

AUSTRALIAN ATOMIC ENERGY COMMISSION

RESEARCH ESTABLISHMENT LUCAS HEIGHTS

keV NEUTRON CAPTURE IN ZIRCONIUM-91

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* Oak Ridge National Laboratory, Oak Ridge, Tennessee, USA. Research sponsored in part by the USAEC under contract to Union Carbide Corporation.

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ABSTRACT

The neutron capture cross section of $^{91}{\rm Zr}$ has been measured with high resolution (Δ E/E ~ 0.2 per cent) between 3 and 30 keV. Values of the ${\rm g}\Gamma_{\rm n}\Gamma_{\gamma}/\Gamma$ for 119 resonances in this energy range have been obtained. The average capture cross section is consistent with values of ${\langle}\Gamma_{\gamma}{\rm s}{\rangle} = {\langle}\Gamma_{\gamma}{\rm p}{\rangle} = 200$.

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CAPTURE; CROSS SECTIONS; ENERGY DEPENDENCE; KEV RANGE 10-100; NEUTRON REACTIONS; RESONANCE; ZIRCONIUM 91 TARGET

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Figure 1 Uncorrected experimental capture yield

Figure 2 Plot of capture kernel versus neutron energy

Figure 3 Staircase plot of number of resonances below energy E

Figure 4 Comparison of experimental data with statistical model calculations of capture cross section assuming $S_0=0.5\times 10^{-4}$, $S_1=5\times 10^{-4}$, $\Gamma_{\gamma S}=\Gamma_{\gamma p}=0.2$ eV and $\langle D_S\rangle=660$ eV

INTRODUCTION

The neutron capture cross section of $^{91}\mathrm{Zr}$ has been measured between 3 and 30 keV using the neutron capture facility at the 40 m flight path of the Oak Ridge Electron Linear Accelerator (ORELA). In previous studies of ⁹⁰,92,94Zr [Boldeman *et al.* 1975a,b], strong non-statistical effects were observed in the neutron capture cross section which were attributed to single-particle effects associated with the N=50 closed shell and the peak in the p-wave strength function at A-95. The non-statistical effect, consisting of significant correlations between the reduced neutron and radiative width for p-wave resonances, found a satisfactory quantitative explanation in terms of the optical model formulation of the valence theory and a doorway state mechanism. In contrast, for neutron capture in odd A $^{91}\mathrm{Zr}$ these effects will be between one and two orders of magnitude less, largely because of smaller reduced neutron widths, and the capture cross section of $^{91}\mathrm{Zr}$ should behave in a statistical manner.

From the data viewpoint, information on $^{91}\mathrm{Zr}$ is of value because of the use of Zr in reactor systems and the occurrence of $^{91}\mathrm{Zr}$ as an important fission product (WRENDA Requests Nos. 691156, 691155).

EXPERIMENTAL METHOD AND DATA ANALYSIS

The experimental system has been described in considerable detail [Boldeman et al. 1975a, Macklin & Allen 1971]. The neutron capture gamma-ray detector consisted of two non-hydrogenous liquid scintillators with low neutron sensitivity which operated under a pulse height weighting scheme such that the efficiency of the device is a function only of the total gamma-ray energy release. The neutron monitor was a thin $^6\mathrm{Li}$ glass scintillator (0.5 mm thick) located 1 metre upstream of the capture detector and therefore operated in the transmission mode. The relative neutron detection efficiency of the $^6\mathrm{Li}$ monitor was based on an evaluation of the $^6\mathrm{Li}(n,\alpha)$ cross section which included the evaluation of Uttley et al. [1971] below 100 keV and recent experimental data from Fort & Marquette [1972], Coates et al. [1972] and Poenitz [1974] above that energy. The capture detector efficiency was deduced using the $^6\mathrm{Li}$ glass detector and observing the saturated capture yield from the 4.9 eV resonance in Au. An enriched target (88.5 per cent $^{91}\mathrm{Zr}$) was used. Resonances in the other Zr isotopes were removed using the data from Boldeman et al. [1975a,b].

The method of analysis has been described in detail by Musgrove et al. [1974] and is based on the ORNL/RFI code of Sullivan et al. [1969]. The Monte Carlo code uses initial guesses for the resonance neutron (Γ_n) and Γ_n

radiative widths (Γ_{γ}) to determine self-shielding and multiple-scattering corrections. An iterative fit is then made to the capture area for each resonance $(2\pi^2\chi^2\mathrm{g}\Gamma_n\Gamma_{\gamma}/\Gamma)$ by varying whichever of the initial values of Γ_{γ} and Γ_n is the smaller. The analysis of the capture cross section data for $^{91}\mathrm{Zr}$ was somewhat restricted by the absence of a suitable total cross section measurement and therefore by the lack of values of Γ_n for the resonances studied. For several resonances in the present data set, shape analysis of the capture data leading to a value of Γ_n was possible. These values of Γ_n have been used in the analytical procedure. Otherwise it has been assumed that $\Gamma_n > \Gamma_{\gamma}$ for all resonances and a value of Γ_n equal to one fifth of the resolution width ($\Gamma_R\!\!\approx\!\!0.2$ per cent) has been used. The broad resonance at 3.991 keV listed by Mughabghab & Garber [1973] was found to have $\Gamma_n \leq 3$ eV in the present data. The absence of any spin data whatsoever for the resonances studied restricts the resonance analysis to the production of the capture kernel only ($\mathrm{g}\Gamma_n\Gamma_{\gamma}/\Gamma$).

3. RESULTS

Figure 1 shows the capture yield data for the energy region between 3 and 30 keV, while Table 1 lists the output data for the 120 resonances identified in this region. Figure 2 is a staircase plot of the observed levels below 30 keV; this shows that few levels have been missed to 30 keV. If the level density is independent of parity and shows a (2J+1) dependence, an s-wave level spacing of 670 eV is obtained.

Figure 3 is a plot of the capture kernel versus neutron energy. There is no feature of the plot which allows any distinction to be made between s- or p-wave resonances or, for that matter, between the allowed values of g. The most interesting feature of Figure 3 is the absence of large values of the capture kernel and the narrow distribution of the measured values. This behaviour contrasts quite strongly with equivalent data for 90 , 92 , 94 Zr, where the variance of the capture kernels was significantly greater than the mean. This feature of the data is entirely expected for 91 Zr as non-statistical effects are estimated to be between one and two orders of magnitude less than for 90 , 92 , 94 Zr. The major reason for the large reduction in non-statistical effects arises because of the reduced neutron widths of the resonances, which are approximately a factor of 40 less than for 90 Zr.

The average capture cross section is shown for $^{91}{\rm Zr}$ in Figure 4. Also shown is the calculated statistical model capture cross section obtained using the published values of Mughabghab & Garber [1973] for

the strength function, the present value of 670 eV for the s-wave level spacing, and assuming an average radiative width of 200 meV for both s-wave and p-wave resonances. The calculated cross section is in agreement with the experimental values. The 30 keV cross section from the present experiment, i.e. 60 ± 6 mb, is in close agreement with the two previous measurements, 60 ± 6 mb from Kapchigashev [1964] and 59 ± 10 from Macklin et al. [1963].

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TABLE 1

RESONANCE PARAMETERS FOR 91Zr

Energy Γ_n $G^{\Gamma}_n \Gamma_{\gamma} \Gamma$,
3.162±0.006 (1.4) a 0.085±0.009 8.510±0.017 0.073±0.000 3.617±0.007 0.052±0.005 8.527±0.017 0.066±0.000 3.648±0.007 0.054±0.005 8.865±0.018 15±5 0.014±0.003 4.012±0.008 0.064±0.006 8.955±0.018 0.046±0.005 4.283±0.008 0.052±0.005 9.045±0.018 0.071±0.005 4.333±0.009 5±2 0.214±0.021 9.239±0.018 5+3 0.075±0.008 4.754±0.010 0.005±0.008	,
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3.869±0.008 4.012±0.008 4.283±0.008 4.283±0.008 4.333±0.009 5±2 0.064±0.006 0.052±0.005 0.052±0.005 0.052±0.005 0.046±0.005 0.071±0.005 0.069±0.007 0.075±0.008	
4.012±0.008	
4.283±0.008 4.333±0.009 5±2 0.052±0.005 0.109±0.018 0.069±0.007 0.075±0.008 0.075±0.008	
4.333±0.009 5±2 0.214±0.021 9.239±0.018 5 ⁺³ -2 0.075±0.008	
14./54±0.010 0.00210.000	
9 310+0 010	
4.985±0.010 0.045±0.005 0.037±0.006	
5.367±0.011 0.012±0.012	
5.533±0.011 15±5 0.052±0.006	
5.640±0.011	
10.14 ±0.02 0.106±0.011	
10.53 ±0.02 0.084±0.009	
6.098±0.012	
6.187±0.012	
6.481±0.013	
6.768±0.014 0.044±0.005 11.08 ±0.02 0.078±0.008	
6.867±0.014 0.050±0.005	
7.047±0.014 5±3 0.151±0.015 11.13.40.00	
7.134±0.014 0.089±0.009 11.24 +0.00	
7.267±0.015 4 ⁺⁴ 0.005+0.015 0.063±0.007	
7.362±0.015	
12.18 ±0.02 0.017+0.004	
7.768±0.015 5^{+3}_{-2} 0.203±0.020 12.23 ±0.02 0.147±0.015	

Ene rgy	Γ _n	gr r/r	Energy	Γ_	gľ _n ľ _/ ľ
(keV)	(eV)	(eV)	(keV)	n (eV)	n γ' (eV)
12.33±0.02		0.056±0.006	18.56±0.04		0.050±0.007
12.56±0.03		0.143±0.014	18.66±0.04		0.155±0.016
12.94±0.03		0.059±0.006	19.52±0.04		0.159±0.016
13.17±0.03		0.077±0.008	19.61±0.04		0.069±0.008
13.27±0.03		0.078±0.008	19.78±0.04		0.086±0.010
13.32±0.03		0.083±0.009	19.83±0.04		0.074±0.009
13.58±0.03		0.048±0.006	20.04±0.04		0.060±0.009
13.71±0.03		0.141±0.014	20.19±0.04		0.144±0.015
13.81±0.03		0.058±0.006	20.20±0.04	20±10	0.168±0.017
14.10±0.03		0.077±0.008	20.27±0.04		0.069±0.009
14.20±0.03		0.050±0.006	20.32±0.04		0.056±0.808
14.26±0.03		0.083±0.009	20.41±0.04		0.071±0.009
14.60±0.03		0.111±0.012	20.65±0.04		0.089±0.010
14.86±0.03		0.071±0.008	20.92±0.04		0.127±0.013
15.19±0.03		0.165±0.017	21.25±0.04		0.075±0.009
15.25±0.03		0.068±0.007	21.32±0.04		0.064±0.009
15.78±0.003		0.064±0.007	21.52±0.04		0.081±0.010
16.20±0.03		0.075±0.008	21.72±0.04	100±30	0.048±0.007
16.71±0.03		0.069±0.008	21.78±0.04		0.106±0.011
16.86±0.03		0.034±0.005	22.19±0.04		0.077±0.009
16.99±0.03		0.070±0.008	22.48±0.04		0.081±0.010
17.07±0.03		0.050±0.007	22.62±0.05		0.062±0.008
17.47±0.03		0.103±0.011	22.76±0.05		0.073±0.009
17.82±0.04		0.132±0.014	22.82±0.05		0.127±0.013
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Energy	Γn	$g\Gamma_n\Gamma_{\gamma}/\Gamma$	Energy	Г	ar r /r
keV)	(eV)	(eV)	(keV)	n	al l \l
23.34±0.05		0.088±0.010	27.13±0.05	(eV)	(eV) 0.196±0.021
23.72±0.05		0.101±0.011	27.39±0.05		0.072±0.010
24.24±0.05	40±15	0.364±0.037	27.61±0.06		0.130±0.014
24.33±0.05		0.077±0.010	27.95±0.06		0.110±0.012
24.81±0.05		0.089±0.011	28.06±0.06		0.131±0.014
24.93±0.05		0.141±0.015	28.77±0.06		0.088±0.010
25.28±0.05		0.093±0.011	28.96±0.05		0.183±0.019
25.72±0.05		0.080±0.010	29.09±0.06		0.148±0.015
26.00±0.05		0.207±0.021	29.18±0.06		0.106±0.011
26.58±0.05 26.95±0.05		0.102±0.012	29.34±0.06		0.064±0.008
27.30±0.05		0.051±0.007	29.55±0.06		0.143±0.015
27.3010.05	į	0.105±0.012	29.74±0.06		0.134±0.014
			30.06±0.06		0.194±0.020

 $()^a$ Γ_n value assumed

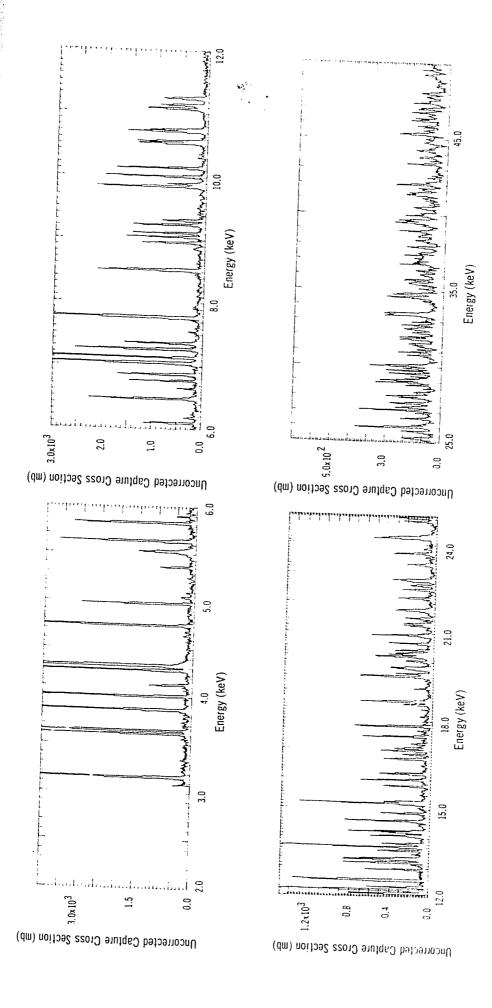


FIGURE 1. UNCORRECTED EXPERIMENTAL CAPTURE YIELD

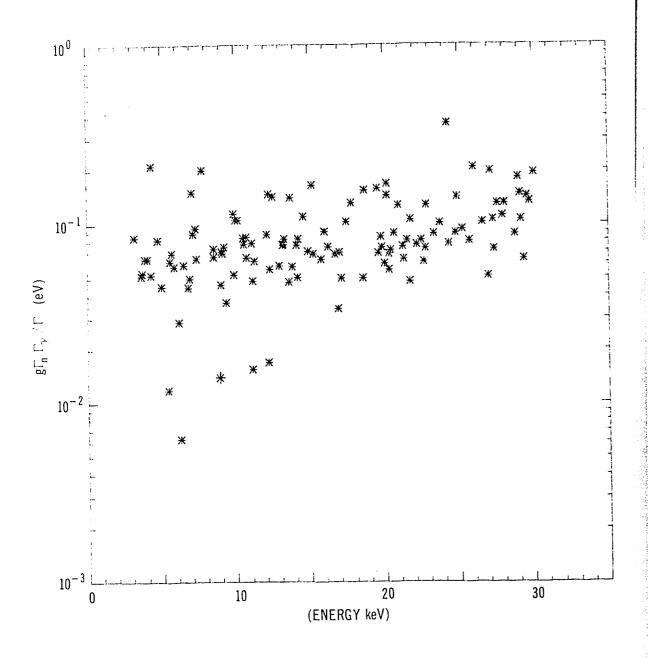


FIGURE 2. PLOT OF CAPTURE KERNEL VERSUS NEUTRON ENERGY

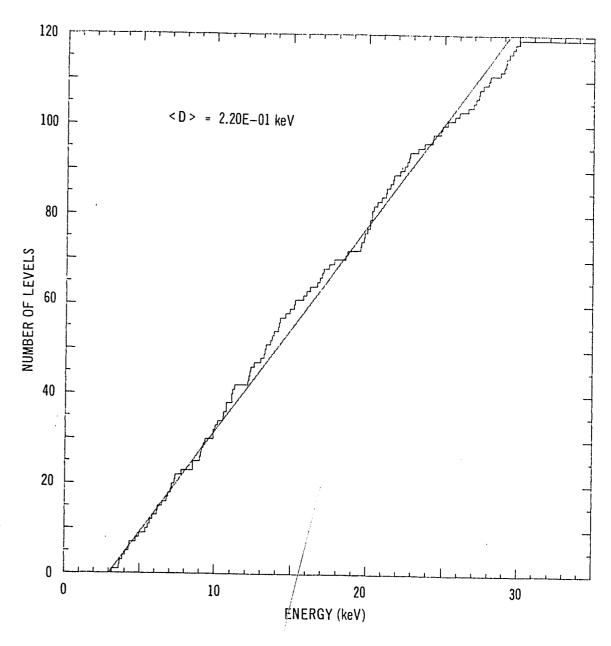


FIGURE 3. STAIRCASE PLOT OF NUMBER OF RESONANCES
BELOW ENERGY E

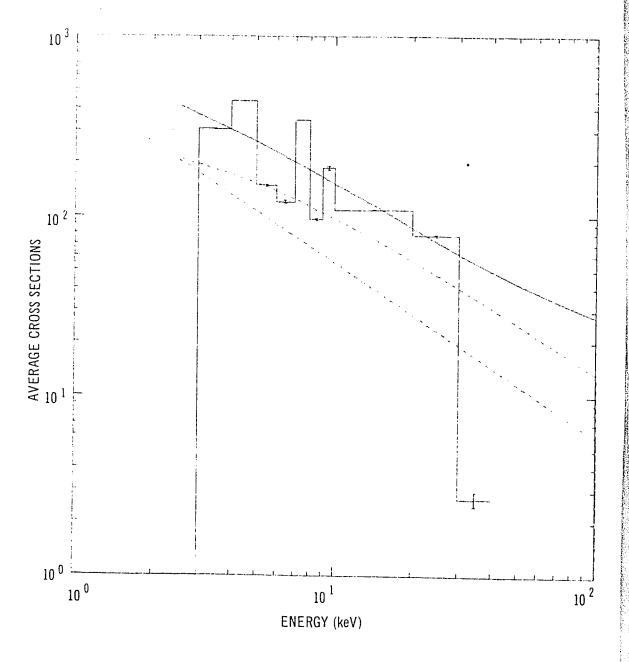


FIGURE 4. COMPARISON OF EXPERIMENTAL DATA WITH STATISTICAL MODEL CALCULATIONS OF CAPTURE CROSS SECTION ASSUMING $s_o=0.5\times10^{-4},\ s_1=5\times10^{-4},\ r_{\gamma\,s}=r_{\gamma\,p}=0.2\,\mathrm{eV}$ and $\langle D_{s.}\rangle=660\,\mathrm{eV}$

