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NUCLEAR PHYSICS RESEARCH IN THE USSR

Collected Abstracts

Issue 21

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I. Institute of Physics and Power Engineering

NEUTRON CROSS-SECTIONS OF DEUTERIUM IN THE ENERGY RANGE  
FROM 0.0001 eV TO 15 MeV

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Paper submitted to the scientific and technical publication  
"Yadernye konstanty" (Nuclear constants), No. 20

In the evaluations of the neutron cross-sections of deuterium that have been carried out abroad, significant discrepancies have been found between the results of various experiments designed to obtain data on the deuterium scattering cross-section in the region of low energies.

New experimental data that have recently appeared make it possible to resolve these contradictions and to select with significantly better justification the value of the total cross-section for deuterium in the low energy region.

In this context it was decided to re-evaluate the deuterium cross-sections for the SOKRATOR library. The present paper describes this evaluation and discusses the results obtained.

The bibliography used for the evaluation is listed in four tables with brief descriptions of the experimental methods.

References to sources of information not included in these tables are given in the list of references at the end of the paper.

At energies below 15 MeV only three interaction reactions of neutrons with deuterium are possible:

- (1) Elastic scattering;
- (2) Radiative capture;
- (3) The  $(n,2n)$  reaction.

The cross-sections for these reactions and also the total cross-section for the interaction of neutrons with deuterium nuclei are discussed in the paper, which also gives the corresponding recommended curves.

The evaluation results are presented in the format of the SOKRATOR library. A 26-group system of constants is likewise given containing a null group corresponding to 14-MeV neutrons.

NEUTRON CROSS-SECTIONS OF NATURAL ERBIUM AND ITS  
STABLE ISOTOPES

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Paper submitted to the scientific and technical publication  
"Yadernye konstanty" (Nuclear constants)

The paper consists of two parts. The first gives an evaluation of the energy dependence of the total cross-sections and of the cross-sections for capture and elastic and inelastic scattering and the elastic scattering anisotropy for natural erbium and its stable isotopes in the energy range 0.0253 eV-15 MeV. The exact behaviour of the cross-sections was first obtained by calculation: in the allowed resonance region on the basis of data on neutron resonance parameters, and in the forbidden resonance regions on the basis of statistical theory. The cross-sections obtained were then compared with experimental data on the energy dependence of cross-sections, on thermal cross-sections, and on resonance integrals of capture. A preliminary evaluation was made of the available data, and the results that seemed most reliable to us were selected.

The calculation in the allowed resonance region was done with the "Uran" program [1], in which the Breit-Wigner formula for an isolated resonance is used taking into account Doppler broadening of the resonance lines and interference between potential and resonance scattering. The calculation was performed to some energy  $E_{lim}$  below which no systematic gap in the levels is observed experimentally. The value of  $E_{lim}$  for each isotope was determined by analysis of the behaviour of the average distances observed between the levels  $\bar{D}_{obs}$  and the average reduced neutron widths  $\bar{G}_n^0$  depending on energy. The average values of the distributions of  $D_{obs}$  and  $G_n^0$  were compared with the theoretical distributions of Wigner and Porter-Thomas. The agreement between the observed and theoretical distributions in the energy range 0- $E_{lim}$  also served as evidence for the absence of a systematic gap of levels in the energy region below  $E_{lim}$ .

The recommended average resonance parameters ( $\bar{D}_{obs}$ ,  $\bar{G}_\gamma$ ,  $S^0$ ) were determined for the energy range 0- $E_{lim}$ .

Above  $E_{lim}$  the cross-sections were calculated with the PRS-2M program [2], in which the formula of the Hauser-Feschbach statistical theory is used with the penetration factors of the nuclear surface from the optical model and the formula for level density from the Fermi-gas model. In this program the single-particle potential contains a term to account for spin-orbital interaction. The centrosymmetric optical potential was taken in the Woods-Saxon form. The parameters of the optical potential are independent of neutron energy. The values of these parameters (except  $r_0$ ) were taken to correspond with the data of Ref. [3], in which the total cross-sections were calculated for a wide range of energies and mass numbers. The values for  $r_0$  were chosen in such a way that the value of the force function, determined in accordance with the optical model, was approximately equal to the value obtained from the allowed resonance region. The values of the average resonance parameters  $\bar{G}_\gamma$  and  $\bar{D}_{obs}$  were also obtained from the latter region.

The cross-sections for the natural mixture of erbium isotopes were found by summation of the isotope cross-sections, taking their percentages into account. Comparison of the cross-sections obtained in this way with data from direct measurement of natural erbium cross-sections constituted an additional verification of the recommended cross-sections.

The second part of the paper contains 26, 21 and 80 group cross-sections for natural erbium and its stable isotopes, plotted on the basis of the energy dependence obtained for the cross-sections. These group cross-sections are intended for reactor calculations. The methods of plotting and averaging the group constants are the same as in Refs. [4] and [5].

The evaluated cross-sections are presented in the format of the SOKRATOR library [6]. The files of these evaluations can be obtained from the Obninsk Nuclear Data Centre.

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II. I.V. Kurchatov Institute of Atomic Energy

MEASUREMENT OF EFFECTIVE CROSS-SECTION FOR CALIFORNIUM-252  
FISSION BY FAST REACTOR NEUTRONS

Eh.F. Fomushkin, E.K. Gutnikova, G.F. Novoselov, V.I. Panin

The effective cross-section for californium-252 fission by neutrons from a fast pulsed reactor were measured. For this purpose a device was set up in which a spring moved the glass plate detector relatively to the fixed layers of fissionable materials at the moment of the neutron pulse. Neptunium-237 was used as reference.

The cross-section for californium-252 fission by reactor neutrons obtained as a result of the measurements was  $1.58 \pm 0.14$  b; within a margin of error this agrees with the data of other authors.

The anomalously high value of the cross-section confirms the significant effect of the neutron shell  $N = 152$  on the fission characteristics of heavy nuclei. The lowering of fission barriers on passing through the shell  $N = 152$  observed in curium isotopes also apparently occurs in californium isotopes.

MEASUREMENT OF THE CROSS-SECTION OF PLUTONIUM-240 FISSION ON  
NUCLEAR EXPLOSION NEUTRONS

Eh.F. Fomushkin, E.K. Gutnikova, G.F. Novoselov, V.I. Panin

A dielectric polymer film was used as the fragment detector in measuring the fission cross-sections for nuclear explosion neutrons by the time-of-flight method. The time base of the detector film relative to the layer of fissionable material was provided mechanically by a reel revolving at a rate of  $\sim 10^4$  rev/min. The time resolution achieved in measuring the plutonium-240 fission cross-sections was 12 ns/m. The cross-section was measured by the relative method and the fission cross-section of plutonium-241 was used as reference. The films were chemically processed by the standard method. They were inspected, and the fission fragment tracks counted, with an optical microscope.

The plutonium-240 fission cross-section measured in the neutron energy range 2-200 keV has a resonance character. The results of averaging the cross-section by energy groups are given in the table.

$E_{\min}$	$E_{\max}$	keV	$\bar{\sigma}_f, b$	$\Delta\sigma_f/\sigma_f, \%$
2	4		0,080	6,8
4	6		0,124	7,0
6	8		0,124	6,8
8	10		0,083	7,2
10	15		0,142	9,3
15	20		0,080	7,4
20	30		0,130	9,4
30	50		0,220	7,1
50	100		0,123	7,5
100	200		0,103	7,4

THE LEVEL POPULATION FOR  $^{104}\text{Pd}$  AND  $^{106}\text{Pd}$  IN INELASTIC SCATTERING  
OF FAST NEUTRONS FROM A REACTOR

L.I. Govor, A.M. Demidov, M.R. Akhmed, Kh.I. Shakarchi,  
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The authors measured the gamma-ray spectra for fast neutrons inelastically scattered on  $^{104}\text{Pd}$  and  $^{106}\text{Pd}$  and evaluated the energies and relative intensities of 105  $\gamma$ -transitions in  $^{104}\text{Pd}$  (see Table 1) and 107  $\gamma$ -transitions in  $^{106}\text{Pd}$ . Schemes of the levels and  $\gamma$ -transitions of these nuclei were prepared.

A comparison was made between the experimental values for the level population in the  $(n, n'\gamma)$  reaction and values calculated with a statistical model. Agreement was found between the experimental and theoretical values for levels with the characteristics  $0^+$  and  $2^+$ , whereas there was a discrepancy of a factor of 2 for levels  $3^-$  and  $4^+$ .

Energy and intensity of  $^{104}\text{Pd}$   $\gamma$ -lines for inelastic scattering of fast neutrons from a reactor

$E_\gamma$ , keV	$\Delta E_\gamma$ , keV	$I_\gamma$ , rel.	$\Delta I_\gamma$	Place in scheme
215,6	0,2	0,7	0,2	
460,0	0,5	0,16	0,03	1793,8 - 1333,7
478,8	0,3	1,4	0,2	1820,8 - 1341,6
497,8	1,0	0,07	0,04	1820,8 - 1323,5
539,4	0,5	0,15	0,02	
555,74	0,07	100,0		555,7 - 0
617		0,32	0,10	1941,2 - 1323,5
623		0,14	0,04	2445,4 - 1820,8
727,4	0,5	0,16	0,04	
740,62	0,12	0,79	0,07	2032,2 - 1341,6
758,69	0,12	0,84	0,07	2032,2 - 1323,5
767,77	0,10	12,1	0,6	1323,5 - 555,7
777,93	0,10	4,2	0,2	1333,7 - 555,7
785,81	0,10	9,2	0,5	1341,6 - 555,7
818,5	0,9	0,059	0,015	
841,0	0,7	0,11	0,03	2181,6 - 1341,6
858,01	0,15	0,80	0,07	2181,6 - 1323,5
879,3	0,8	0,067	0,016	
889,4	0,6	0,12	0,04	
895,1	1,0	0,05	0,02	
902,4	0,6	0,13	0,03	2244,7 - 1341,6
910,2	0,4	0,22	0,04	2992,4 - 2032,2
922,5	0,8	0,13	0,05	2265,3 - 1341,6
925,9	0,2	0,65	0,09	2249,4 - 1323,5
934,8	0,9	0,062	0,015	2276,7 - 1341,6
941,78	0,15	0,74	0,06	2265,3 - 1323,5
953,9	0,7	0,076	0,016	3135,8 - 2181,6
974,4	0,2	0,40	0,05	
978,3	0,4	0,22	0,03	2799,1 - 1820,8
996,0	0,7	0,10	0,03	2338,3 - 1341,6
1028,1	0,9	0,09	0,03	2351,6 - 1323,5

Continuation of table

$E_{\gamma}$ , keV	$\Delta E_{\gamma}$ , keV	$I_{\gamma}$ , rel.	$\Delta I_{\gamma}$	Place in scheme
1077,2	1,0	0,08	0,03	
1088,1	0,5	0,14	0,03	
1121,6	0,6	0,30	0,10	
1132,1	0,3	0,38	0,04	2465,8 - 1333,7
1167,56	0,13	1,00	0,07	
1179,3	0,2	0,49	0,05	2521,1 - 1341,6
1230,7	0,2	0,41	0,08	2572,5 - 1341,6
1238,05	0,09	4,9	0,3	1793,8 - 555,7
				1792,9 - 555,7
1247,4	0,5	0,24	0,05	2570,9 - 1323,5
1265,09	0,1	4,0	0,3	1820,8 - 555,7
1271,69	0,16	0,59	0,05	2613,3 - 1341,6
1284,1	0,4	0,18	0,03	3104,9 - 1820,8
1360,0	0,6	0,10	0,03	2641,7 - 1341,6
1318,2	0,3	0,30	0,03	2641,7 - 1323,5
1341,73	0,09	6,0	0,4	1341,6 - 0
1372,6	0,9	0,09	0,03	2714,5 - 1341,6
1381,4	0,8	0,10	0,03	2714,5 - 1333,7
1390,1	1,1	0,07	0,03	
1397	line boundary	0,26	0,03	
1409	line boundary	0,26	0,03	
1450,5	0,4	0,19	0,03	2774,0 - 1323,5
1480,4	1,0	0,08	0,03	
1514,9	1,0	0,08	0,03	
1525,4	0,2	0,60	0,05	2082,2 - 555,7
1542	line boundary	0,16	0,03	2866 - 1323,5
1551,0	0,7	0,11	0,03	2874,5 - 1323,5
1563,0	0,2	0,77	0,07	2138,7 - 555,7
1599,9	0,4	0,18	0,03	2923,4 - 1323,5
1612,6	1,4	0,04	0,02	
1622,8	0,2	1,36	0,14	2178,5 - 555,7
1615,2	0,4	0,31	0,06	2181,6 - 555,7
1631,10	0,14	3,2	0,3	2193,2 - 555,7
1652,6	0,8	0,15	0,03	3000,3 - 1323,5

Continuation of table

$E_{\gamma}$ , keV	$\Delta E_{\gamma}$ , keV	$I_{\gamma}$ , rel.	$\Delta I_{\gamma}$	Place in scheme
1688,92	0,14	1,38	0,14	2244,7 - 555,7
1696,8	0,6	0,13	0,03	3020,3 - 1323,5
1720,7	0,3	0,35	0,04	2276,7 - 555,7
1762	line boundary	0,18	0,03	3086 - 1323,5
1782	line boundary	0,59	0,10	2338,3 - 555,7
1794	line boundary	0,48	0,06	1793,8 - 0
				2351,6 - 555,7
1812,5	0,7	0,12	0,03	3135,8 - 1323,5
1839,7	0,8	0,1	0,03	
1895,8	0,6	0,13	0,03	
1900,8	0,2	0,55	0,06	2456,6 - 555,7
1909,8	1,0	0,03	0,03	2465,8 - 555,7
1926,2	0,5	0,16	0,03	2521,1 - 555,7
1977,4	0,3	0,52	0,06	2533,1 - 555,7
2016,9	0,3	0,33	0,04	2572,5 - 555,7
2070,0	0,3	0,43	0,05	2625,7 - 555,7
2106,5	0,7	0,14	0,03	
2138,7	0,3	0,35	0,04	2694,4 - 555,7
2164,7	1,4	0,05	0,02	
2210,6	0,3	0,43	0,05	2766,3 - 555,7
2232,8	1,1	0,07	0,02	
2276,9	0,3	0,46	0,05	2276,7 - 0
2338,3	0,3	0,31	0,04	2338,3 - 0
2362,4	0,4	0,31	0,06	2917,9 - 555,7
2420	line boundary	0,21	0,03	2975 - 555,7
2452,2	0,7	0,14	0,03	3007,7 - 555,7
2523,6	0,7	0,12	0,03	3079,3 - 555,7
2534,4	0,7	0,14	0,03	
2540,4	0,6	0,15	0,03	3096,1 - 555,7
2593,1	0,6	0,15	0,04	
2598,8	1,0	0,10	0,04	
2623,6	0,4	0,54	0,06	3179,5 - 555,7
2705,5	0,7	0,13	0,03	
2715,8	0,6	0,17	0,04	
2726,1	1,1	0,09	0,03	3281,5 - 555,7
2916,2	1,0	0,11	0,03	
3001,6	1,2	0,03	0,03	3000,3 - 0
3007,6	1,0	0,1	0,03	3000,7 - 0
3013,5	0,9	0,12	0,03	
3032,0	0,8	0,14	0,04	
3281,3	1,1	0,07	0,03	3281,5 - 0

LEVEL POPULATION FOR  $^{92}\text{Mo}$  IN INELASTIC SCATTERING OF FAST NEUTRONS FROM A REACTOR

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Paper presented at Third Conference on Neutron Physics

Table 1

Energies and intensities of  $\gamma$ -lines of  $^{92}\text{Mo}$  from the  $(n, n'\gamma)$  reaction for fast neutrons from a reactor

$E_{\gamma}$ , keV	$I_{\gamma}$ , rel. $\theta = 90^\circ$	Place in scheme
146,6 $\pm$ 0,6	0,7 $\pm$ 0,2	2759,4 - 2610,8
243,6 $\pm$ 0,3	17,3 $\pm$ 1,0	2525,0 - 2281,6
304,0 $\pm$ 0,5	1,5 $\pm$ 0,2	3366,4 - 3062,2
329,2 $\pm$ 0,7	1,3 $\pm$ 0,2	2610,8 - 2281,6
361,4 $\pm$ 0,8	1,0 $\pm$ 0,2	3366,4 - 3004,8
479,8 $\pm$ 0,2	4,0 $\pm$ 0,6	3004,8 - 2525,0
537,3 $\pm$ 0,2	5,9 $\pm$ 0,5	3062,2 - 2525,0
772,0 $\pm$ 0,2	25,9 $\pm$ 1,5	2281,6 - 1509,6
941,7 $\pm$ 0,6	0,55 $\pm$ 0,10	3944,8 - 3004,8
946,3 $\pm$ 0,8	0,6 $\pm$ 0,2	4008,5 - 3062,2
1009,4 $\pm$ 0,4	2,2 $\pm$ 0,4	2519,0 - 1509,6
1340,0 $\pm$ 0,2	6,5 $\pm$ 0,5	2849,6 - 1509,6
1509,0 $\pm$ 0,1	100	1509,6 - 0,0
1581,8 $\pm$ 0,7	0,80 $\pm$ 0,15	3091,4 - 1509,6
2033,0 $\pm$ 0,4	2,2 $\pm$ 0,3	3542,7 - 1509,6
2113,0 $\pm$ 0,3	1,6 $\pm$ 0,3	3622,6 - 1509,6
2179,3 $\pm$ 0,8	0,6 $\pm$ 0,2	3683,9 - 1509,6
2417,6 $\pm$ 1,5	0,3 $\pm$ 0,1	3927,2 - 1509,6
3091,4 $\pm$ 0,5	2,50 $\pm$ 0,15	3091,4 - 0,0
3927,2 $\pm$ 1,0	0,35 $\pm$ 0,15	3927,2 - 0,0
3944,8 $\pm$ 1,0	0,40 $\pm$ 0,15	3944,8 - 0,0
4022,1 $\pm$ 0,7	0,75 $\pm$ 0,25	

Table 2

Experimental and theoretical level populations for  $^{92}\text{Mo}$  in the  $(n, n'\gamma)$  reaction for fast neutrons from a reactor

$E_{\text{lev}}$ , keV	$I^{\pi}$	$P_s$ rel. (exp)	$P_s^M$ (theor.), rel.
1509,6	$2^+$	73,0 $\pm$ 4,0	73,0
2281,6	$4^+$	9,5 $\pm$ 2,2	11,9
2519,0	$0^+$	2,2 $\pm$ 0,4	3,8
2525,0	$5^-$	6,7 $\pm$ 1,5	2,5
2610,8	$6^+$	0,9 $\pm$ 0,5	0,54
2759,4	$8^+$	1,0 $\pm$ 0,3	0,27
2849,6	$3^-$	5,9 $\pm$ 0,5	5,2
3004,8	$4^+$	2,4 $\pm$ 0,8	2,7
3062,2	$4^+$	3,6 $\pm$ 0,5	2,7
3091,4	$2^+$	3,9 $\pm$ 0,4	4,1

The  $P_s$  values are given after deduction of the cascade population.

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UPPER LIMIT OF CROSS-SECTION OF  $^{226}\text{Ra}$  FISSION BY THERMAL NEUTRONS

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The aim of the work was to verify the findings published in Ref. [1], which reported the comparatively large cross-section of 4.7 mb for fission of  $^{226}\text{Ra}$  by thermal neutrons, and a symmetrical distribution of the fragment masses. Using mica detectors and glass detectors with an aluminium filter it was found by the authors of the present paper that in a flux of thermal neutrons the upper limit for the radium fission cross-section is 0.05 mb. This value is in agreement with that given earlier in Ref. [2]. The spectrum of diameters of fragment tracks in the glass detector, analogous to the spectrum of kinetic energy, was found to be asymmetrical and was probably due to the fission of uranium impurities in the radium target. Thus the results of Ref. [1] were not confirmed in this experiment.

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## SCATTERING OF 4.7-MeV NEUTRONS ON Al AND Fe

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By now much material has been collected on inelastic scattering of neutrons with an initial energy  $E_n = 14$  MeV, but little work has been done on neutrons with lower initial energies.

This paper presents the results of measurements of neutron scattering spectra on Al and Fe when  $E_n = 4.7$  MeV at three scattering angles:  $20^\circ$ ,  $97^\circ$  and  $120^\circ$ .

We used the time-of-flight method based on a pulsed operating regime of the EhG-5 accelerator, which had the following parameters: pulse repetition frequency 2 MHz, average current on target  $1 \mu\text{A}$ , resolution in elastic scattering peak 3-4 ns, distance scatterer - target 10 cm, scatterer - detector 250-310 cm. The scatterers were aluminium and iron cylinders with the dimensions  $l = 4$  cm and  $d = 5$  cm. The detector was a plastic scintillator of  $d = 10$  cm and  $h = 3$  cm coupled with an FEhU-63 photomultiplier. The detector was placed in a mobile shield made of paraffin with boron carbide and was also surrounded with a 5-cm layer of lead. The overall dimensions of the shield were  $80 \times 120$  cm. A standard titanium-deuterium target with a thickness of  $1 \text{ mg/cm}^2$  was used.

The energy calibration of the device was performed using  $^{252}\text{Cf}$  spontaneous fission neutrons, the spectrum of which was taken in the form:  $N(E) = N_0 e^{-0.88E} \text{sh} \sqrt{2E}$  [3]. As detector of the fission fragments giving zero time we used a scintillating film with a thickness of  $1 \text{ mg/cm}^2$  together with an FEhU-36 photomultiplier, in front of which the  $^{252}\text{Cf}$  source was placed at a distance of 0.3 mm. The calibration was performed over the neutron energy range 0.7-6 MeV with an accuracy of 3-4%, which was taken into account in the final result. After calibration we measured the spectra of neutrons with  $E_n = 4.7$  MeV which were scattered on Fe and Al at the three angles (laboratory system of co-ordinates):  $20^\circ$ ,  $97^\circ$  and  $120^\circ$ . After these measurements had been completed the calibration was checked. The spectra of neutrons scattered at an angle of  $120^\circ$  are given in the figure.

In determining the inelastic scattering cross-sections the cross-sections for elastic interaction were presumed to be known, their values being taken from Ref. [4]. The results of the cross-section measurements for the identified levels are given in the table.



$^{27}\text{Al}$

Scattering cross-section (mb/sr)			
Scattering angle	Level 1.01 MeV	Level 2.2 MeV	Level 3 MeV
$20^\circ$	$8,3 \pm 0,5$	$4,3 \pm 0,3$	$4,5 \pm 0,4$
$97^\circ$	$9,4 \pm 0,6$	$3,9 \pm 0,3$	$4 \pm 0,3$
$120^\circ$	$12 \pm 0,8$	$4,2 \pm 0,3$	$3,5 \pm 0,3$

Fe

Scattering cross-section (mb/sr)			
Scattering angle	Level 0.84 MeV	Level 2.65 MeV	Total cross-section for inelastic scattering with $E_n = 1-4$ MeV
$20^\circ$	$15 \pm 0,9$	$9 \pm 0,5$	$60 \pm 3,5$
$97^\circ$	$13 \pm 0,8$	$11 \pm 0,7$	$52 \pm 3$
$120^\circ$	$14 \pm 0,9$	$8,8 \pm 0,5$	$54 \pm 3$

Thus the scattering is practically spherically symmetrical for the levels identified. It is striking that the spectra of neutrons scattered on  $^{27}\text{Al}$  lack a noticeable level with an energy of 2.73 MeV. However, in a study of inelastic scattering of protons in the energy range  $E_p = 5-6.5$  MeV [5] this level appeared clearly; the same study also revealed a significant energy dependence of proton inelastic scattering. Possibly the absence of this level in our results is due to a strong energy dependence of the formation cross-section of this level.

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IV. Atomic Reactor Research Institute

TOTAL NEUTRON CROSS-SECTION AND NEUTRON RESONANCE PARAMETERS OF  $^{241}\text{Am}$   
IN THE ENERGY RANGE 0.004-30 eV

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(Atomic energy)

On the ITEhF heavy-water reactor the transmissions in the energy range 0.004-8 eV were measured by the time-of-flight method with a resolution of 140 ns/m. The neutron pulse was formed by a selector chopper with synchronously revolving rotors suspended in a magnetic field [2]. Analogous measurements were made on our Institute's SM-2 reactor up to an energy of 30 eV. The experimental conditions and procedure for processing the results are described in detail in Ref. [1]. The present paper generalizes the results of the experiments mentioned above.

The  $\text{AmO}_2$  powder samples with a thickness of  $0.63 \times 10^{21}$  and  $3.3 \times 10^{21}$  at/cm<sup>2</sup> contained 99.99% of the isotope  $^{241}\text{Am}$ . The resonance parameters for 13 levels up to 8 eV were calculated by the form method using the single-level Breit-Wigner formula; two of these - 2.538 and 6.65 eV - were discovered for the first time. In order to make the calculated transmission values agree with the experimental values in the thermal region a negative level  $E_0 = -0.425$  eV was introduced. The existence of five powerful levels at  $E_0 = 2.36, 4.40, 6.78, 7.97$  and 16.02 eV was not confirmed. Analysis of the square root distribution of the neutron widths given showed that it corresponded to a Porter-Thomas distribution with a degree of freedom  $\nu = 1.07 \pm 0.23$ . The best agreement of this distribution with the experimental histogram was obtained for a level number  $n = 44$  and  $2g\Gamma_n = 0.096$  MeV. It follows from this that the loss of six weak levels below 26 eV must be postulated. Wigner distributions for one system of levels describe the histogram of the experimental distribution better, although two level systems must exist in the  $^{241}\text{Am}$  nucleus. The experimental data were analysed for correlation of a long-range order, using the  $\Delta_3$  test of Dyson and Metha [3] for the orthogonal assembly. The experiments gave  $\Delta_3 = 0.367$ ; the theoretical value for

two level systems is  $\Delta_3 = 0.723 \pm 0.219$ , for one system it is  $\Delta_3 = 0.362 \pm 0.11$ . By analysis of the statistical properties of the nucleus an average level distance  $\bar{D}_0 = 0.67 \pm 0.1$  eV,  $\overline{2g\Gamma_n} = 0.102$  MeV,  $S_0 = (0.76 \pm 0.18) \times 10^{-4}$  was obtained. The numerical data on the total cross-section were transmitted to the Nuclear Data Centre, and the resonance parameters were published in Ref. [1].

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Parameters of  $^{241}\text{Am}$  neutron resonances

$E_0$ , eV	$\Gamma$ , MeV	$2g\Gamma_n$ , MeV
-0,425	40	$2g\Gamma_n^0 = 1,0$
$0,306 \pm 0,002$	$45 \pm 1$	$0,0556 \pm 0,0004$
$0,573 \pm 0,004$	$43 \pm 1$	$0,0928 \pm 0,0016$
$1,283 \pm 0,004$	$41 \pm 2$	$0,330 \pm 0,016$
$1,916 \pm 0,005$	$46 \pm 2$	$0,107 \pm 0,002$
$2,533 \pm 0,008$	$41 \pm 2$	$0,070 \pm 0,001$
$2,581 \pm 0,009$	$38 \pm 2$	$0,150 \pm 0,004$
$3,956 \pm 0,009$	$28 \pm 3$	$0,230 \pm 0,008$
$4,947 \pm 0,010$	$31 \pm 5$	$0,176 \pm 0,005$
$5,390 \pm 0,012$	$38 \pm 7$	$0,844 \pm 0,114$
$6,100 \pm 0,013$	$42 \pm 14$	$0,116 \pm 0,005$
$6,650 \pm 0,015$		$0,05 \pm 0,03$
$7,53 \pm 0,02$		$0,07 \pm 0,04$
$8,17 \pm 0,02$	$42 \pm 5$	$0,096 \pm 0,004$
$9,11 \pm 0,02$	$48 \pm 3$	$0,358 \pm 0,006$
$9,84 \pm 0,03$	$48 \pm 3$	$0,370 \pm 0,007$
$10,11 \pm 0,03$		$0,025 \pm 0,004$
$10,39 \pm 0,03$	$45 \pm 4$	$0,294 \pm 0,007$
$10,99 \pm 0,04$	$52 \pm 4$	$0,0382 \pm 0,003$
$11,58 \pm 0,05$		$0,018 \pm 0,003$
$12,06 \pm 0,06$		$0,007 \pm 0,003$
$12,86 \pm 0,06$	$44 \pm 5$	$0,116 \pm 0,009$
$14,32 \pm 0,06$		$0,066 \pm 0,012$
$14,66 \pm 0,07$	$44 \pm 5$	$2,30 \pm 0,13$
$15,66 \pm 0,07$	$32 \pm 12$	$0,215 \pm 0,012$
$16,35 \pm 0,07$	$44 \pm 5$	$1,185 \pm 0,033$
$16,81 \pm 0,07$	$31 \pm 8$	$0,575 \pm 0,020$
$17,69 \pm 0,07$	$40 \pm 10$	$0,373 \pm 0,016$
$18,09 \pm 0,07$		
$19,39 \pm 0,07$	$37 \pm 12$	$0,182 \pm 0,016$
$20,28 \pm 0,07$		$0,050 \pm 0,010$
$20,84 \pm 0,09$		$0,064 \pm 0,011$
$21,72 \pm 0,09$		$0,067 \pm 0,012$
$22,74 \pm 0,09$		$0,070 \pm 0,012$
$23,08 \pm 0,09$		$0,39 \pm 0,05$
$23,33 \pm 0,09$		$0,40 \pm 0,05$
$24,17 \pm 0,09$		$1,27 \pm 0,08$
$25,05 \pm 0,10$		
$25,60 \pm 0,10$		$1,21 \pm 0,08$
$26,50 \pm 0,10$		
$26,67 \pm 0,10$		
$27,52 \pm 0,10$		
$27,65 \pm 0,10$		
$28,31 \pm 0,11$		$0,40$
$28,82 \pm 0,12$		$0,35$
$29,43 \pm 0,12$		$0,61$

THE TOTAL NEUTRON CROSS-SECTION AND NEUTRON RESONANCE PARAMETERS  
OF  $^{243}\text{Am}$  IN THE ENERGY RANGE 0.4-35 eV

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(Atomic energy)

On the SM-2 reactor the  $^{243}\text{Am}$  transmissions for neutrons with energies between 0.4 and 30 eV were measured by the time-of-flight method with a path length of 92 m and a resolution of 70 ns/m. The neutron burst was formed by a spectrum with three synchronously turning rotors suspended in a magnetic field. The samples of  $\text{AmO}_2$  powder contained 79.8% americium. The isotopic content of the samples was 96.9%  $^{243}\text{Am}$ , 3.32%  $^{241}\text{Am}$  and 0.08%  $^{244}\text{Cm}$ . The samples were placed in an aluminium cylinder with a wall thickness of 1 mm and annealed in an oxygen atmosphere at 400°C for 3 hours. The sample thicknesses were 0.8 mm ( $0.45 \times 10^{21}$  at/cm<sup>2</sup>) and 14 mm ( $7.9 \times 10^{21}$  at/cm<sup>2</sup>).

A battery of helium counters was used as a neutron detector. A statistical accuracy of 0.5-1.5% was maintained in the measurements. The error also contained a component connected with the determination of the shape of the resolution function, and the neutron background varied between 0.7 and 4%. The measured transmissions were corrected for scattering on the oxygen (up to 3% in a thick sample) and for the  $^{241}\text{Am}$  contribution.

The results were processed on a BEhSM-6 computer. The numerical material on  $\sigma_t(E)$  was transmitted to the Nuclear Data Centre. The resonance parameters for 48 levels were obtained using the single-level Breit-Wigner formula and the form method in the region below 26 eV and the area method above 26 eV. The measurements did not confirm the existence of a weak level with an energy of 22.011 eV, but a weak level with  $E_0 = 31.49$  eV was found. The value of the total resonance integral was calculated as 1740 b. The analysis of the statistical properties of the  $^{243}\text{Am}$  nucleus yielded average distances between the levels of  $\bar{D}_0 = 0.71 \pm 0.1$  eV and a force function  $S_0 = (0.89 \pm 0.19) \times 10^{-4}$ . The potential scattering cross-section could not be determined owing to a 20% impurity of reactive elements in the sample.

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Neutron resonance parameters for  $^{243}\text{Am}$

$E, \text{ eV}$	$G, \text{ MeV}$	$2g_n^2, \text{ MeV}$	$\sigma_0, \text{ b}$
$0,416 \pm 0,002$	$39 \pm 2$	$0,00089 \pm 0,00005$	$67 \pm 2$
$0,977 \pm 0,004$	$37 \pm 2$	$0,0134 \pm 0,003$	$460 \pm 20$
$1,355 \pm 0,0043$	$56 \pm 1$	$0,890 \pm 0,007$	$15130 \pm 300$
$1,743 \pm 0,005$	$39 \pm 1$	$0,208 \pm 0,002$	$3980 \pm 50$
$3,133 \pm 0,008$	$47 \pm 3$	$0,012 \pm 0,003$	$100 \pm 10$
$3,424 \pm 0,017$	$45 \pm 2$	$0,253 \pm 0,008$	$2130 \pm 160$
$3,844 \pm 0,017$	$22 \pm 5$	$0,009 \pm 0,001$	$137 \pm 30$
$5,120 \pm 0,03$	$63 \pm 2$	$0,260 \pm 0,006$	$1040 \pm 45$
$6,561 \pm 0,04$	$50 \pm 3$	$0,794 \pm 0,044$	$3100 \pm 380$
$7,063 \pm 0,04$	$46 \pm 3$	$0,072 \pm 0,011$	$286 \pm 14$
$7,856 \pm 0,05$	$36 \pm 9$	$1,580 \pm 0,136$	$7330 \pm 1800$
$8,39 \pm 0,02$	40	0,010	40
$8,77 \pm 0,02$	46	0,113	360
$9,32 \pm 0,02$	43	0,133	430
$10,31 \pm 0,03$	47	0,433	1150
$10,87 \pm 0,04$	40	0,013	40
$11,27 \pm 0,05$	49	0,267	630
$11,68 \pm 0,06$	35	0,094	290
$12,12 \pm 0,06$	41	0,152	400
$12,87 \pm 0,06$	43	2,20	5060
$13,15 \pm 0,06$	45	1,00	2200
$15,12 \pm 0,07$	33	0,070	180
$15,39 \pm 0,07$	37	0,36	3100
$16,20 \pm 0,07$	39	0,512	1060
$16,56 \pm 0,07$	27	0,174	510
$17,84 \pm 0,07$	35	0,210	449
$18,14 \pm 0,07$	27	0,046	120
$19,50 \pm 0,07$	27	0,193	470
$19,88 \pm 0,07$	40	0,085	140
$20,94 \pm 0,08$	29	0,54	1200
$21,09 \pm 0,08$	16	0,86	3200
$21,85 \pm 0,09$	27	0,14	300
$22,011 \pm 0,09$			
$22,59 \pm 0,09$	33	1,00	1700
$22,72 \pm 0,09$	19	0,65	2200
$24,39 \pm 0,09$	(22)	0,73	2400
$25,38 \pm 0,10$	(40)	0,14	180
$26,30 \pm 0,10$		0,06	80
$28,75 \pm 0,10$	(30)	1,16	1900
$27,34 \pm 0,10$	(30)	0,43	500
$28,73 \pm 0,10$	(30)	0,97	1145
$29,29 \pm 0,11$	(30)	0,68	1000
$30,12 \pm 0,12$	(30)	0,486	696
$31,06 \pm 0,12$	(30)	0,7	968
$31,49 \pm 0,13$	(30)	0,12	161
$32,43 \pm 0,14$	(30)	0,88	1140
$33,19 \pm 0,14$	(30)	1,9	2350
$33,92 \pm 0,15$	(30)	0,8	986

NEUTRON RESONANCES OF THE ISOTOPES  $^{244}\text{Cm}$ ,  $^{245}\text{Cm}$ ,  $^{246}\text{Cm}$  AND  $^{248}\text{Cm}$

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Neutron choppers with synchronously revolving rotors suspended in a magnetic field were installed at the ITEhF and NIIAR reactors for measuring the transmission of samples containing the curium isotopes 244, 245, 246 and 248. The resolution of the spectrometers was 70 ns/m on the 92-m base at NIIAR [1] and 140 ns/m on the 50-m base at ITEhF [2].

A dehydrated powder of stable curium oxide with known oxygen content ( $\text{Cm}_2\text{O}_3$ ) was studied. Before being filled into the target the powder was annealed at a temperature of 900-1100°C. Two samples were prepared with different curium isotope contents. The sample characteristics are given in Table 1. Sample 1 contained the impurities  $^{243}\text{Am}$  (2%) and  $^{240}\text{Pu}$  (1.6%); sample 2 contained 0.13%  $^{243}\text{Am}$  and 0.8%  $^{240}\text{Pu}$ .

The transmissions of the samples were measured in the energy region above 0.5 eV; a statistical accuracy within 1-2% was maintained, and the neutron background varied between 0.5 and 2%. Correction was made for the scattering of neutrons on oxygen.

The resonance parameters were calculated by the form and area method using the single-level Breit-Wigner formula. The resolution function was not approximated by an analytical expression, but directly calculated by computer, taking into account the factors that characterize the resolving power of the device. The location of the resonances  $E_0$ , and the values of the neutron ( $2g\Gamma_n$ ) and total ( $\Gamma$ ) widths were determined.

1. Curium-244: Measurements were carried out for several  $^{244}\text{Cm}$  sample thicknesses, the greatest being  $n_0 = 0.24 \times 10^{22}$  atoms/cm<sup>2</sup>. Table 2 gives the resonance parameters for  $^{244}\text{Cm}$  up to 171 eV. In processing the resonances by the area method the radiation width  $\Gamma_\gamma$  was taken to be 37 MeV.

The average distance between levels  $\bar{D} = 14.6 \pm 2.2$  eV and the force function  $S_0 = (0.68 \pm 0.30) \times 10^{-4}$  were calculated. It was also shown



that for  $^{244}\text{Cm}$  the neutron widths quoted followed the Porter-Thomas distribution for one degree of freedom, while the distances between levels followed the Wigner distribution.

2. Curium-245: Table 3 gives tentative values of the resonance parameters for the isotope  $^{245}\text{Cm}$ . A sample with thickness  $0.26 \times 10^{21}$  atoms/cm<sup>2</sup> was used. Up to an energy of 50 eV 32 levels were found. They were all calculated by the area method, it being taken that  $\Gamma_{\gamma} = 40$  MeV.

3. Curium-246: For a  $^{246}\text{Cm}$  sample  $0.21 \times 10^{22}$  atoms/cm<sup>2</sup> thick 6 levels were identified in the energy range up to 157 eV. The resonance parameters were calculated by the form and area method, taking  $\Gamma_{\gamma} = 37$  MeV. In the present paper the level 26.88 eV is not identified unambiguously, but attributed to the isotope  $^{246}\text{Cm}$ . It may also be ascribed to  $^{248}\text{Cm}$ . The parameters calculated for it in this case are  $\Gamma_n = 36.7 \pm 3.1$  MeV,  $\Gamma_n = 21.7$  MeV; the neutron width agrees with the  $\Gamma_n = 19.6 \pm 0.9$  MeV of Ref. [3], in which this level is assigned to  $^{248}\text{Cm}$ .

4. Curium-248: The resonance parameters of the levels of  $^{248}\text{Cm}$  for energies up to 100 eV are given in Table 5. The sample thickness was  $0.27 \times 10^{21}$  atoms/cm<sup>2</sup>. A value  $\Gamma_{\gamma} = 40$  MeV was chosen in processing the resonances by the area method.

A new resonance with an energy of 84 eV is given for  $^{248}\text{Cm}$ . The level with an energy of 35.0 eV [3] was not found. The neutron width  $\Gamma_n = 3.5 \pm 0.3$  MeV of the 35.0 eV resonance in  $^{244}\text{Cm}$  [3] does not permit the conclusion that  $^{248}\text{Cm}$  has a level of the same energy; it should apparently be ascribed to  $^{244}\text{Cm}$ . In measurements taken in an atomic explosion [4], no resonance with an energy of 35.0 eV was found for  $^{248}\text{Cm}$ .

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Table 1

Sample characteristics

Number of sample	Weight, mg	Isotopes (%)					
		<sup>242</sup> Cm	<sup>244</sup> Cm	<sup>245</sup> Cm	<sup>246</sup> Cm	<sup>247</sup> Cm	<sup>248</sup> Cm
1	82,8	0,31	88,62	9,57	1,49	-	-
2	116,4	-	39,28	0,47	51,85	1,61	6,79

Table 2

Resonance parameters of <sup>244</sup>Cm

E <sub>0</sub> , eV	Γ, MeV	Γ <sub>n</sub> , MeV	E <sub>0</sub> , eV	Γ <sub>n</sub> , MeV
7,67	44 ± 3	10,4 ± 0,4	95,6	26,0 ± 4,8
16,77	37 ± 5	1,90 ± 0,30	95,5	7,8 ± 2,2
22,85	36 ± 10	0,84 ± 0,10	132	16 ± 8
35,0	33 ± 5	5,1 ± 0,8	139	2,2 ± 0,9
52,8	-	0,56 ± 0,15	171	3,6 ± 1,8
69,8	-	0,44 ± 0,25		

Table 3

Resonance parameters of <sup>245</sup>Cm

E <sub>0</sub> , eV	2Γ <sub>n</sub> , MeV	E <sub>0</sub> , eV	2Γ <sub>n</sub> , MeV	E <sub>0</sub> , eV	2Γ <sub>n</sub> , MeV
1,93	-	25,0	2,4 ± 0,4	36,3	3,6 ± 1,8
4,69	1,73 ± 0,35	26,9	-	40,9	1,8 ± 0,9
9,25	0,32 ± 0,05	27,1	1,0 ± 0,3	42,9	3,0 ± 1,5
11,4	0,50 ± 0,20	29,6	4,2 ± 0,7	43,5	-
14,0	0,25 ± 0,08	31,4	0,5 ± 0,2	44,9	1,7 ± 0,5
15,9	-	32,4	0,4 ± 0,2	47,8	5,7 ± 1,4
21,6	2,6 ± 0,4	35,2	-	49,2	2,2 ± 1,2
				50,5	1,8 ± 0,7

Table 4

Resonance parameters of <sup>246</sup>Cm

E <sub>0</sub> , eV	Γ, MeV	Γ <sub>n</sub> , MeV	E <sub>0</sub> , eV	Γ <sub>n</sub> , MeV
4,32	27 ± 2	0,34 ± 0,01	84,5	-
15,29	28 ± 3	0,52 ± 0,01	91,5	9,9 ± 2,5
26,88*	26 ± 3	2,85 ± 0,04	157	34,1 ± 7,6

\* Possibly <sup>248</sup>Cm.

Table 5

Resonance parameters of <sup>248</sup>Cm

E <sub>0</sub> , eV	Γ, MeV	Γ <sub>n</sub> , MeV	E <sub>0</sub> , eV	Γ <sub>n</sub> , MeV
7,26	36 ± 3	1,90 ± 0,04	84	-
75,6	-	102,5 ± 13,6	98,6	169 ± 18

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MEASUREMENT OF THE ENERGY DEPENDENCE OF  
 $\eta$  FOR  $^{233}\text{U}$  IN THE REGION 0.02-1 eV

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Exact data on the effective number of fission neutrons  $\eta$  for  $^{233}\text{U}$  in the low-energy region are important for planning slow-neutron breeder reactors, as the fundamental indeterminacy of the breeding ratio in such reactors is due to the indeterminacy of  $\eta$  for  $^{233}\text{U}$ . For this reason the energy dependence of  $\eta$  for  $^{233}\text{U}$  in the range 0.02-1 eV was determined, with an accuracy of 1-2%, on the Institute's WWR-M reactor. The measurements were made by the time-of-flight method with a resolution of  $\sim 12 \mu\text{s/m}$  and normalized to the value at  $E = 0.0253 \text{ eV}$ , it being taken that at this point  $\eta = 2.297$ . The neutron beam from the reactor was transformed into a pulsed beam by a mechanical chopper, the rotor of which was 300 mm in diameter and had a slit width of 2 mm. The sample and detectors were located at a distance  $L = 503 \pm 5 \text{ cm}$  from the chopper. The fission neutrons were recorded in a geometry close to  $2\pi$  by a battery of 50 SNM-37 helium counters with a moderator. In order to reduce the neutron lifetime a cadmium screen was placed between the moderator and the counters. The incident neutron flux was measured from the  $\gamma$ -rays of neutron capture by cadmium or indium samples. Such  $(n, \gamma)$  detectors consist of a NaI(Tl) crystal with  $\phi 70 \times 70 \text{ mm}$  and an FEhU-49 photomultiplier. The transmission was measured with an SNM-5 boron counter or three SNM-37 helium counters.

Calculated corrections for the resolution, the energy dependence of the flux detector efficiency and the multiple scattering of neutrons in the sample were made to the experimental results obtained. The magnitude of the correction for multiple processes in the fissionable sample was calculated by the Monte Carlo method; in the 0.025 eV energy region it was about 2.5%, increasing to 3.5% in the 1 eV region. The correction for resolution was about 2% in the 0.2 eV region in relation to the value at  $E = 0.0253 \text{ eV}$ . An inaccurate knowledge of this correction introduced a systematic error of the order of 0.5% almost comparable to the statistical error of measurement of  $\sim 0.7\%$ .

Table 1 gives the energy dependence obtained for  $\eta$  of  $^{233}\text{U}$  with flux measurements using an  $(n, \gamma)$ -detector with cadmium and indium samples. The energy dependence of the total cross-section for  $^{233}\text{U}$ , together with the absolute error, is given in Table 2. Table 3 shows the energy dependence of the  $^{233}\text{U}$  fission cross-section with flux measurement by  $(n, \gamma)$ -detectors with cadmium and indium samples. Table 4 contains the recommended values of  $\eta$  and the fission cross-section of  $^{233}\text{U}$  in the energy range 0.02-3 eV, obtained by averaging the measured values, which corresponds to a deterioration of resolution by a factor of about 2. The values of  $\eta$  from 1 eV to 3 eV are taken from our 1964 experiments, which showed that in this region  $\eta$  is independent of changes in resolution from 3 to 6  $\mu\text{s/m}$ . The energy dependence of the fission cross-section was obtained under the assumption that  $v = \text{const}$  in the neutron energy range studied; it is normalized to a value  $\sigma_f = 525.1$  b at  $E = 0.0253$  eV.

Table 1

Values of  $\eta$  for  $^{233}\text{U}$  (the incident neutron flux was measured using an  $(n, \gamma)$  detector and indium and cadmium samples)

E, eV	Flux for In		Flux for Cd	
	$\eta$	$\pm \Delta \eta$	$\eta$	$\pm \Delta \eta$
1,0672	2,244	0,037		
0,8968	2,310	0,029		
0,7641	2,332	0,030		
0,6588	2,277	0,032		
0,5739	2,308	0,032		
0,5044	2,325	0,029		
0,4463	2,372	0,036		
0,3935	2,323	0,028		
0,3577	2,326	0,031	2,372	0,027
0,3228	2,357	0,030	2,347	0,029
0,2923	2,322	0,029	2,355	0,029
0,2660	2,253	0,031	2,300	0,027
0,2441	2,240	0,030	2,269	0,025
0,2242	2,149	0,030	2,257	0,029

E <sub>γ</sub> , eV	Flux for In		Flux for Cd	
	$\bar{I}$	$\pm 4\bar{I}$	$\bar{I}$	$\pm 4\bar{I}$
0,2066	2,252	0,030	2,242	0,024
0,1910	2,219	0,025	2,215	0,026
0,1771	2,234	0,024	2,216	0,027
0,1647	2,206	0,027	2,235	0,024
0,1535	2,206	0,026	2,183	0,024
0,1434	2,215	0,026	2,175	0,023
0,1343	2,248	0,024	2,213	0,023
0,1261	2,240	0,023	2,230	0,023
0,1185	2,243	0,022	2,219	0,022
0,1117	2,239	0,022	2,243	0,022
0,1054	2,228	0,022	2,230	0,022
0,0996	2,256	0,025	2,235	0,021
0,0943	2,276	0,022	2,263	0,022
0,0894	2,283	0,023	2,275	0,022
0,0849	2,276	0,022	2,285	0,022
0,0807	2,262	0,023	2,256	0,022
0,0768	2,267	0,022	2,267	0,021
0,0732	2,272	0,022	2,285	0,022
0,0698	2,302	0,022	2,300	0,022
0,0667	2,287	0,022	2,294	0,021
0,0637	2,289	0,022	2,304	0,022
0,0610	2,281	0,022	2,290	0,022
0,0584	2,272	0,021	2,294	0,020
0,0560	2,290	0,022	2,288	0,021
0,0537	2,301	0,022	2,302	0,022
0,0516	2,285	0,024	2,303	0,022
0,0496	2,291	0,020	2,298	0,020
0,0477	2,294	0,022	2,304	0,019
0,0459	2,306	0,023	2,303	0,022
0,0442	2,311	0,021	2,305	0,021
0,0427	2,305	0,021	2,308	0,021
0,0411	2,292	0,021	2,299	0,021
0,0397	2,290	0,048	2,294	0,020
0,0383	2,296	0,021	2,310	0,020
0,0371	2,287	0,021	2,303	0,019
0,0358	2,302	0,020	2,306	0,020
0,0347	2,312	0,019	2,306	0,020
0,0335	2,306	0,020	2,299	0,020
0,0325	2,290	0,020	2,295	0,020
0,0315	2,285	0,020	2,286	0,020
0,0305	2,307	0,021	2,295	0,019
0,0298	2,292	0,021	2,316	0,021
0,0287	2,295	0,019	2,303	0,020
0,0279	2,300	0,021	2,288	0,020
0,0271	2,308	0,023	2,298	0,028
0,0263	2,293	0,021	2,302	0,021
0,0256	2,323	0,022	2,301	0,020
0,0249	2,290	0,020	2,300	0,019
0,0242	2,308	0,021	2,307	0,020
0,0235	2,304	0,021	2,293	0,020
0,0229	2,310	0,021	2,315	0,020
0,0223	2,305	0,021	2,282	0,019
0,0217	2,304	0,023	2,297	0,020
0,0212	2,279	0,023	2,288	0,020
0,0207	2,290	0,023	2,289	0,025
0,0201	2,285	0,022	2,292	0,029

Table 2  
Total cross-section of  $^{233}\text{U}$

$E, \text{ eV}$	$\sigma_t, \text{ b}$	$\pm \Delta \sigma_t, \text{ b}$	$E, \text{ eV}$	$\sigma_t, \text{ b}$	$\pm \Delta \sigma_t, \text{ b}$
1,0672	194,96	3,85	0,0610	382,84	1,91
0,8968	161,60	3,66	0,0584	389,09	1,96
0,7641	150,15	2,13	0,0560	396,62	1,88
0,6506	155,57	2,25	0,0537	403,41	2,48
0,5739	160,42	2,98	0,0516	412,46	2,01
0,5044	165,73	4,20	0,0496	420,67	2,02
0,4468	169,93	2,84	0,0477	429,42	2,21
0,3985	172,91	2,20	0,0459	435,77	2,45
0,3577	184,58	1,54	0,0442	446,33	2,10
0,3126	190,72	1,25	0,0427	453,81	2,16
0,2728	197,67	1,67	0,0411	462,10	2,42
0,2366	209,51	1,63	0,0397	468,67	2,61
0,2441	217,46	3,55	0,0383	477,93	2,55
0,2242	224,00	1,25	0,0371	485,49	3,05
0,2086	230,71	2,98	0,0358	492,12	2,29
0,1910	236,04	2,35	0,0347	502,71	2,73
0,1771	238,54	1,20	0,03335	510,63	2,48
0,1647	246,76	1,40	0,0325	521,23	2,65
0,1535	250,79	1,53	0,0315	526,12	3,01
0,1434	253,00	1,30	0,0305	537,03	2,57
0,1343	260,13	1,72	0,0296	543,73	2,72
0,1261	268,58	1,63	0,0287	551,69	3,07
0,1185	277,96	1,32	0,0279	559,00	2,59
0,1117	283,82	1,71	0,0271	570,27	2,95
0,1054	283,47	2,17	0,0263	581,45	4,66
0,0996	301,44	1,67	0,0256	591,13	3,00
0,0943	309,71	1,85	0,0249	594,77	3,15
0,0894	315,67	1,90	0,0242	595,98	3,43
0,0849	323,75	1,91	0,0235	610,78	3,86
0,0807	335,56	2,95	0,0229	613,92	3,52
0,0768	340,77	1,71	0,0223	628,15	3,08
0,0732	350,22	1,63	0,0217	631,74	3,03
0,0698	354,37	1,85	0,0212	640,02	3,28
0,0667	364,32	2,31	0,0207	645,42	3,53
0,0637	372,36	2,00	0,0201	657,56	3,61
0,0196	665,82	3,63	0,0114	839,50	7,40
0,0192	669,43	3,55	0,0122	843,14	6,63
0,0187	682,13	3,92	0,0119	849,66	4,15
0,0183	691,74	4,34	0,0117	852,15	5,00
0,0178	702,69	4,50	0,0115	865,24	4,62
0,0174	703,59	3,70	0,0112	875,38	8,00
0,0170	707,91	3,27	0,0110	876,40	6,05
0,0166	721,93	3,51	0,0108	884,65	5,88
0,0163	732,52	4,05	0,0106	899,63	5,36
0,0159	740,53	5,62	0,0104	915,45	7,10
0,0155	741,56	5,57	0,0102	914,22	8,21
0,0152	753,42	6,81	0,0101	931,45	11,22
0,0149	763,50	4,24	0,0099	943,46	6,12
0,0146	760,77	3,62	0,0097	936,92	9,06
0,0143	776,54	6,03	0,0095	937,53	11,68
0,0140	794,48	4,23	0,0094	952,74	8,17
0,0137	798,42	4,65	0,0092	979,67	10,54
0,0134	806,23	5,24	0,0091	977,07	6,60
0,0131	811,00	4,68	0,0089	992,29	12,25
0,0129	822,75	9,08	0,0088	989,98	8,27
0,0126	825,04	6,00	0,0086	1012,22	7,65

Table 3

Fission cross-section of  $^{233}\text{U}$  (incident neutron flux measured using  $(n, \gamma)$  detector with indium and cadmium samples)

E, eV	Flux for In		Flux for Cd	
	$\sigma_f, b$	$\pm \sigma_f, b$	$\sigma_f, b$	$\pm \sigma_f, b$
1,0672	164	4,0		
0,8968	138	3,3		
0,7641	128	2,1		
0,6523	130	2,3		
0,5739	136	2,8		
0,5044	142	3,8		
0,4408	149	3,1		
0,3985	154	2,2		
0,3677	159	2,0	157	1,7
0,3428	167	1,9	166	1,8
0,2,28	170	2,1	173	2,1
0,2568	176	2,3	180	2,0
0,2441	182	3,5	184	3,3
0,2242	188	2,3	189	2,3
0,2066	195	3,3	194	2,9
0,1910	197	2,5	196	1,6
0,1771	200	1,8	199	2,2
0,1647	206	2,3	204	2,0
0,1535	210	2,4	207	2,2
0,1434	217	2,4	213	2,0
0,1343	222	2,3	220	2,2
0,1261	229	2,3	228	2,3
0,1185	237	2,0	234	2,0
0,1117	242	2,4	242	2,3
0,1054	249	2,7	250	2,6
0,0996	260	2,9	257	2,4
0,0943	269	2,6	266	2,6
0,0894	276	2,8	275	2,6
0,0849	282	2,7	283	2,7
0,0807	291	4,1	290	3,0
0,0768	296	2,8	296	2,6
0,0732	306	2,7	307	2,7
0,0698	314	2,9	314	2,9
0,0667	321	3,2	321	3,0
0,0637	328	3,1	330	3,2
0,0610	336	3,1	338	3,1
0,0581	341	3,2	344	3,1
0,0560	350	3,2	350	3,0
0,0537	358	3,5	353	3,5
0,0516	364	3,7	367	3,3

Continuation of Table 3

E, eV	Flux for In		Flux for Cd	
	$\sigma_{f,b}$	$\pm \sigma_{f,b}$	$\sigma_{f,b}$	$\pm \sigma_{f,b}$
0,0436	372	3,1	374	3,2
0,0477	381	3,7	382	4,0
0,0459	388	3,9	389	3,8
0,0442	399	3,5	398	3,5
0,0427	404	3,6	405	3,5
0,0411	409	3,3	411	3,6
0,0397	415	4,0	415	3,8
0,0383	424	3,3	426	3,6
0,0371	430	4,2	434	3,9
0,0358	439	3,7	439	3,4
0,0347	450	3,8	449	3,9
0,0335	453	3,8	455	3,9
0,0325	463	4,1	464	4,1
0,0315	466	4,2	466	4,2
0,0305	480	4,3	473	3,9
0,0296	484	4,4	488	4,3
0,0287	492	4,3	493	4,4
0,0279	499	4,4	499	4,2
0,0271	511	5,1	509	6,3
0,0263	513	5,7	513	5,7
0,0256	525	5,1	522	4,4
0,0249	530	4,6	531	4,5
0,0242	534	5,1	534	4,3
0,0235	547	5,4	544	5,2
0,0229	552	5,2	552	5,0
0,0223	563	5,1	557	4,6
0,0217	566	5,6	564	4,7
0,0212	567	5,7	569	4,9
0,0207	575	6,0	575	6,5
0,0201	584	5,8	586	7,6



Table 4

Recommended values of  $\eta$  and  $\sigma_f$  for  $^{233}\text{U}$  in the energy range 3-0.2 eV

E, eV	$\eta$	$\pm \Delta \eta$	$\sigma_f, \text{b}$	$\pm \Delta \sigma_f, \text{b}$	Remarks
2.76	1.64	0.04			
2.71	1.71	0.04			
1.67	2.08	0.05			
1.35	2.14	0.05			
1.12	2.25	0.07			
0.764	2.293	0.02	139	1.0	
0.799	2.341	0.02	154	1.0	
0.593	2.328	0.02	171	1.0	
0.244	2.284	0.02	182	1.0	
0.207	2.240	0.02	192	1.0	
0.191	2.223	0.02	197	1.0	
0.165	2.201	0.02	205	1.0	
0.119	2.229	0.02	235	1.0	
0.094	2.263	0.02	275	1.0	
0.090	2.290	0.02	314	1.0	
0.061	2.296	0.02	349	1.5	
0.040	2.304	0.02	389	2.0	
0.0384	2.304	0.02	424	2.0	
0.0325	2.295	0.02	463	2.0	
0.0279	2.301	0.02	501	2.0	
0.0242	2.302	0.02	537	2.0	
0.0212	2.290	0.02	571	2.0	

Normalization region,  
mean values:  
 $\eta = 2.297$   
 $\sigma_f = 525.1$

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EVALUATION OF NUCLEAR DATA FOR  $^{235}\text{U}$  IN THE NEUTRON ENERGY RANGE  $10^{-4}$ -1 eV

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Paper submitted to the journal "Vestsi Akademii Navuk BSSR, Seryya Fizika-Ehnergetychnykh Navuk"

Data on  $\sigma_a$ ,  $\sigma_f$ ,  $\alpha$  and  $\eta$  for  $^{235}\text{U}$  were evaluated for the neutron energy range  $10^{-4}$ -1 eV on the basis of analysis of experiments conducted by various authors. The "best" curves for  $\sigma_a(E)$  and  $\sigma_f(E)$  were plotted, and  $\eta(E)$  curves measured directly and obtained from the ratio of  $\sigma_f$  to  $\sigma_a$  and from  $\alpha$  were compared.

EVALUATION OF NUCLEAR DATA FOR  $^{235}\text{U}$  IN THE FORBIDDEN RESONANCE REGION

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Seriya fiziko-ehnergeticheskikh nauk"

The average  $^{235}\text{U}$  resonance parameters in the forbidden resonance range of energies (0.15-100 keV) were obtained both by averaging the parameters in the forbidden resonance range ( $\langle \Gamma_\gamma \rangle$  and  $\langle D \rangle_\tau$ ) and by fitting to the experimental cross-section data. In order to take into account the structure of the cross-sections  $\sigma_t$  and  $\sigma_f$ , the parameters  $\langle \Gamma_n \rangle_3$  and  $\langle \Gamma_n \rangle_4$  were fitted to the mean cross-sections  $\langle \sigma_t \rangle$  for the energy range chosen, while  $\langle \Gamma_f \rangle_4$  was fitted to the cross-sections  $\langle \sigma_f \rangle$ . In obtaining  $\langle \Gamma_f \rangle_8$  the competition of neutron inelastic scattering was taken into account. In the calculations of the cross-sections the force function  $S_1$  was taken to be constant, and  $S_0$  was determined by fitting to the experimental values of  $\sigma_t$ . With the parameters obtained it was possible to describe fairly well all the experimental data on  $\sigma_t$ ,  $\sigma_f$  and  $\alpha$ . The calculation with the values  $\langle \Gamma_n \rangle_\tau^l$  obtained without fitting of the force function  $S_0$  to the data on  $\sigma_t$  showed that in the region below 10 keV no satisfactory description of the experimental data on  $\alpha$  is achieved if the fluctuations in  $\sigma_t$  are not taken into account.

EVALUATION OF NUCLEAR CONSTANTS FOR  $^{235}\text{U}$  IN THE NEUTRON ENERGY  
RANGE  $10^{-4}$  eV-15 MeV

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Paper submitted to the scientific and technical publication  
"Yadernye konstanty" (Nuclear constants)

The authors briefly describe the results of evaluating the complete file of constants for  $^{235}\text{U}$  in the neutron energy range  $10^{-4}$  eV-15 MeV. For the evaluation a detailed analysis was made of all available experimental data for  $^{235}\text{U}$ , and a number of methods of evaluating constants was developed and applied for  $^{235}\text{U}$ , including: obtaining of self-consistent resonance parameters of the Adler parameter type, obtaining of average statistical parameters of the nucleus in the forbidden resonance region of energies (0.15-100 keV), by means of which it was possible to calculate the average cross-sections in this region ( $\sigma_f$ ,  $\sigma_t$ ,  $\alpha$ ), analysis of the angular distributions of elastically scattered neutrons using expansion in the Bessel and Legendre functions, calculation of spectra of  $\gamma$ -rays emitted by  $^{235}\text{U}$  on capture, fission and inelastic scattering of neutrons, and calculation of inelastic scattering cross-sections taking into account the competition of fission. The evaluated  $^{235}\text{U}$  data were transmitted to the Obninsk Nuclear Data Centre in the SOKRATOR format.

ANALYSIS OF NEUTRON INELASTIC SCATTERING ON THE  $^{235}\text{U}$  NUCLEUS

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Paper submitted to "Izvestiya Akad. Nauk BSSR,  
Seriya fiziko-ehnergeticheskikh nauk"

The authors calculated the cross-sections for neutron inelastic scattering by the  $^{235}\text{U}$  nucleus involving excitation of both discrete and continuous level spectra in the energy range up to 3.5 MeV. The calculations took into account effects of competition from the radiative neutron capture and fission reactions and, in the region of allowed levels, the effect of width fluctuation. The calculations were performed for 19 levels of the target nucleus. The results agree closely with the experimental data in the energy region in question, unlike those of the British and German nuclear data libraries.

EVALUATION OF THE NEUTRON INELASTIC SCATTERING CROSS-SECTION FOR  $^{235}\text{U}$

G.V. Antsipov, V.A. Kon'shin, V.P. Korennoj

Presented at the Third Conference on Neutron Physics (Kiev, June 1975)

Using the Hauser-Feshbach formalism taking into account the fluctuations in neutron and fission widths the authors calculated the neutron inelastic scattering cross-sections in the region of allowed and overlapping levels of the target nucleus (0-3.5 MeV). In this calculation the competition of capture and fission was taken into account, as was that of inelastic scattering accompanied by excitation of a continuous spectrum in the calculation of the excitation cross-sections of the various levels. The calculations were performed for 19 levels (415 keV). On the basis of the theoretical calculations and analysis of the direct and indirect experimental information on  $\sigma_n'$  the neutron inelastic scattering cross-sections for  $^{235}\text{U}$  were evaluated in the energy range up to 15 MeV.

EVALUATION OF THE NEUTRON INELASTIC SCATTERING CROSS-SECTIONS FOR  $^{235}\text{U}$

G.V. Antsipov, V.A. Kon'shin, V.P. Korennoj, E.Sh. Sukhovitskij

Paper submitted to the scientific and technical publication  
"Yadernye konstanty" (Nuclear constants)

The authors evaluated the cross-sections  $\sigma_n'$  for  $^{235}\text{U}$  in the range from threshold to 15 MeV. The original data for the evaluation were the results of calculations in the energy range up to 3.5 MeV, experimental data on  $\sigma_n'$  and also indirect experimental information. The evaluated data obtained are compared with the experimental results and the data of the British and German nuclear data libraries. Agreement with experiment was found over the entire energy range. In the region below 1.5 MeV there is a discrepancy between the results of the present work and the data of the other libraries, which are significantly lower than the experimental values. The paper gives the numerical values of the evaluated data.

RESONANCE PARAMETERS OF  $^{235}\text{U}$  IN THE ENERGY RANGE UP TO 140 eV

V.A. Kon'shin, G.B. Morogovskij, E.Sh. Sukhovitskij

Presented at Third Conference on Neutron Physics (Kiev, June 1975)

In order to obtain the resonance parameters we analysed the available experimental data on fission and capture cross-sections and the total cross-section of  $^{235}\text{U}$  in the energy range up to 140 eV. The presence of powerful interference between the levels due to their proximity, the existence of four fission channels and the asymmetry of the individual resonances made it necessary to use an Adler-Adler-type formalism.

From the experimental data which, it seems to us, contain no significant errors, we obtained Adler-type parameters for the independent description of each type of cross-section, and also self-consistent parameters with which the data on  $\sigma_t$ ,  $\sigma_f$  and  $\sigma_c$  can be simultaneously described. The Breit-Wigner parameters were also obtained for consistent description of the cross-sections  $\sigma_t$ ,  $\sigma_f$  and  $\sigma_c$ . The cross-sections reconstructed from the multilevel parameters faithfully reproduce the behaviour of  $\sigma_i(E)$  over the entire energy range considered. The mean statistical parameters  $\langle D \rangle$ ,  $\langle \Gamma_\gamma \rangle$  and  $\langle \Gamma_f \rangle$  were determined.

EVALUATION OF THE ANGULAR DISTRIBUTIONS OF NEUTRONS ELASTICALLY SCATTERED ON  $^{235}\text{U}$

E.Sh. Sukhovitskij, A.R. Benderskij, V.A. Kon'shin

Paper submitted to the journal "Izvestiya Akad. Nauk BSSR, Seriya fiziko-ehnergeticheskikh nauk"

The paper presents an analysis of experimental data on the angular distributions of elastically scattered neutrons for  $^{235}\text{U}$  in the neutron energy range 0.05-15 MeV. The results are given in the form of a Legendre polynomial expansion. This expansion was obtained using the representation of angular distributions by Bessel function expansion. It was shown that by using Bessel functions certain disadvantages of expansion in Legendre polynomials can be avoided - the negative cross-sections, the large number of expansion factors, and the need for experimental data for the angles  $0^\circ$  and  $180^\circ$ .

EVALUATION OF THE NUCLEAR REACTION CROSS-SECTIONS FOR  $^{239}\text{Pu}$  IN THE  
RESONANCE RANGE OF ENERGIES DURING THE PREPARATION OF A  
COMPLETE FILE OF CONSTANTS

V.A. Kon'shin, G.B. Morogovskij, E.Sh. Sukhovitskij

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Seriya fizika-ehnergetychnykh navuk"

The authors analysed the available experimental data on  $\sigma_f$ ,  $\sigma_t$  and  $\sigma_c$  in order to determine the resonance parameters of  $^{239}\text{Pu}$ . The resonance parameters obtained by simultaneous fitting are given in the paper. Comparison of the mean cross-sections with experimental data shows satisfactory agreement between them. The mean values of the resonance parameters are given in the paper.

NUCLEAR DATA FOR  $^{239}\text{Pu}$  IN THE FORBIDDEN RESONANCE REGION  
OF NEUTRON ENERGIES

G.V. Antsipov, A.R. Benderskij, V.A. Kon'shin, E.Sh. Sukhovitskij

Paper submitted to "Izvestiya Akad. Nauk BSSR,  
Seriya fiziko-ehnergeticheshykh nauk"

The authors obtained the mean resonance parameters for  $^{239}\text{Pu}$  in the energy range 0.3-30 keV. The data were obtained from parameters in the allowed resonance region and also by fitting to evaluated data on the total interaction cross-section  $\sigma_t$  and the fission cross-section  $\sigma_f$ . The quality criterion for the results obtained is comparison of the calculated values of  $\alpha$  with the experimental values. With the mean resonance parameters obtained it is possible to calculate  $\alpha$  with satisfactory accuracy. Some discrepancy is observed in the region above 10 keV, perhaps because the neutron inelastic scattering reaction was not taken into account.

ANALYSIS OF NEUTRON INELASTIC SCATTERING ON  $^{239}\text{Pu}$

G.V. Antsipov, V.P. Korennoj, V.I. Martynyuk

Presented at Third Conference on Neutron Physics  
(Kiev, June 1975)

Using the framework of the statistical model,  $\sigma_n^i$  on the  $^{239}\text{Pu}$  nucleus was calculated for the neutron energy range  $E = 0.1-4$  MeV, taking into account the competition from the  $(n, \gamma)$  and  $(n, f)$  reactions and the width fluctuation effect. At energy  $E_n > 0.5$  MeV the contribution of the continuous spectrum was taken into consideration. The fission widths were evaluated up to an energy of 0.5 MeV on the basis of the channel theory of fission, and at higher energies in an approximation of independence of  $\tau$ . The calculation results were compared with experimental data and existing values for  $\sigma_n^i$  and  $\sigma_n^{Ej}$ .

EVALUATION OF THE CROSS-SECTIONS OF THE  $(n, 2n)$  AND  $(n, 3n)$   
REACTIONS FOR  $^{239}\text{Pu}$

E.Sh. Sukhovitskij, V.A. Kon'shin

Paper submitted to "Izvestiya Akad. Nauk BSSR,  
Seriya fiziko-ehnergeticheskikh nauk"

The authors propose a model for calculating the cross-sections of the  $(n, 2n)$  and  $(n, 3n)$  reactions. Experimental data on the interaction of neutrons with nuclei were used in constructing the model. The cross-sections of the  $(n, 2n)$  and  $(n, 3n)$  reactions for  $^{235}\text{Pu}$  and  $^{238}\text{U}$  are given, and the results of this work are compared with other evaluations.

EVALUATION OF THE NEUTRON CROSS-SECTIONS OF  $^{240}\text{Pu}$  IN  
THE FORBIDDEN RESONANCE REGION

G.V. Antsipov, A.R. Benderskij, V.A. Kon'shin, E.Sh. Sukhovitskij

Presented at Third Conference on Neutron Physics  
(Kiev, June 1975)

On the basis of the available experimental data the mean resonance parameters were evaluated taking into account the competition of inelastic scattering.

In the evaluation the concept of a double-humped structure of the fission barrier was used, which made it possible correctly to evaluate the fission penetration and width fluctuation factors. The evaluated parameters of the two-humped fission barrier are given. With the mean resonance parameters obtained it is possible to calculate the cross-sections  $\sigma_t$ ,  $\sigma_{n\gamma}$ ,  $\sigma_f$ ,  $\sigma_{nn}$  and their errors in the energy range 1-100 keV. This paper is part of the work on compilation of a complete file of constants for  $^{240}\text{Pu}$ .

EVALUATION OF THE NUCLEAR CONSTANTS OF  $^{240}\text{Pu}$  FOR COMPILING  
A COMPLETE FILE

G.V. Antsipov, A.R. Benderskij, V.A. Kon'shin, E.Sh. Sukhovitskij

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(Kiev, June 1975)

The paper contains a compilation of all available experimental data on  $^{240}\text{Pu}$ . On the basis of this compilation the neutron cross-sections are evaluated for the range  $10^{-4}$  eV-15 MeV using the concept of a double-humped fission barrier structure. The cross-sections in the resonance region are described by resonance parameters.

For the forbidden resonance region the mean resonance parameters are given as functions of energy. In evaluating the cross-sections and angular distributions in the high-energy region, where experimental data are lacking, the statistical model of the nucleus and the connected channel model are used. The spectra of neutrons and gamma rays from inelastic processes are evaluated. The cross-sections are presented in the usual format.



NUCLEAR PHYSICS CONSTANTS FOR TRANSPLUTONIUM ELEMENTS

V.A. Kon'shin

Paper published in "Izvestiya Akad. Nauk BSSR, Seriya fiziko-ehnergeticheskikh nauk"

The paper contains evaluated nuclear data for transplutonium elements at thermal energies. The available experimental data on  $\bar{\nu}$ ,  $\sigma_f^{2200}$ ,  $\sigma_c^{2200}$ ,  $I_f$  and  $I_c$  for  $^{247}\text{Cm}$ ,  $^{249}\text{Cf}$  and  $^{251}\text{Cf}$  are surveyed and analysed. A table gives nuclear data for the chain of formation of the elements from  $^{239}\text{Pu}$  to  $^{257}\text{Fm}$ .

CALCULATION OF THE ACCUMULATION OF TRANSPLUTONIUM ELEMENTS  
IN THERMAL AND RESONANCE NEUTRON FLUXES

G.V. Antsipov, V.A. Kon'shin, G.A. Sarnadskij

Presented at Third Conference on Neutron Physics  
(Kiev, June 1975)

The accumulation of transplutonium elements in thermal and resonance neutron fluxes was calculated. The dependence of the yields of these elements on the relation of the thermal flux to the epithermal flux in a constant full flux and on the neutron temperature was studied. The influence of errors in the experimental data for  $\sigma_f$ ,  $\sigma_c$ ,  $I_\gamma$  and  $I_c$  on the values for transuranium element accumulation was studied.

EVALUATION OF  $\alpha(^{235}\text{U})$  IN THE ENERGY RANGE 0.1 keV-15 MeV

G.V. Antsipov, V.A. Kon'shin

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(Kiev, June 1975)

The available experimental data on  $\alpha(^{235}\text{U})$  in the energy range above 100 eV were analysed from the point of view of possible systematic errors affecting six aspects: standardization,  $\sigma_f$  comparison, function of gamma-ray and fission detectors, background determination and energy resolution.

Using the mean statistical parameters evaluated by us we calculated the value of  $\alpha$ . The influence of fluctuation in the force function  $S_0$  and of change in  $\langle \Gamma_\gamma \rangle$  on the results of the calculation of  $\alpha$  was studied.