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FUNCTIONALS OF NEUTRON CROSS-SECTIONS IN THE  
UNRESOLVED RESONANCE REGION**

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METHOD FOR ANALYTICAL CALCULATION OF GROUP-AVERAGED  
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ABSTRACT

A new method is suggested for analytical calculations of group-averaged experimentally measurable functionals of neutron cross-sections (transmission functions) in the unresolved resonance region for one neutron channel and many radiative channels. It can be used for correct calculation of these functionals for even-even heavy nuclei such as  $^{238}\text{U}$ ,  $^{240}\text{Pu}$ , etc.

The analytical approach suggested in Refs [1, 2] to construction of average cross-sections in the unresolved region can be generalized formally to the group characteristics used to describe the resonance self-shielding effects if it is taken into account that the total cross-section and the radiative capture cross-section are expressed only in terms of two possible combinations of statistically reproducible quantities  $(1 - iR_{nn})^{-1}$  and  $(1 + iR_{nn}^*)^{-1}$

$$\sigma = 2\pi k^{-2} (1 - \text{Re } S_{nn}) = \sigma_o \left[ \cos^2 \varphi - \frac{1}{2} \left( \frac{e^{-2i\varphi}}{1 - iR_{nn}} + \frac{e^{2i\varphi}}{1 + iR_{nn}^*} \right) \right], \quad (1)$$

where  $\sigma_o = 4\pi k^{-2}$ , and the radiative capture cross-section

$$\sigma_y = \sigma_o \left[ \text{Re} \frac{1}{1 - iR_{nn}} - \frac{1}{1 - iR_{nn}} \cdot \frac{1}{1 + iR_{nn}^*} \right]. \quad (2)$$

Then any physical cross-section functional  $\psi(\sigma_r, \sigma)$  can be formally represented as a function of variables  $g = (1 - iR_{mn})$  and  $g^* = (1 + iR_{mn}^*)$ . Let us assume that there exists an integral relation:

$$\psi\left(\frac{1}{g}; \frac{1}{g^*}\right) = \int_0^\infty \int_0^\infty dt dt' f(t, t') e^{-gt} e^{-g^* t'} = \int_0^\infty \int_0^\infty dt dt' f(t, t') e^{-(t+t')} e^{iR_{mn}t} e^{-iR_{mn}^* t'}, \quad (3)$$

then the average value of the functional will be

$$\langle \psi \rangle = \int_0^\infty \int_0^\infty dt dt' f(t, t') e^{-(t+t')} F(t, t'), \quad (4)$$

where  $F$  is the function defined and studied in Ref. [2] with  $t = \frac{u+v}{2}$  and  $t' = \frac{u-v}{2}$ .

Thus, the problem of finding  $\langle \psi \rangle$  is reduced here to constructing the corresponding function  $f(t, t')$  ("original"  $\psi$ ).

Let us consider the functional

$$\begin{aligned} \left\langle \frac{1}{\sigma_R + \sigma} \right\rangle &= \left\langle \frac{1}{\sigma_R + \sigma_o \cos^2 \varphi - \frac{1}{2} \sigma_o (g^{-1} e^{-2i\varphi} + g^{*-1} e^{2i\varphi})} \right\rangle = \\ &= \frac{1}{\sigma_R + \sigma_o \cos^2 \varphi} \left\langle \frac{gg^*}{gg^* - hg^* - h^*g} \right\rangle, \end{aligned} \quad (5)$$

where  $\sigma_R$  is the dilution cross-section,  $g = 1 - iR_{mn}$ ,  $g^* = 1 + iR_{mn}^*$

$$\eta = \frac{1}{2} \frac{\sigma_o}{\sigma_R + \sigma_o \cos^2 \varphi} e^{-2i\varphi} = \eta e^{-2i\varphi}. \quad (6)$$

We denote  $p = g-h$  and  $q = g^* - h^*$ , then

$$\frac{gg^*}{gg^* - hg^* - h^*g} = \frac{(p+\eta)(q+h^*)}{pq-\eta^2} = 1 + \frac{h^*}{q} + \frac{h}{p} + \frac{\eta^2}{pq-\eta^2} \left(2 + \frac{h^*}{q} + \frac{h}{p}\right). \quad (7)$$

Now averaging the terms in (7) one-by-one, we obtain:

$$\left\langle \frac{1}{q} \right\rangle = \left\langle \frac{1}{1+iR_{nn}^* - h^*} \right\rangle = \int_0^\infty e^{-\kappa(1-h^*)} \langle e^{-iR_{nn}^* t} \rangle dt = \frac{1}{1-h^* + S_n}, \quad (8)$$

where we used the result of the work, and similarly

$$\left\langle \frac{1}{p} \right\rangle = \frac{1}{1-h + S_n}. \quad (9)$$

Moreover, the representations of the functions in the form of a Laplace transform with respect to two variables are known [3]:

$$\begin{aligned} \frac{1}{pq-\eta^2} &= \int_0^\infty \int_0^\infty I_0(2\eta\sqrt{tt'}) e^{-tp} e^{-r'q} dt dt' \\ \frac{1}{p(pq-\eta^2)} &= \frac{1}{\eta} \int_0^\infty \int_0^\infty \frac{\sqrt{t'}}{t} I_1(2\eta\sqrt{tt'}) e^{-tp} e^{-r'q} dt dq \\ \frac{1}{q(pq-\eta^2)} &= \frac{1}{\eta} \int_0^\infty \int_0^\infty \frac{\sqrt{t'}}{t} I_1(2\eta\sqrt{tt'}) e^{-tp} e^{-r'q} dt dt', \end{aligned} \quad (10)$$

so that

$$\left\langle \frac{gg^*}{gg^* - hg^* - h^*g} \right\rangle = 1 + \frac{h^*}{1-h^*+S_n} + \frac{h}{1-h+S_n} + \eta^2 \int_0^\infty \int_0^\infty e^{-t(1-h)} e^{-t'(1-h^*)} \times \quad (11)$$

$$dF(t,t') \left[ 2I_0(2\eta\sqrt{tt'}) + \frac{h^*t' + ht}{\eta\sqrt{tt'}} I_1(2\eta\sqrt{tt'}) \right] dt dt' ,$$

i.e. we arrive at a representation of form (4).

This result can be used to construct the average values for a group of transmissions, regarding it as a Laplace image of  $e^{-n\sigma}$ :

$$\int_0^\infty \langle e^{-n\sigma} \rangle e^{-np} dn = \left\langle \frac{1}{p+\sigma} \right\rangle . \quad (12)$$

The resonance integral in the group can be represented as

$$\left\langle \frac{\sigma_r}{\sigma_R + \sigma} \right\rangle = i\eta \left\langle \frac{R_{nn}^* - R_{nn}}{gg^* - \eta g^* - h^*g} \right\rangle = i\eta \left\langle \frac{R_{nn}^* - R_{nn}}{pq - \eta^2} \right\rangle = i\eta \left\langle \frac{R_{nn}^* - R_{nn}}{(pq - \eta^2)} \right\rangle - \eta \int_0^\infty \int_0^\infty dt dt' I_0(2\eta\sqrt{tt'}) \times$$

$$e^{-t(1-h)} e^{-t'(1-h^*)} \left( \frac{\partial}{\partial t} + \frac{\partial}{\partial t'} \right) F(t,t') = S_n \eta \int_0^\infty du e^{-u(1-\eta \cos 2\varphi)} \times \quad (13)$$

$$\int_{-\nu}^{\nu} dv I_0(2\eta\sqrt{u^2 - v^2}) \cos(\eta v \sin 2\varphi) sh 2f \Phi(u, \nu) .$$

where  $u = t + t'$ ,  $\nu = t - t'$  and function  $\Phi$  was introduced and studied in Ref. [2]. In the same way we can also determine other group-averaged characteristics in the unresolved region: average cross-sections in experiments with filtered beams (self-indication) and the

$$\text{quantity } \left\langle \frac{1}{(\sigma_R + \sigma)^2} \right\rangle .$$

## REFERENCES

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