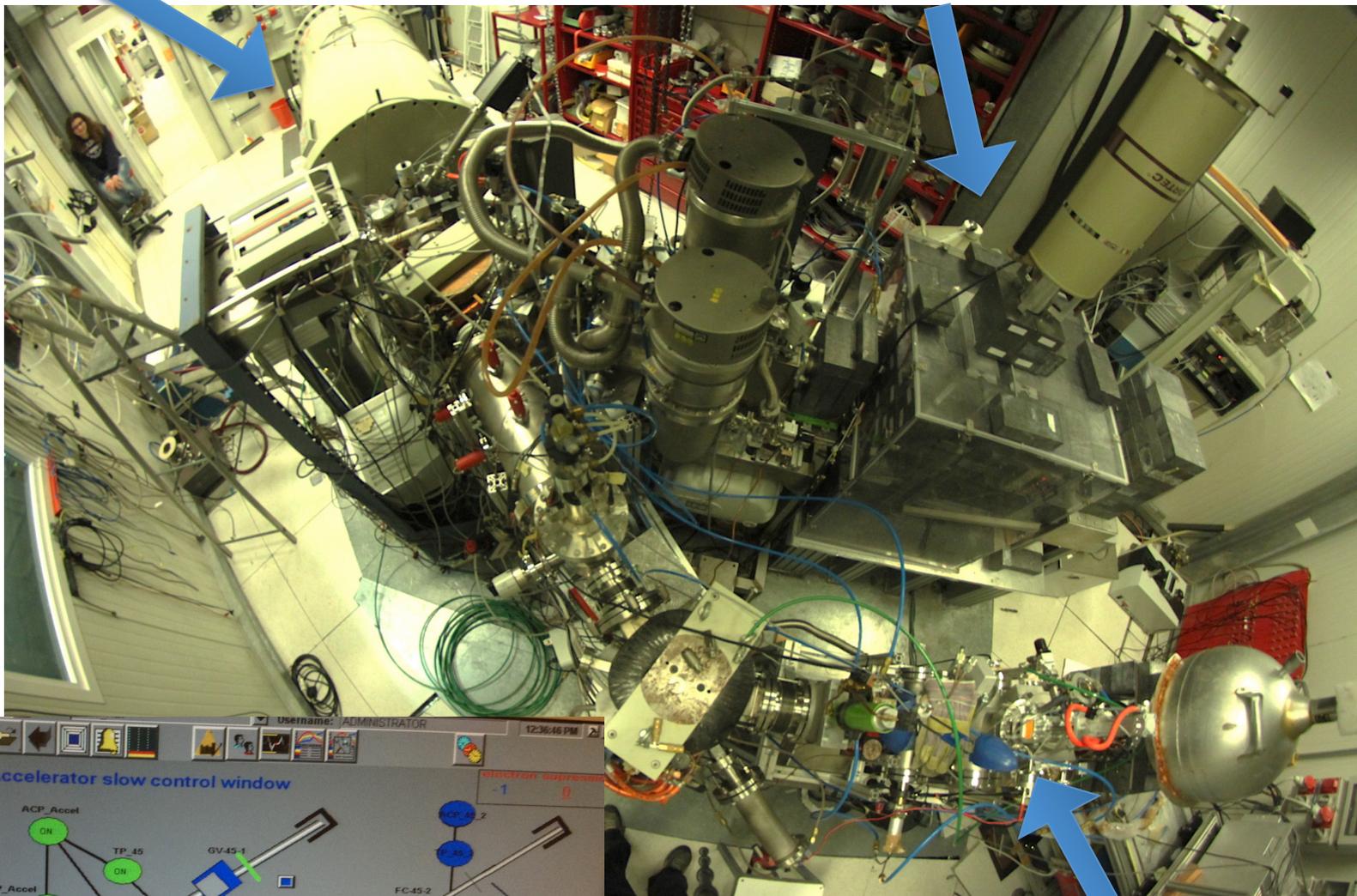
Gran Sasso National Laboratory (INFN)

LUNA-II 400 kV



400 kV accelerator

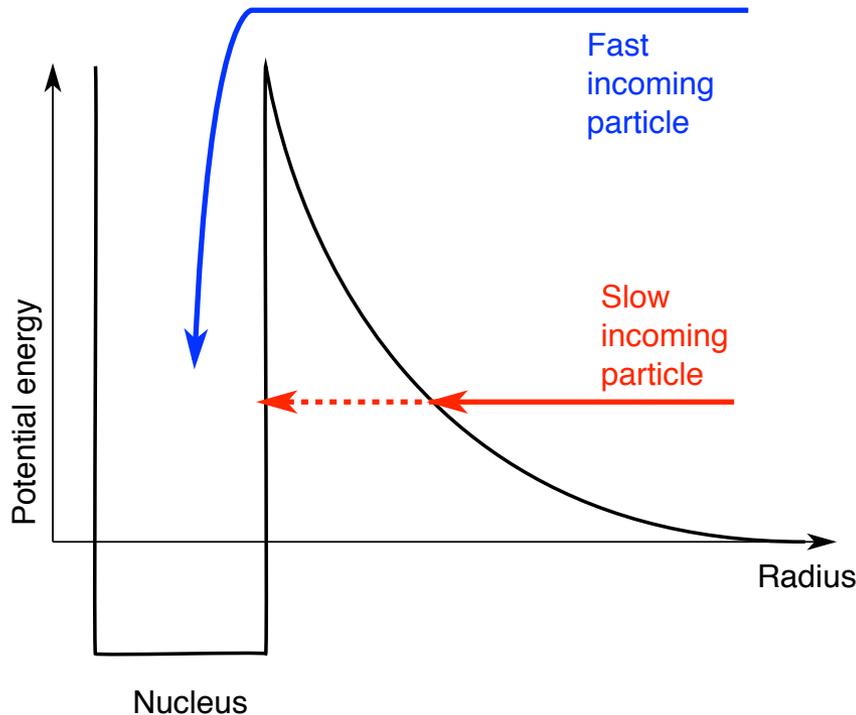
gas target beamline



solid target beamline



Nuclear reaction cross section σ for low-energy charged particles



Typical Coulomb barrier height : \sim MeV

Typical stellar temperature $k_B * T \sim$ keV



Definition of the astrophysical S-factor $S(E)$:

$$\sigma(E) = \frac{S(E)}{E} \exp\left[-2\pi Z_1 Z_2 \alpha \left(\frac{\mu c^2}{2E}\right)^{0.5}\right]$$

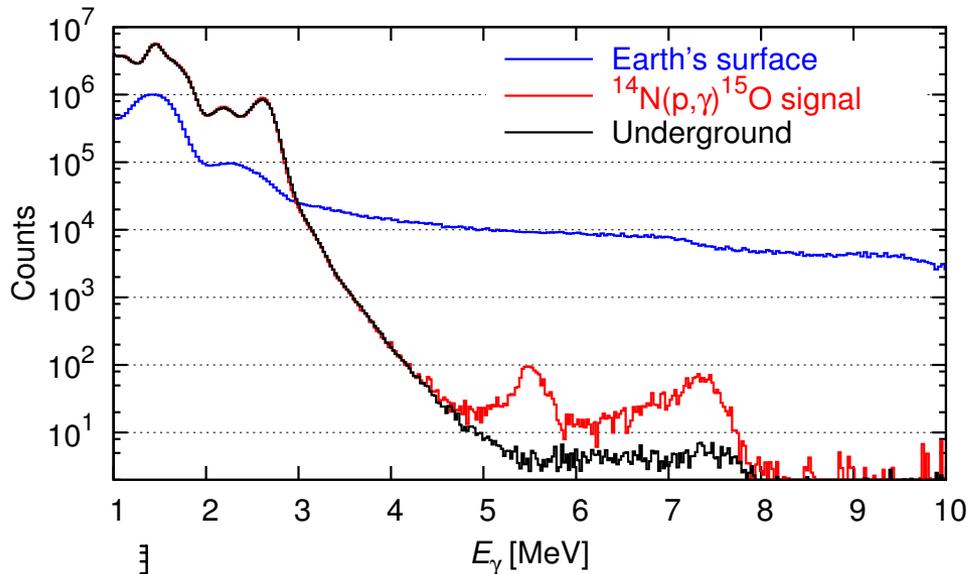


Very low cross sections to be measured!



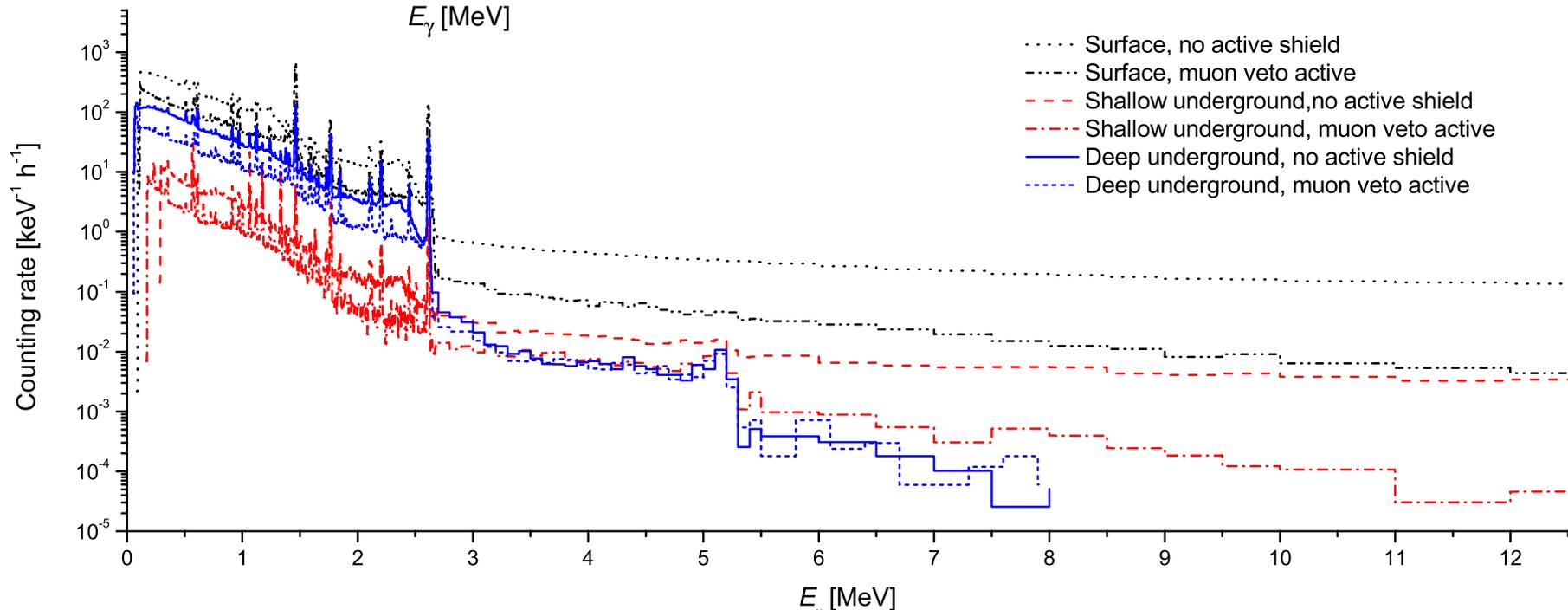
Need for low background!

Uniquely low no-beam background at LUNA



BGO γ -detector (free-running)

HPGe γ -detector
(with cosmic-ray veto)



Thermonuclear reaction rate

- Nuclear reaction represent the main energy source of stars



- Total energy released \sim rate of nuclear reaction

$$R_{01} = \frac{N_R}{Vt} = \sigma(v) \frac{N_0}{V} \frac{N_1}{V} v = n_0 n_1 v \sigma(v)$$

- In the stellar plasma the relative velocity of nuclei **0** and **1** is not constant, but follows Maxwell-Boltzmann distribution

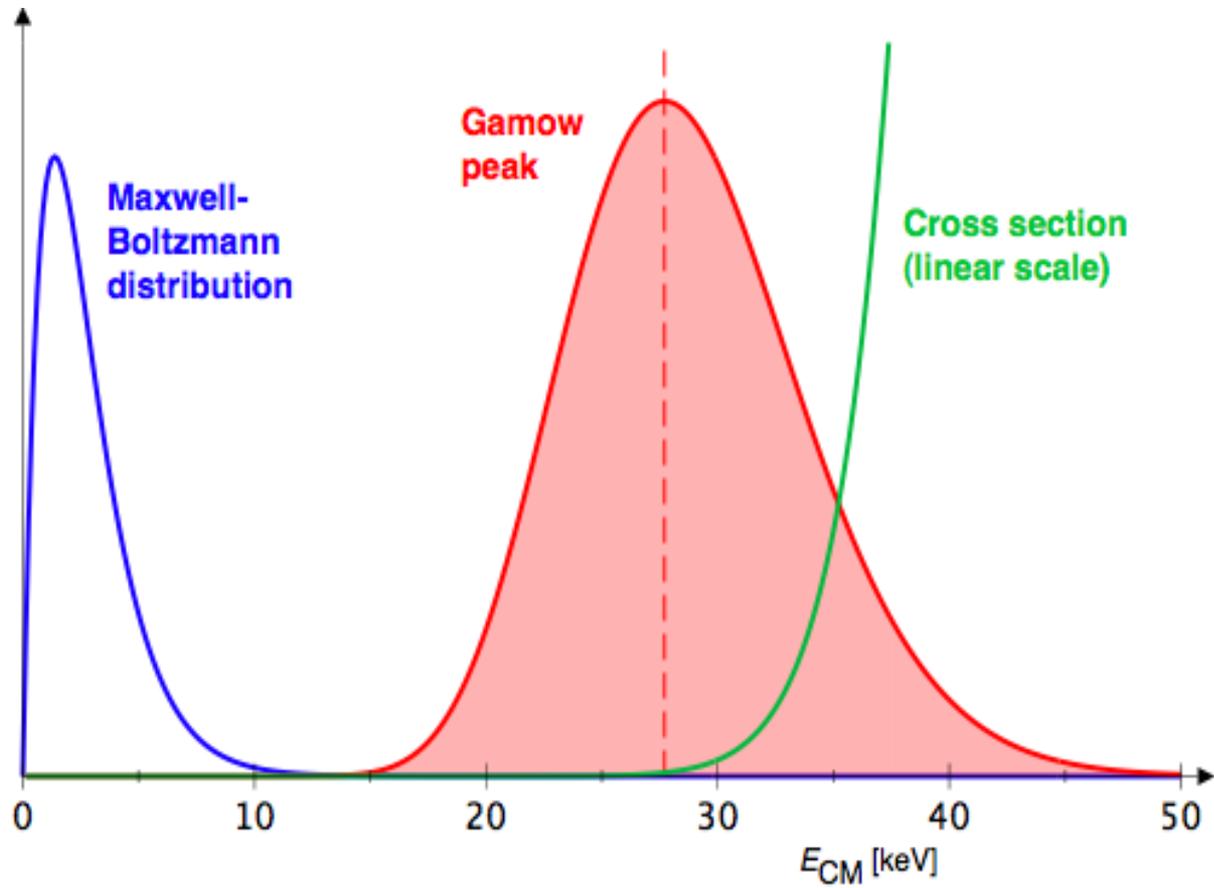
$$P(v)dv = \left(\frac{m_{01}}{2\pi kT} \right)^{3/2} e^{-\frac{m_{01}v^2}{2kT}} 4\pi v^2 dv$$

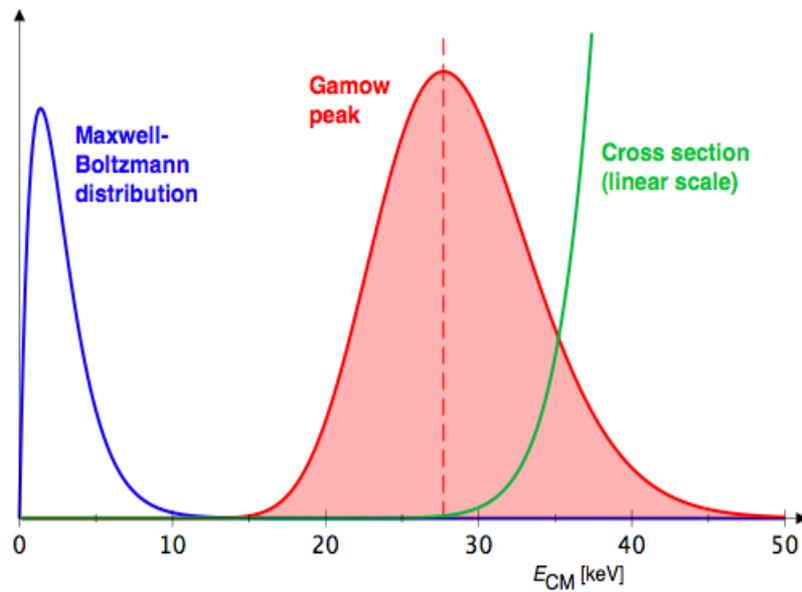
$$P(E)dE = \frac{2}{\sqrt{\pi}} \frac{1}{(kT)^{3/2}} \sqrt{E} e^{-\frac{E}{kT}} dE$$

- Reaction rate per particle pair

$$R_{01} = n_0 n_1 \int v P(v) \sigma(v) dv \equiv n_0 n_1 \langle \sigma v \rangle$$

$$\langle \sigma v \rangle = \sqrt{\frac{8}{\pi m_{01}}} \frac{1}{(kT)^{3/2}} \int_0^{\infty} \underbrace{E}_{\text{green}} \underbrace{\sigma(E)}_{\text{blue}} e^{-\frac{E}{kT}} dE$$



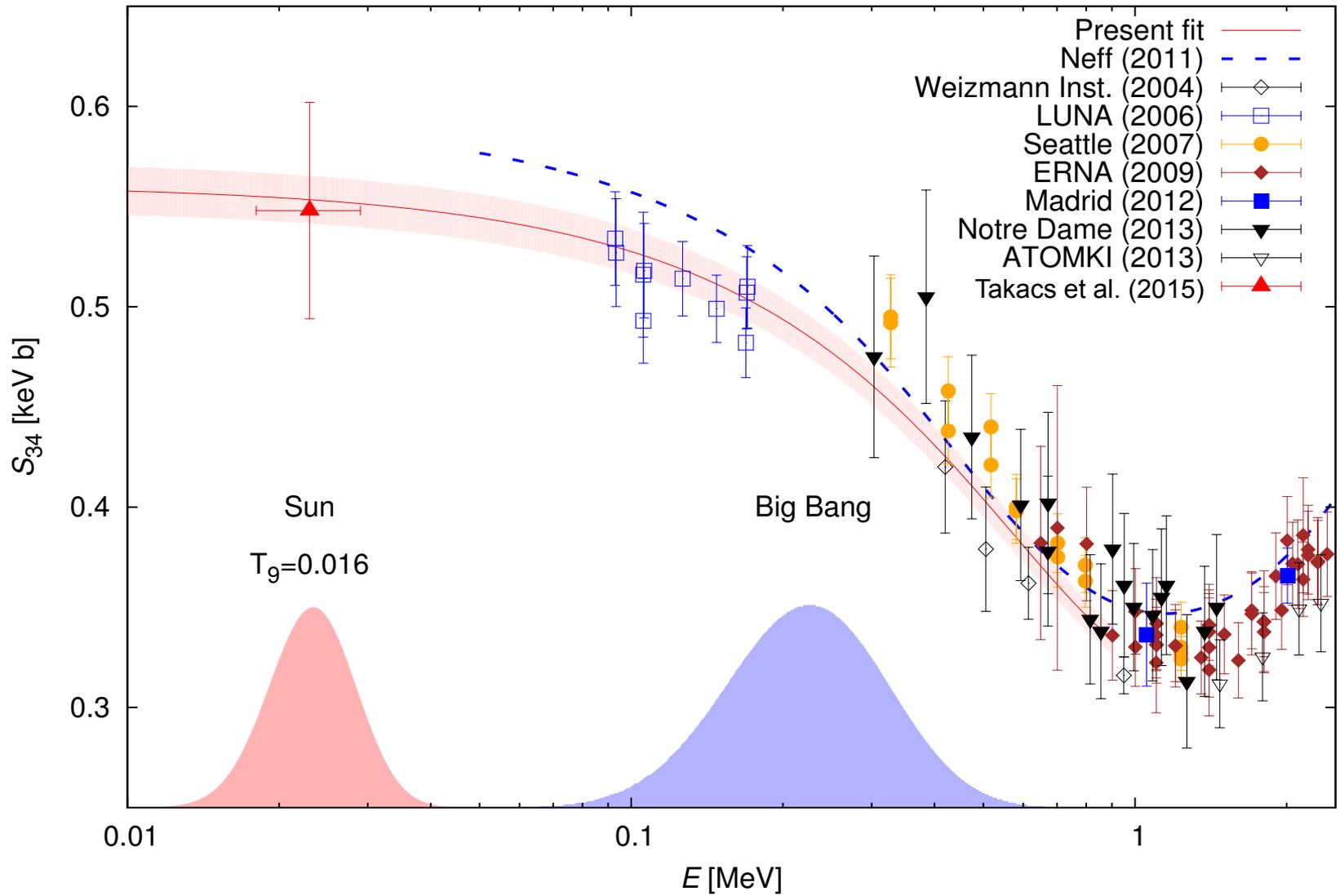


At which energies do the reactions take place in a plasma?

Assuming 10^{16} s^{-1} beam
 10^{18} at/cm^2 target
 10^{-2} detection efficiency

Scenario	Reaction	E_G [keV]	σ [barn]	Detected events/hour
Sun (16 MK)	$^3\text{He}(\alpha, \gamma)^7\text{Be}$	23	10^{-17}	10^{-9}
	$^{14}\text{N}(p, \gamma)^{15}\text{O}$	28	10^{-19}	10^{-11}
AGB stars (80 MK)	$^{14}\text{N}(p, \gamma)^{15}\text{O}$	81	10^{-12}	10^{-4}
Big bang (300 MK)	$^3\text{He}(\alpha, \gamma)^7\text{Be}$	160	10^{-9}	10^{-1}
	$^2\text{H}(\alpha, \gamma)^6\text{Li}$	96	10^{-11}	10^{-3}

Too low, even for LUNA



The presence of resonances in the Gamow window can greatly enhance the cross section



Their contribution will dominate the reaction rate

E_{proton}	$^{22}\text{Ne}(p,\gamma)^{23}\text{Na}$	E_{level}	
661 keV	_____	9426	$3/2^-$
638 keV	_____	9405	$1/2^-$
634 keV	_____	9401	
630 keV	_____	9396	$7/2^-$
552 keV	_____	9322	
514 keV	_____	9286	$3/2, 5/2^+$
479 keV	_____	9252	$1/2^+$
436 keV	_____	9211	$3/2^+$
394 keV	_____	9171	
369 keV	_____	9147	
333 keV	_____	9113	
323 keV	_____	9103	
291 keV	_____	9072	
260 keV	=====	9042	$7/2, 9/2^+$
256 keV	=====	9039	$15/2^+$
215 keV	_____	9000	
189 keV	_____	8975	$5/2^+$
158 keV	=====	8945	$7/2^-$
156 keV	=====	8944	$3/2^+$
105 keV	-----	8895	$1/2^+$
71 keV	-----	8862	$1/2^+$
37 keV	=====	8830	$1/2^+$
29 keV	=====	8822	$9/2^-$
3 keV	_____	8797	

Supernovae

Novae

AGB

$$\langle \sigma v \rangle = \sqrt{\frac{8}{\pi m_{01}}} \frac{1}{(kT)^{3/2}} \int_0^{\infty} E \sigma(E) e^{-\frac{E}{kT}} dE$$

- For isolated, narrow resonances the cross section is given by the Breit-Wigner formula

$$\sigma_{BW} = \frac{\lambda^2}{4\pi} \frac{(2J+1)}{(2j_0+1)(2j_1+1)} (1 + \delta_{01}) \frac{\Gamma_a \Gamma_b}{(E_r - E)^2 + \frac{\Gamma^2}{4}}$$

λ : de Broglie wavelength $\lambda = \frac{2\pi\hbar}{\sqrt{2m_{01}E}}$

j_i : spin of nuclei 0 and 1

J : spin of the resonance (compound nucleus)

Γ_i : partial width of the entrance (a) and exit (b) channel

Γ : total width of the resonance

$$\omega = \frac{(2J+1)}{(2j_0+1)(2j_1+1)} (1 + \delta_{01})$$

$$\langle \sigma v \rangle = \frac{\sqrt{2\pi\hbar^2}}{(m_{01}kT)^{3/2}} \omega \int_0^\infty \frac{\Gamma_a \Gamma_b}{(E_r - E)^2 + \frac{\Gamma^2}{4}} e^{-\frac{E}{kT}} dE$$

- For narrow resonance the partial width (Γ_a, Γ_b) and the Maxwell-Boltzmann factor approx. constant over the resonance width Γ

$$\langle \sigma v \rangle = \frac{\sqrt{2\pi\hbar^2}}{(m_{01}kT)^{3/2}} \omega \frac{\Gamma_a \Gamma_b}{\Gamma} 2e^{-\frac{E_r}{kT}} \int_0^\infty \frac{\frac{\Gamma}{2}}{(E_r - E)^2 + \frac{\Gamma^2}{4}} dE$$

$$\langle \sigma v \rangle = \left(\frac{2\pi}{m_{01}kT} \right)^{3/2} \hbar^2 e^{-\frac{E_r}{kT}} \omega \underbrace{\frac{\Gamma_a \Gamma_b}{\Gamma}}$$

- Resonance strength:

$$\omega\gamma = \omega \frac{\Gamma_a \Gamma_b}{\Gamma}$$

Experimental determination of $\omega\gamma$

- In nuclear physics experiments the measured quantity is often the **yield**:

$$Y = \frac{\text{total number of reactions}}{\text{total number of incident particles}} = \frac{N_R}{N_b}$$

- Yield and cross section $Y = \sum_i \sigma_i n_i \Delta x_i$

$$Y = \int_{\text{target}} \sigma(x) n(x) dx = \int_{\text{target}} \sigma(x) n(x) \frac{dE}{dx} \frac{dx}{dE}$$

$$Y = \int_{E_p - \Delta E}^{E_p} \frac{\sigma(E)}{\varepsilon(E)} dE$$

Stopping power:

$$\varepsilon(E) = -\frac{1}{n_t} \frac{dE}{dx} \quad n_t = \frac{N_t}{V}$$

- Assuming Breit-Wigner cross section the integral can be solved analytically

$$Y = \frac{\lambda_r^2}{2\pi \epsilon_r} \omega \gamma \left[\arctan\left(\frac{E_p - E_r}{\Gamma/2}\right) - \arctan\left(\frac{E_p - E_r - \Delta E}{\Gamma/2}\right) \right]$$

- The maximum yield is achieved if $E_p = E_r + \frac{\Delta E}{2}$

$$Y_{\max} = \frac{\lambda_r^2}{\pi \epsilon_r} \omega \gamma \arctan\left(\frac{\Delta E}{\Gamma}\right)$$

- For charged particle induced capture reactions - for example: (p, γ)

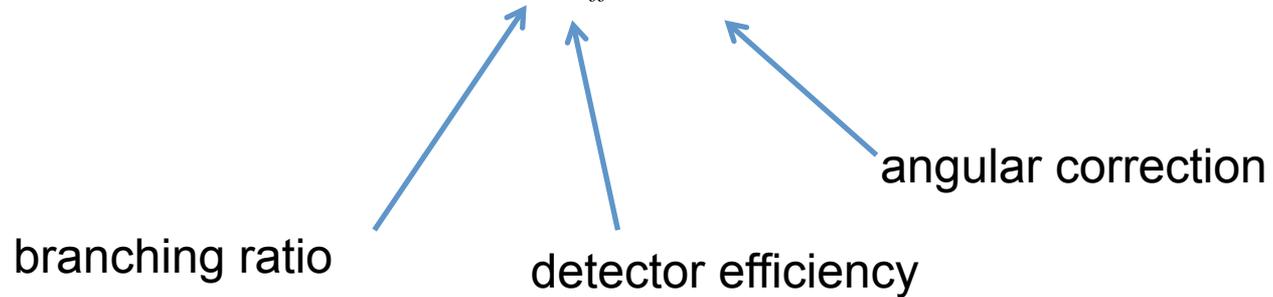
$$\Gamma_p \ll \Gamma_\gamma \quad \longrightarrow \quad \Gamma = \Gamma_p + \Gamma_\gamma \approx \Gamma_\gamma \approx (1 \text{ meV} - 1 \text{ eV})$$

$$\Delta E \approx (10 \text{ keV})$$

$$Y_{\Delta E \rightarrow \infty} = \frac{\lambda_r^2}{2 \epsilon_r} \omega \gamma$$

- “Real” experimental yield calculation

$$Y = \frac{N_R}{N_b} = \frac{N_{peak}}{N_b B \eta_{eff} W(\vartheta)}$$



- Final formula for the experimental resonance strength

$$\omega\gamma = \frac{2}{\lambda_r^2} \epsilon_r \frac{N_{peak}}{N_b B \eta_{eff} W(\vartheta)}$$

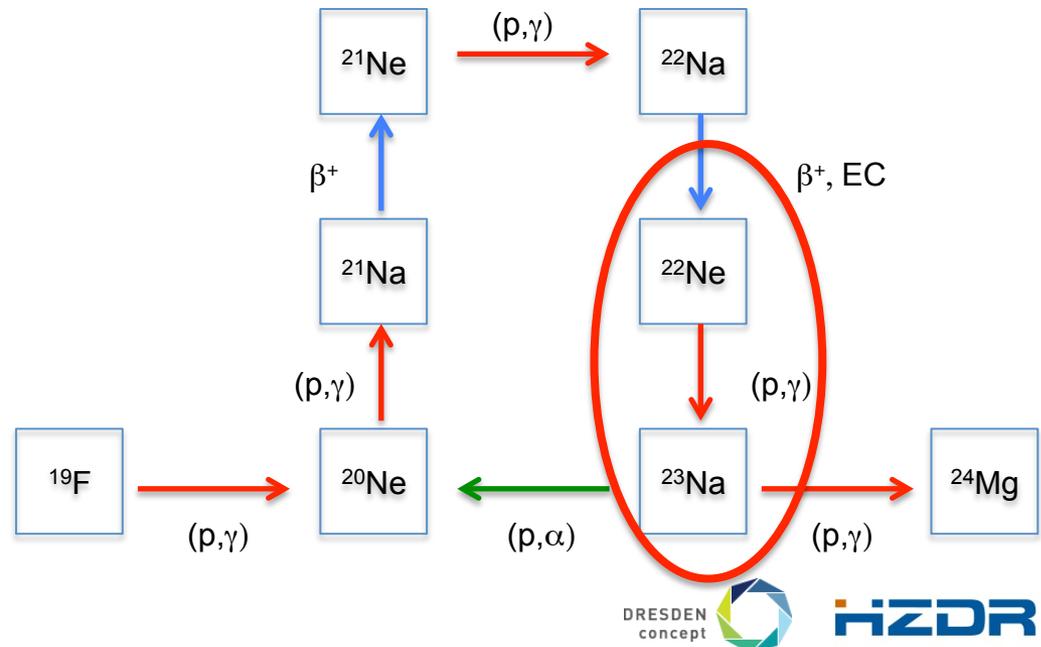
$^{22}\text{Ne}(p,\gamma)^{23}\text{Na}$

■ Importance

- Second to the slowest reaction in the Ne-Na cycle of hydrogen burning
- Strong effect on the abundances of ^{22}Ne and ^{23}Na
- ^{22}Ne is important for neutron-capture driven nucleosynthesis: $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$
- Large uncertainties regarding its reaction rate

■ Astrophysical sites

- AGB stars
- Novae
- Supernovae



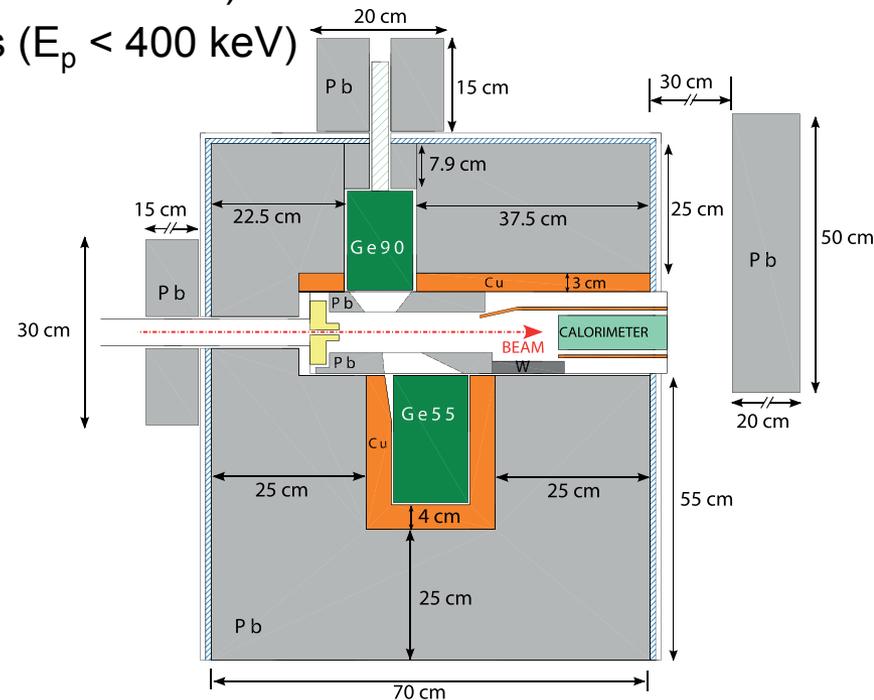
Experiment at LUNA

■ Differentially pumped gas target system

- Windowless system
- 3 pumping stages
- Gas recirculation with purifier

■ HPGe-based first phase

- 2 ultra low background (ULB) detectors
- Placement: 55° (137% rel. eff.) and 90° (90% rel. eff.)
- Study of selected low energy resonances ($E_p < 400$ keV)
- Limit uncertainty by angular distribution



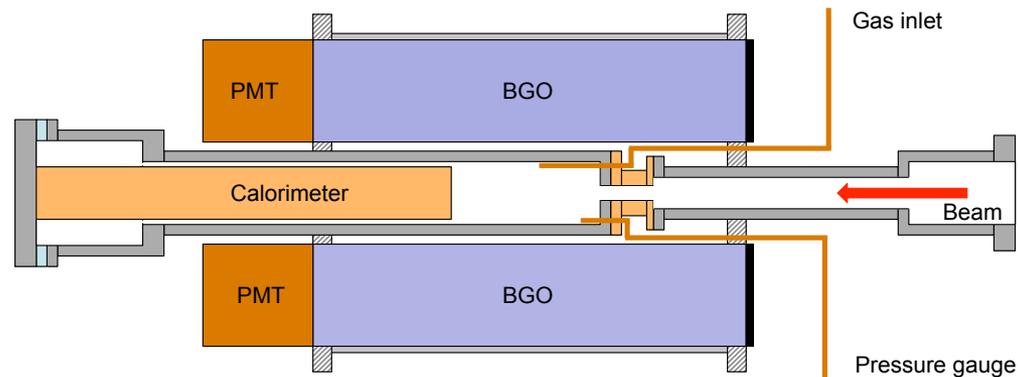
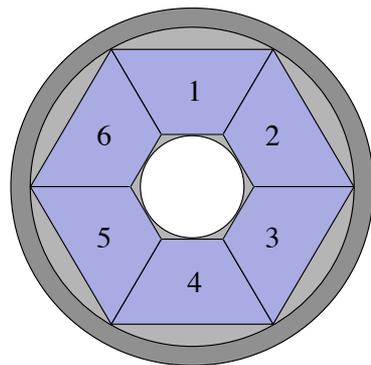
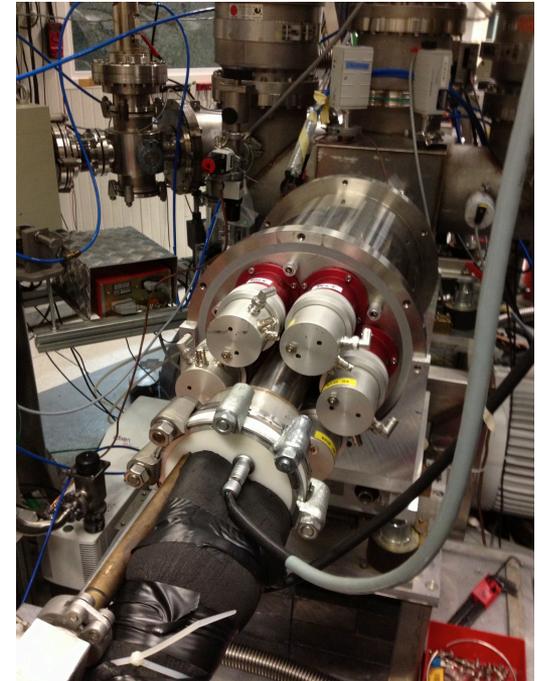
E_{proton}	$^{22}\text{Ne}(p,\gamma)^{23}\text{Na}$	E_{level}	
	394 keV	9171	
	369 keV	9147	
	333 keV	9113	
	323 keV	9103	
	291 keV	9072	
$\omega\gamma = (6.89 \pm 0.16) \times 10^{-6} \text{ eV}$	260 keV	9042	$7/2, 9/2^+$
	256 keV	9039	$15/2^+$
	215 keV	9000	
$\omega\gamma = (1.87 \pm 0.06) \times 10^{-6} \text{ eV}$	189 keV	8975	$5/2^+$
$\omega\gamma = (1.48 \pm 0.10) \times 10^{-7} \text{ eV}$	158 keV	8945	$7/2^-$
	156 keV	8944	$3/2^+$
$\omega\gamma \leq 7.6 \times 10^{-9} \text{ eV}$	105 keV	8895	$1/2^+$
$\omega\gamma \leq 1.5 \times 10^{-9} \text{ eV}$	71 keV	8862	$1/2^+$
	37 keV	8830	$1/2^+$
	29 keV	8822	$9/2^-$
	3 keV	8797	

LUNA-HPGe

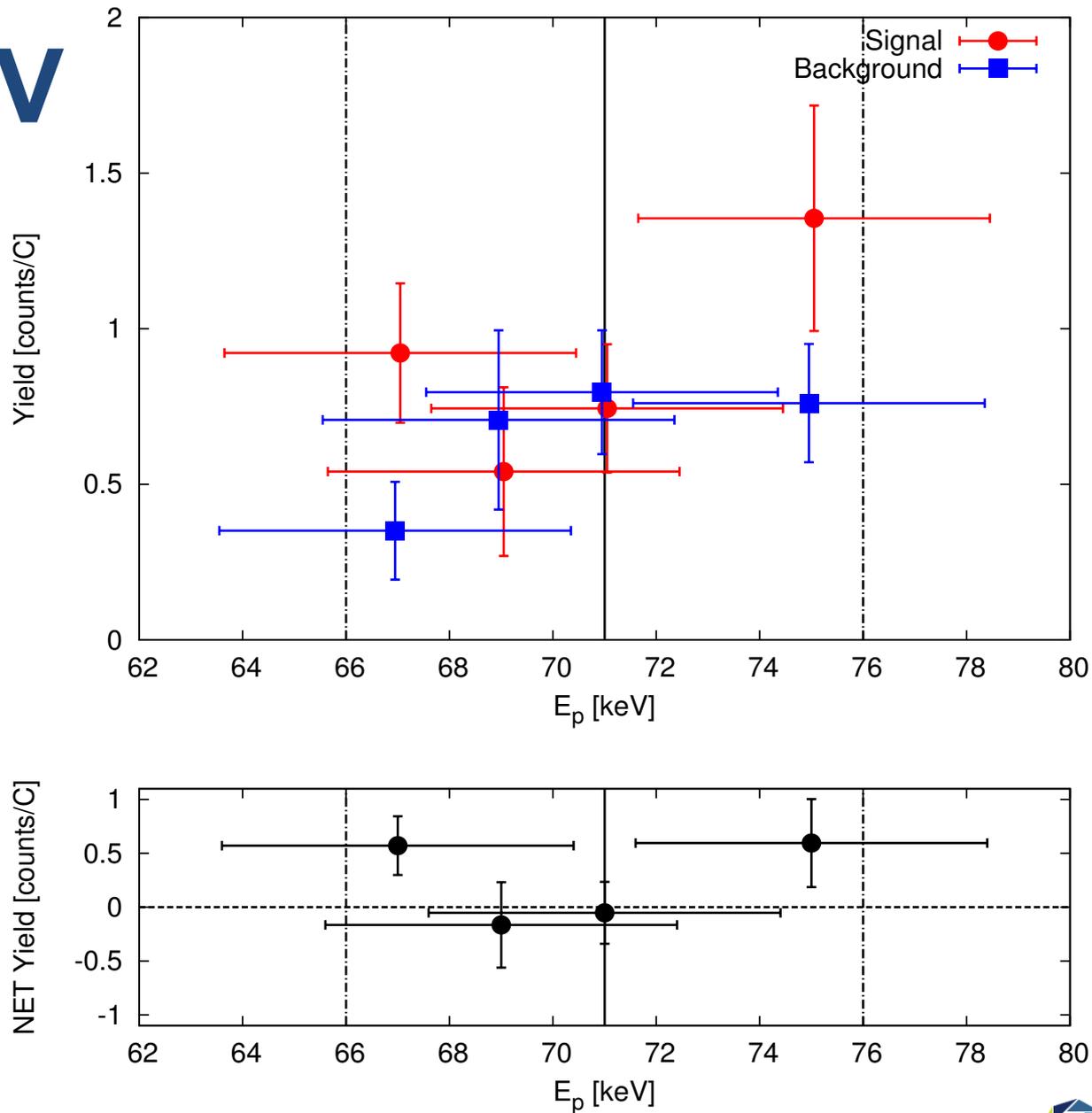
Experiment at LUNA

■ BGO-based second phase

- 4π bismuth germanate summing crystal
- 6 sectors with separate PMTs
- Positioned around the gas target chamber
- Lowest energy resonances ($E_p < 200$ keV)
- Study of direct capture reaction



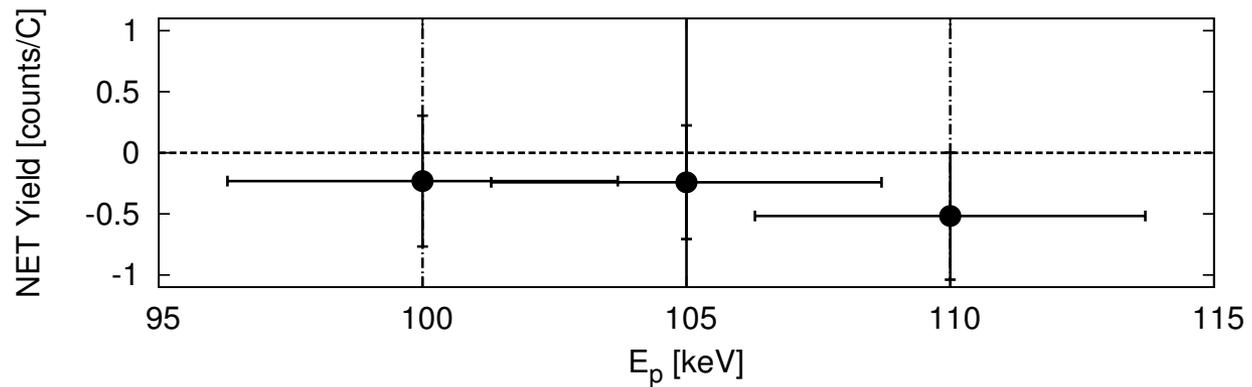
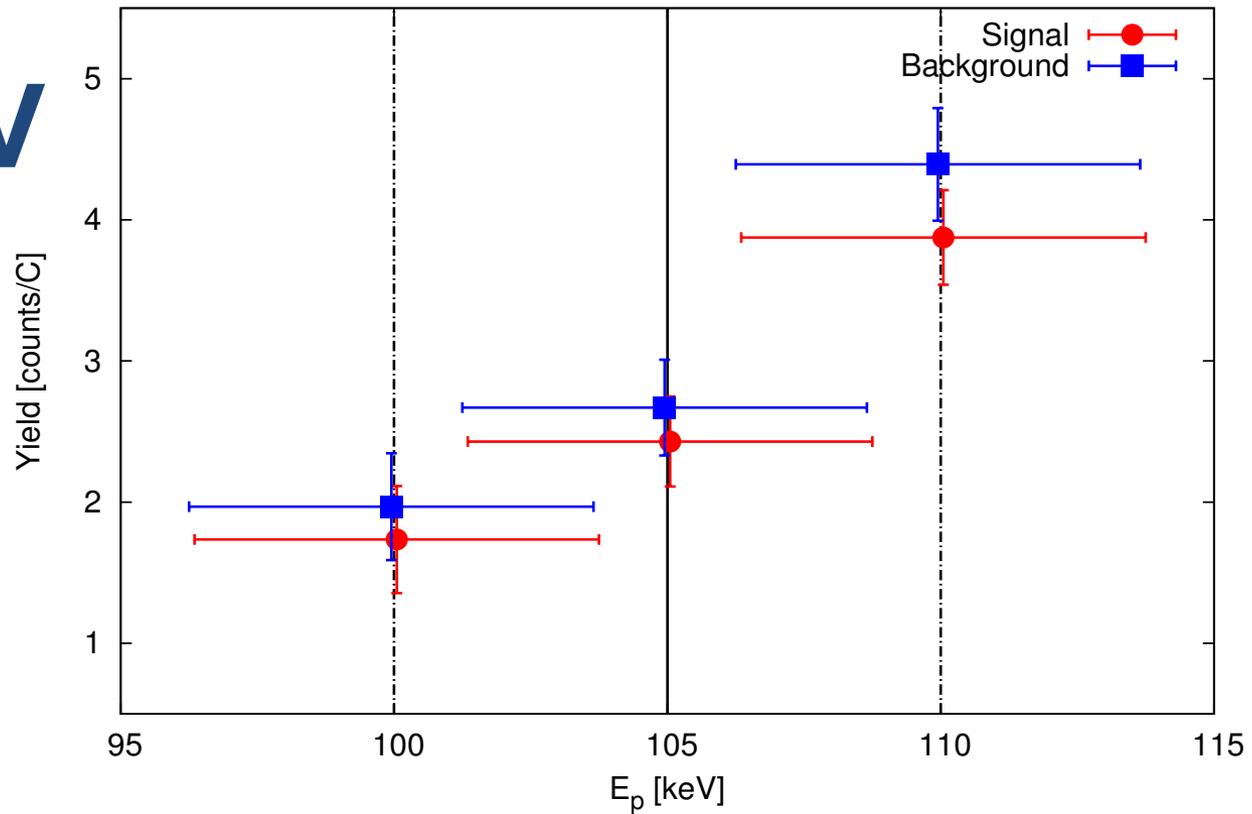
71 keV



Preliminary

(quoted errors are statistical only)

105 keV



Preliminary

(quoted errors are statistical only)

Results – Resonant capture

E_p [keV]	Upper limit LUNA-BGO	Upper limit LUNA-HPGe
	$\omega\gamma$ [eV]	$\omega\gamma$ [eV]
67	$\leq 5.4 \times 10^{-11}$	$\leq 1.5 \times 10^{-9}$
69	$\leq 3.1 \times 10^{-11}$	
71	$\leq 2.5 \times 10^{-11}$	
75	$\leq 8.1 \times 10^{-11}$	
100	$\leq 6.0 \times 10^{-11}$	$\leq 7.6 \times 10^{-9}$
105	$\leq 4.9 \times 10^{-11}$	
110	$\leq 4.1 \times 10^{-11}$	

Preliminary

(F. Cavanna et al., PRL 115, 252501, 2015)

Numbers based on unbound / bound profile likelihood method (90% CL)

W. Rolke et al., NIM A 551 (2005) 493

Creed of LUNA

1. Reduce the background!
2. Measure nuclear reaction at (or near) the relevant energies!
3. Use high beam intensity!
4. Have great patience!

