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Neutron Activation Cross Sections Measured at KRI in Neutron Energy Region 13.4 – 14.9 MeV

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December 2016

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Abstract

Results of measurements carried out at KRI for reaction cross sections induced by neutrons with energy around 14 MeV are gathered in the present publication. In total, 1547 experimental cross section values obtained for 242 reactions are given in Tables and Figures of this paper. All the recollected measurement results have been recalculated using the latest reference data on properties of target nuclei and reaction products. Also the reference cross-section values used in KRI experiments earlier were changed by the corresponding values taken from the recent IRDFF v1.05 evaluation. The revised experimental results supersede the corresponding KRI data published earlier. The work was sponsored by NDS of IAEA.

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1. INTRODUCTION

The present work is aimed for taking together results of the systematic activation cross section measurement for the reactions of importance for fusion and other applications. The experiments were carried out on the edge of the $20^{\text{th}} - 21^{\text{st}}$ centuries by a group of experimentalists at the Khlopin Radium Institute (KRI). The most important data were earlier published in [1 - 5].

Most KRI measurements were performed in the standard, rigidly fixed conditions that provided decreasing the random data uncertainties. The systematic uncertainty component related to the instrumentation used was reduced by obtaining the cross sections that are known with high accuracy using the same conditions.

When it was possible, the cross-sections for the excitation of metastable and ground states of the product nuclide were measured separately. In total, the **1547** experimental cross-section values obtained for **242** reactions are presented in Tables and Figures of this paper.

All the collected data has been revised with the latest reference parameters [26 - 29] which are the reference cross-sections, isotopic abundance of the target nuclei, and half-lives and gamma intensities of the reaction products. In this paper, the uncertainty contributions of the reference data are given explicitly in an additional column of the Tables that contain the cross-section measurement results. This was done to facilitate a possible next result recalculation that will be required if the reference data would be further changed.

One page for one reaction was chosen as the format of the experimental data presentation. There are the numeric information on measurement results given in Tables and the intercomparison of available experimental data and evaluations demonstrated in Figures. The relevant decay schemes and numeric decay data used are presented on the corresponding pages too.

The reference cross-section data on ${}^{27}Al(n, \alpha)^{24}Na$ and ${}^{93}Nb(n, 2n)^{92m}Nb$ is tabulated in Attachment. Both the new IRDFF v1.05 values used in this paper and the old KRI values used earlier are given in the Attachment Tables.

In the report body, the applied experimental methods are described in detail. Some original methods developed for special measurements are also considered.

The experimental data values given in Tables supersede the corresponding experimental KRI data published earlier.

The work was sponsored by NDS of IAEA.

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Some misprints revealed in the report version of March 2016 have been corrected in the present version (December 2016). The author thanks very much Marina Mikhailukova who noted and reported these errors to me.

The revealed and corrected misprints are shown in the table lower.

Page number	Position in the list	Version of March 2016	Version of December 2016
p.37	Data table for 29 Si(n, 2p) 28 Mg, incident energy E _n	14.8	14.80
p.39	Table at the page bottom, ³⁹ K abundance	95.2581 44	93.2581 44
p.104	Data table for ^{nat} Mo(n, x) ⁹⁶ Nb	Data of columns 2,3,4 corresponds to the real data of columns 1,2,3 but the real data of column 4 is absent	Corrected
p.105	Data table for ^{nat} Mo(n, x) ⁹⁷ Nb	Data of columns 2,3,4 corresponds to the real data of columns 1,2,3 but the real data of column 4 is absent	Corrected
p.123	Table at the page bottom, reaction product $T_{1/2}$	35.36 d 6	35.36 h 6
p.126	Table at the page bottom, reaction product	¹⁰⁵ Rh	¹⁰⁵ Ru
p.137	Table at the page bottom, reaction product $T_{1/2}$	7.45 h 1	7.45 d 1
p.208	Table at the page bottom, reaction product $T_{1/2}$	42.39 h 6	42.39 d 6
p.275	Data table for ${}^{27}Al(n, \alpha){}^{24}Na$ reference cross section	Data at 13.48 MeV are absent.	Corrected
p.276	Data table for ⁹³ Nb(n, 2n) ^{92m} Nb reference cross section	Data at 13.48 MeV are absent.	Corrected.

Mistakes revealed in the version of March 2016 and corrected in the present version

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

2. EXPERIMENTAL PROCEDURES

Activation cross-section measurements were carried out using a conventional, widely used scheme: irradiation – cooling – gamma counting. Briefly, the main features of the measurements considered are as follows:

- a) most reaction cross sections were measured at several neutron energy values inside the interval 13.4 14.9 MeV;
- b) the 27 Al(n, α) 24 Na and 93 Nb(n, 2n) 92m Nb reaction cross sections known with the accuracy of 0.4 0.8% were used for the neutron fluence determination as the reference data;
- c) the geometry of irradiation and gamma counting was fixed for most measurements, and some standard reaction cross sections known with an accuracy of 1-2% were also measured in the same conditions;
- d) the scattered neutron contribution was minimized by the use of thin-wall constructions and air cooling of the target;
- e) the changes of neutron flux during irradiation were measured by two independent scintillation detectors. One of them was continuously scanning around the target and monitoring both angular and energy distributions of neutrons;
- f) the neutron field parameters were calculated in detail using a special code that took into account the real target properties and experimental geometry.
- g) the data accumulation, treatment and presentation were computerized. All information about the measurements was stored in a structured data bank provided with service codes.
- h) some efforts were made to create a computer code that could estimate the gamma radiation of samples irradiated by neutrons. This code was used as an auxiliary tool for choosing the optimal parameters of measurements.

The issues outlined above are considered in more detail in the following section of this paper.

2.1. Neutron irradiation

A Neutron Generator NG-400 (of ICT type) was used for sample irradiation with neutrons produced in the T(d, n)⁴He reaction. Depending on requirements of the particular experiment, the accelerating voltage was 220 - 320 kV, the deuteron current was 0.01 - 0.5 mA and the Ti-T target thickness was $1 - 2 \text{ mg/cm}^2$.

The irradiation arrangement is shown in Fig. 2.1.1. The samples were gathered into assemblies that were mounted relative to the beam direction at the angles of 0, 40, 60, 80, 100, 120 and 140 degrees (the ring No. 1 for seven assemblies) or at 0, 40, 60, 75, 91, 107, 125 and 150 degrees (the ring No. 2 for eight assemblies). In this arrangement, the neutron energy interval 13.4 - 14.9 MeV was covered at the accelerator parameters indicated above.

A special attention was paid to minimizing the contribution of scattered neutrons to the primary neutron spectrum. The accelerator target chamber was constructed using thin details that have the wall thickness of 0.2 - 0.5 mm. The target warm was withdrawn by a focused air jet. The sample holder was also made from thin-wall constructions where the maximal thickness had the sample assemblies (~1 - 3 mm). Additionally, the role of scattered neutrons was reduced by using during irradiations a short target - sample distance (30 mm) that was much less than the distance between the target and other massive elements of the accelerator and experimental hall.





A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

Typically, 2 - 4 different materials were irradiated simultaneously. The samples were gathered in packets in which the first and the last positions were occupied by the monitor foils (Al or Nb). Occasionally, an additional monitor foil was inserted in the middle of the assembly. The fluence data determined by all monitor foils were treated jointly taking into account the real foil distance to the neutron source, the angular distribution of d-T neutrons and the real law of neutron flux reduction with distance. This decreased the uncorrelated part of monitor errors by a factor approximately 2 - 3 in comparison with the conventional method where the fluence is independently determined for each packet. The method is described in more detail below.

Variations of the neutron flux during irradiation were measured by two independent scintillation detectors, one of which was fixed rigidly and the other one was rotated around the target at the distance 1 m, over the angular range from -160 to +120 degrees. A standard period of 20 s was set for a rotation cycle, however, if necessary, that could be changed from 5 s to 2 min.

Detailed information on the angular and energy distribution of neutrons, and their variations during the irradiation, was written to a PC disk and used later during data processing.

2.2. Calculation of neutron field parameters

The space and energy distributions of neutrons generated in the ${}^{3}H(d,n)^{4}He$ reaction were calculated with taking into account the realistic experimental conditions, such as finite sizes of the beam and the sample, inhomogeneity of the tritium distribution in the target, changes of the energy and angle parameters of the beam due to slowing down, etc. As the base model, the method described in [20] was used. In the present work, this approach was extended to the events of non-coaxial geometry.

Calculations described are very important when the sizes of the irradiated samples are comparable with the sample distance to the neutron source. There are a number of circumstances when the samples are to be placed in close proximity to the target. This could be necessary, for example, if the measured cross-section is small or if the reaction product has long half-life or low gamma-ray intensity or if only a small amount of material is available for irradiation, etc.

Additionally, it is worth noticing that the role of scattered neutrons which distort the original neutron field is much less near the target.

Last but not least, the allocation of the irradiated samples close to the target allows accumulating the needed neutron fluence without excessive activation of the target chamber, accelerator elements, etc.

Some examples of calculation results are shown below. The dependence of neutron field parameters on the irradiation geometry is presented in Figs. 2.3.1 - 2.3.4. The influence of deuteron energy on neutron characteristics is shown in Figs. 2.3.5 - 2.3.6.

Most noticably, the finite geometry effects appear in a remarkable broadening of neutron energy distribution near 90° (Fig. 2.3.2) and in significant deviation of neutron flux changing with distance from the $1/r^2$ law (Fig. 2.3.4).



Fig. 2.3.1. Neutron spectra calculated for different sample diameters, 14 and 1 mm. Beam diameter is 3 mm and sample-to-target distance is 30 mm.



Fig. 2.3.2. Neutron spectrum full width at half of maximum (FWHM) dependence on the sample-to-target distance. Beam diameter is 3 mm and sample diameter is 14 mm.

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Fig. 2.3.3. Dependence of the averaged neutron energy on the sample-to-target distance. Sample diameter is 14 mm.



Fig. 2.3.5. The energy dependence of the 4π -solid angle integrated neutron flux.



Fig. 2.3.4. Change of the averaged neutron flux multiplied by the squared distance to target. Sample diameter is 14 mm.



Fig. 2.3.6. Neutron spectra calculated for different deuteron energies (solid line corresponds to $E_d = 200$ keV).

The real FWHMs are about of 200 keV for every irradiation angle at the sample-to-target distance 30 mm. Although the neutron energy spread is remarkable, it does not influence the cross sections measured because the change of the majority of cross sections around the neutron energy 14 MeV is smooth and can be well approximated by the linear dependence. It is worth mention that the method of the cross section approximation by linearization was successfully used in our previous publications.

Generally, if the cross section change is linear and the neutron energy distribution is symmetric then the average neutron energy is the value that sufficiently characterizes the cross section measured. These conditions are valid almost everywhere in the present measurements. This is why the uncertainty of the average neutron energy (not FWHM!) is given in our present data illustrated in the Figures.

2.3. Reference cross sections

Cross-sections measured in the present work were determined relatively to the reference cross-sections which are either ${}^{27}Al(n,\alpha){}^{24}Na$, or ${}^{93}Nb(n,2n){}^{92m}Nb$, or both. The second reaction was favorable in practice because the ${}^{92m}Nb$ has eminently suitable decay properties: one dominant gamma-line 934.4 keV that is in the well calibrated region of gamma detector and the convenient half-life 10.15 d which provides gamma counting without hurrying. The ${}^{27}Al(n,\alpha){}^{24}Na$ reaction was the preferred reference at short irradiations.

Systematic cross-section measurements of reactions induced by 14-MeV neutrons were started at KRI in the early 1990s. At that time, improvements in nuclear data libraries were demanded, and several projects for conducting new experiments and extending evaluations were initiated. The ²⁷Al(n, α)²⁴Na cross-section around neutron energy 14 MeV has remained stable throughout the period of nuclear data revision. This is illustrated with the Fig. 2.2.1.

As can be seen in the Fig. 2.2.1 the "old" $^{27}Al(n,\alpha)^{24}Na$ cross-section values used at KRI measurements as the reference ones are in the good agreement with each of the modern evaluations. Most closely, the KRI data are to FENDL-3.0 and IRDFF v.1.05 evaluations.

The situation with the ⁹³Nb(n,2n)^{92m}Nb cross-section evaluation was quite different when the systematic cross-section measurements were started at KRI. Firstly, the experimental data was scattered widely (see Fig. 2.2.2). Secondly, there are only few evaluations that contained the split data for the ground and isomeric states. Therefore, obtaining ⁹³Nb(n,2n)^{92m}Nb cross-section data reliable enough for use as reference values during the determination of other measured cross-sections was a rather complicated problem.



Fig. 2.2.2. The ⁹³Nb(n,2n)^{92m}Nb cross section data of EXFOR and Neutron Libraries in the time when the KRI systematic cross-section measurement was started.



Fig. 2.2.1. Modern evaluations of the ${}^{27}Al(n,\alpha){}^{24}Na$ cross section in comparison with the data used at KRI measurements as the reference cross-section values.

In the circumstances outlined above, it was decided to bind the ${}^{93}Nb(n,2n){}^{92m}Nb$ cross section that was studied insufficiently to the ${}^{27}Al(n,\alpha){}^{24}Na$ cross section that was known more precisely.

A large body of ratios of the 93 Nb(n,2n) 92m Nb to the 27 Al(n, α) 24 Na cross sections were obtained during the initial stages of our measurements, as most of the KRI standard sample assemblies contained both niobium and aluminum foils that were sandwiched with the irradiated samples,

e.g. Nb - sample X1 - Al - sample X2 - Nb or Al - sample Y1 - Nb - sample Y2 - Al.

About of 20 independent ratio values were obtained using the model described above for all neutron energies. The ratios were A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI averaged to reduce random uncertainty in the data. Then the ${}^{93}Nb(n,2n)^{92m}Nb$ values were tabulated and used as the reference cross-section standard of the KRI measurements.

A good quality of the ${}^{93}Nb(n,2n)^{92m}Nb$ cross-section data used in the KRI experiments is confirmed by the contemporary evaluations shown in Fig. 2.2.3.

 27 Al(n, α)²⁴Na and 93 Nb(n,2n)^{92m}Nb reference cross-sections were stated once, at the very beginning of KRI experiments. These remained unchangeable during the total cycle of the KRI activation cross-section measurements.



Fig. 2.2.3. Contemporary evaluations of ⁹³Nb(n,2n)^{92m}Nb cross-section in comparison with the data used during KRI measurements as reference cross-section values.

One of the aims of this paper is to bring the reference decay and cross-section data used at KRI activation cross-section measurements into compliance with recommendations of up-to-date evaluated data files. Using this opportunity, the old (KRI) $^{27}Al(n,\alpha)^{24}Na$ and $^{93}Nb(n,2n)^{92m}Nb$ reference cross-section data should be changed with the values of the modern evaluations.

IRDFF v1.05 seems to be most preferable for this purpose because this evaluation is one of the latest (released 09 October, 2014), and contains detailed information on both crosssections and uncertainties. Additionally, the IRDFF v1.05 data for the two cross-sections considered here are in a total accordance with the data of the last available TENDL-2014 evaluation. Furthemore, the IRDFF v1.05 data for ${}^{27}Al(n,\alpha)^{24}Na$ and ${}^{93}Nb(n,2n)^{92m}Nb$ crosssections agree well with the reference data used at the KRI measurements. (Fig. 2.2.1 and Fig. 2.2.3). This means that the exchange of the old reference data (KRI) with new ones (IRDFF v1.05) does not cause a drastic change in the KRI activation data published earlier.

The old and the new reference cross-section data are tabulated in Table I and Table II of ATTACHMENT (pp.275-276).

2.4. Gamma counting

Two detectors were used for γ -counting of irradiated samples. These were a Ge(Li)-detector of the 160 cm³ sensitive volume and the energy resolution 4.0 keV at $E_{\gamma} = 1332.5$ keV, and a HPGe-detector that had a relative efficiency 24.7% and energy resolution of 1.8 keV at $E_{\gamma} = 1332.5$ keV. The HPGe-detector had a thin beryllium entrance window that extends the energy interval of gamma rays acceptable to detection up to 5 keV.

The detectors were enclosed in heavy shields consisting of consecutive layers of the lead, cadmium and steel (the Ge(Li) shield) or lead, copper and aluminum (the HPGe shield). The typical background gamma spectra of the shielded detectors are shown in Fig. 2.4.1.



Fig. 2.4.1. Background count rate of the shielded HPGe- and Ge(Li)-detectors used in the experiment.

A thorough knowledge of background spectrum details is often required for an accurate determination of low activities. To provide a good statistics of background spectra, these were measured by many small portions using the time windows when the detectors were free of other gamma measurements. In order to avoid broadening the peaks during summation of disconnected portions, a certain procedure was developed. First, the energy scale of every spectrum was recalibrated using well known background peak positions. Then, the spectra were deformed to the unified energy scale (e.g. 0.5 keV per channel as in Fig. 2.4.1) by a special computer code. After this, the spectra portions were summed. The background spectra shown in Fig. 2.4.1 corresponds to the total exposition of approximately 20 days for the HPGe and 9 days for the Ge(Li) detectors.

The relative full-energy peak efficiencies of the detectors were measured at different distances from the detector using the standard gamma sources (OSGI) as well as several radioactive nuclides produced at the Neutron Generator. The dependence of efficiency on energy and distance were approximated by analytical formulas. Two distance points were calibrated most carefully and set to the standard counting points for which the absolute full-energy peak efficiency was determined. The distances from the upper surface of the detector cup to the lower side of the sample were 33 and 112 mm for these two points. The absolute full-energy peak efficiency measured at the remote point for the HPGe detector is presented in Fig. 2.4.2.



Fig. 2.4.2. Absolute full-energy peak efficiency of the high purity germanium (HPGe) detector with the thin beryllium entrance measured at the distance of 112 mm from the detector cover top. A black solid curve is the analytical efficiency approximation used in activation cross-section calculations.

The HPGe detector used in the experiment had not only good energy resolution but also high registration efficiency, especially in the region of low gamma energies. The detector was operated in complicated measurements, e.g., detection of low activities embedded in high background.

The high detector sensitivity that is an undoubtedly positive property often has a negative effect related to the simultaneous registration of X-rays and other gamma-rays emitted in cascade with the defined gamma-line. The pulses correlated in time are summed in the detector and leave the peak area. The full-energy peak efficiency should then be corrected for the effect of gamma cascades.

The total registration efficiency is needed for calculation of gamma coincidence corrections. Experimental measurements of this quantity require the use of mono-energetic gamma sources. However, the sources of this category such as ²⁴¹Am, ¹⁰⁹Cd, ¹³⁹Ce, ²⁰³Hg, ¹³⁷Cs, ⁵⁴Mn and others are, strictly speaking, not mono-energetic because they emit the X-Rays also. These are registered by the HPGe detector and should be excluded by an additional correction procedure.

The problem of the total gamma registration efficiency could be solved by means of MCNP emulations. In the calculations we carried out, the real geometry of gamma counting was modeled. This included gamma detector crystal, detector capsule, sample holder, shield layers and other details. The MCNP emulations were performed for both HPGe- and Ge(Li)- detectors and also for both 112 mm and 33 mm standard gamma counting positions.

Results of MCNP calculations performed for the Ge(Li) detector is shown in Fig. 2.4.3. A gamma radiation source was supposed to be placed in the standard measurement point distanced at 33 mm from the top of the detector capsule.

It is seen in in Fig. 2.4.3 that the full-energy peak efficiencies calculated by means of MCNP are in a good agreement with those measured experimentally. Besides, it was discovered that the emulated response functions of various gamma sources correlate closely with corresponding gamma spectra measured in experiments. These facts could be considered as a strong evidence for the reliability of the total efficiency calculated by MCNP and, therefore, the accuracy of corrections for gamma cascade summing that was calculated in the present measurements.



Fig. 2.4.3. Absolute total, full-energy peak, one escape and double escape gamma efficiencies of the Ge(Li) detector calculated with MCNP for the distance of 33 mm from the detector cover top. A black solid curve is an approximation of the full-energy peak efficiency measured experimentally. The colored points correspond to the results of MCNP calculations,

One more conclusion can be drawn from the comparison of the calculated and measured characteristics of the detectors presented in Fig. 2.4.2 and Fig. 2.4.3. This concerns the accuracy of the full-energy peak efficiency. Since the methods used in the experiment and the MCNP calculation were totally independent, then the analysis of small differences between results of the two different approaches can shorten the uncertainty of the gamma peak efficiency approximation by the value $\Delta \varepsilon_{\gamma} \leq 1.2\%$ for both detectors and almost for the total gamma ray energy interval that was used in the measurements. A little uncertainty increasing was only admitted at the ends of the used gamma energy region.

After processing the gamma spectra of the samples activated by neutrons, the procedure of activation cross-section determination was started.

2.5. Algorithm of cross section determination

By definition, the cross section of the nuclear reaction A(n, x)B is determined by the relation:

$$\sigma_{AB} = \frac{N_B}{N_A \cdot \Phi_n},\tag{2.5.1}$$

where N_{B} is the number of nuclei **B** produced in the reaction;

 N_{A} is the number of the target sample nuclei **A**;

 Φ_n is the neutron fluence accumulated by the sample.

Now, consider the procedures used to determine each of the three values in the formula (2.5.1).

2.5.1. Determination of the target nuclei amount

The number of the initial nuclei, N_A , is connected with the mass of the sample through:

$$N_{A} = \frac{m \cdot Aw \cdot p_{1} \cdot p_{2}}{At}, \qquad (2.5.2)$$

where m is the sample mass, [g];

Aw is the Avogadro's number;

 p_1 is the abundance of the isotope A in the chemical element;

 p_2 is the share of the chemical elements in the chemical formula;

At is the atomic weight of atoms of the sample, [a.e.m];

2.5.2. Neutron fluence determination

There are several methods used to determine the integrated flux. The approach used in this work involves the use of ratios, whereby many uncertainty factors can be cancelled and the total data uncertainty is diminished.

In this approach, the integrated neutron fluence accumulated by the reference sample is determined as:

$$\Phi_n = \frac{N_{rB}}{N_{rA} \cdot \sigma_{rAB}},$$
(2.5.3)

where the formula is similar to (2.5.1) and "r" means "reference".

Usually, σ_{rAB} is known with good accuracy; N_{rA} and N_{rB} should be determined by the same procedures as N_A and N_B .

The fluence accumulated by each sample depends on the sample position in the assembly. This dependence is described by the $\frac{1}{r^2}$ -law if the neutron scattering and sizes of sample and target can be neglected. For close geometry, as shown in chapter 2.3, the accumulated fluence

changes with distance more slowly (see Fig. 2.3.4, for example). More correctly, the dependence $\frac{1}{r^a}$ would be used, where $a \le 2$.

For the actual conditions of the experiment (the target-sample distance is 30 mm, beam diameter 5 mm, diameter of the sample 14 mm, and sample thickness 1 mm), it was estimated that $a \approx 1.97$ which approximates closely to the "normal" value 2. Nevertheless, for more generality, let us consider that:

$$\Phi_n \sim \frac{1}{r^a} \tag{2.5.4}$$

To simplify the following formulas one can introduce a new variable, the so called "reverse fluence" φ_n that is connected with the normal fluence Φ_n by:

$$\varphi = \Phi_n^{-\frac{1}{a}} \tag{2.5.5}$$

Then the reverse fluence $\varphi(r_i)$, accumulated by the sample i, placed at the distance r_i from the neutron source will be calculated as:

$$\varphi(r_i) = \varphi(r_1) \cdot \frac{r_i}{r_1}$$
(2.5.6)

where $\varphi(r_1)$ is the reverse fluence accumulated by the reference sample located at the distance r_1 from the neutron source.

Often, as in the present case, the two reference foils occupy the first and last positions in the assembly. For them,

$$\begin{cases} \varphi(r_1) = \varphi_0 \cdot r_1 \\ \varphi(r_2) = \varphi_0 \cdot r_2 \end{cases}$$
(2.5.7)

where φ_0 is the reverse fluence at the unitary distance from the neutron source;

 r_1 and r_2 are the positions of the two reference foils.

The solution of the system (2.5.7) is equivalent to drawing a straight line in the plane $\{r, \varphi_n\}$. The straight line goes out from the point [0,0] and must pass through the regions of two other points with coordinates $[r_1, \varphi(r_1)]$ and $[r_2, \varphi(r_2)]$ with minimal deviations. It is natural to use the method of weighted least squares for finding the slope of the line φ_0 :

$$w_1 \cdot [\varphi(r_1) - \varphi_0 \cdot r_1]^2 + w_2 \cdot [\varphi(r_2) - \varphi_0 \cdot r_2]^2 = \min$$
(2.5.8)

where w_1 and w_2 are the weights of measurements for the first and second reference samples.

Then,

$$\varphi_{0} = \frac{W_{1} \cdot \varphi(r_{1}) \cdot r_{1} + W_{2} \cdot \varphi(r_{2}) \cdot r_{2}}{W_{1} \cdot r_{1}^{2} + W_{2} \cdot r_{2}^{2}}$$
(2.5.9)

The reverse fluence accumulated by the sample located at the distance r_i will be determined as:

$$\varphi(r_i) = \varphi_0 \cdot r_i \tag{2.5.10}$$

If n sample assemblies arranged at different angles to the deuteron beam are irradiated simultaneously (as shown in Fig. 2.1.1) and every assembly contains two reference foils, then we will have a system of n equations like (2.5.8).

$$w_{1k} \cdot [\varphi(r_{1k}) - \varphi_0 \cdot r_{1k}]^2 + w_{2k} \cdot [\varphi(r_{2k}) - \varphi_0 \cdot r_{2k}]^2 = \min$$
(2.5.11)

where k runs from 1 to n.

Additionally, it is necessary to take into account that the effective center of the accelerator beam can be shifted from the geometrical center of the target. This will cause a change to the distance set used in the formulas above. The real distance from the neutron source to the sample, r_{ik}^* , is connected with the mechanically fixed distance, r_{ik} , from the target to the sample through:



$$r_{ik}^{*^{2}} = r_{ik}^{2} + u^{2} + 2r_{ik} \cdot u \cdot \sin(\theta_{k})$$
(2.5.12)

where u is the beam shift projection on the plane of sample assemblies; \mathcal{G}_k is the angle between the beam direction and the assembly k axis.

Additionally, the neutron source anisotropy $V(\varphi)$ is to be considered. This is defined as:

$$V(\mathcal{G}) = \frac{\Phi(\mathcal{G})}{\Phi_0} \tag{2.5.13}$$

where $\Phi(\mathcal{G})$ is the fluence of neutrons emitted in direction \mathcal{G} , and Φ_0 is the mean neutron fluence:

$$\Phi_{0} = \frac{N_{n_tot}}{4\pi}$$
(2.5.14)

where $N_{n tot}$ is the total number of neutrons emitted in all directions during the irradiation.

The "anisotropic reverse fluence", φ_k^* , is bound with the "isotropic reverse fluence", φ_k with relation:

$$\varphi_k^* = \varphi_k \cdot V_k^{-\frac{1}{a}} \tag{2.5.15}$$

where the index k corresponds to the angle \mathcal{G}_k .

Finally, the system of n equations which corresponds to the zero approach (2.5.11) can be rewritten with the modified variables:

$$w_{1k} \cdot [\varphi_k^*(r_{1k}^*) - \varphi_{0k} \cdot r_{1k}^*]^2 + w_{2k} \cdot [\varphi_k^*(r_{2k}^*) - \varphi_{0k} \cdot r_{2k}^*]^2 = \min$$
(2.5.16)

where k runs from 1 to n.

The variables marked with asterisk include the information on neutron source anisotropy and the possible shift of the beam position on the target. Therefore, the equations (2.5.16) describe the problem more realistic.

The system (2.5.16) has no analytical solution similar to (2.5.9). However, it can be easily solved with a computer using an iteration procedure. As a result, the total number of neutrons emitted during the irradiation time, $N_{n_{-tot}}$, and the real position of the beam on the target, u, will be deduced. It is worth noting that for determination of these two parameters, more than two independent equations are used. This means, the uncorrelated error of the neutron fluence determination decreases $\sqrt{n_r - 2}$ times, where n_r is the number of reference samples.

Finally, the neutron fluence accumulated by the i-sample in the k-assembly will be calculated by the formulas:

$$\boldsymbol{\varphi}_{ik} = \boldsymbol{\varphi}_{0k} \cdot \boldsymbol{r}_{ik}^* \tag{2.5.17}$$

$$\varphi_{ok} = \left[\Phi_0 \cdot V(\mathcal{G}_k)\right]^{-\frac{1}{a}} \tag{2.5.18}$$

$$r_{ik}^{*2} = \sqrt{r_{ik}^{2} + u^{2} + 2r_{ik} \cdot u \cdot \sin(\varphi_{k})}$$
(2.5.19)

2.5.3. Determination of the amount of nuclei produced in the reaction

As indicated above, this analysis is restricted to activation cross section measurements. In this case, we need to know the number of nuclei of type **B**, N_B , that were produced in the reaction studied (formula 2.5.1). It is related to the induced sample activity, A_B , measured experimentally through the simple relation:

$$A_{B} = \lambda_{B} \cdot N_{B} \tag{2.5.20}$$

where λ_{B} is the decay constant.

The formula (2.5.20) would be quite enough for determination of the N_B value if all the **B**-nuclei were produced by a very short neutron flash and the sample activity was measured very quickly, just at the moment of the flash. Unfortunately, this ideal situation can never happen.

Firstly, the nuclei we are interested in are radioactive and amount of them is changing in the time according to the radioactive decay law:

$$N_{R}(t) = N_{R}(t_{00}) \cdot e^{-\lambda_{B} \cdot t}$$
(2.5.21)

where $N_{B}(t_{00})$ is the number of **B**-nuclei produced in the neutron flash;

t is the time gone after the irradiation moment.

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Secondly, various nuclei are generated at the neutron irradiation of the sample. Some of them can also produce the **B**-nuclei in the process of own decay. These are "**B**-parents". Let us denote the **B**-parents by index B^m (mother). Then, the change in the time of the **B**-nuclei amount will be expressed by the more complicated dependence:

$$N_{B}(t) = N_{B}(t_{00}) \cdot e^{-\lambda_{B} \cdot t} \cdot \left[1 + \frac{\sigma_{B^{m}}}{\sigma_{B}} \cdot \alpha \cdot \frac{\lambda_{B^{m}}}{\lambda_{B^{m}} - \lambda_{B}} \cdot \left(1 - e^{-(\lambda_{B^{m}} - \lambda_{B}) \cdot t}\right)\right]$$
(2.5.22)

where σ_{B^m} and σ_B are cross-sections of production of the corresponding nuclei at neutron irradiation;

 α is a part of mother nuclide decays that is resulting in generation of **B**-nuclides;

 $\lambda_{P^{m}}$ and λ_{B} are decay constants of the corresponding nuclei.

The formula (2.5.22) covers almost all situations that have practical significance because the production of the radioactive chains consisting of three or more members is rather unusual for irradiations involving fast neutrons.

Generally speaking, the ratio $\frac{\sigma_{B^m}}{\sigma_B}$ is unknown and has to be determined from experimental data. Moreover, it may depend on neutron energy and, hence, is specific for each sample packet.

data. Moreover, it may depend on neutron energy and, hence, is specific for each sample packet. There are a number of methods for its determination.

The situation is most easy when decays of the mother and daughter nuclei are accompanied by emission of gamma rays having different spectra. If the mother nuclide does not indicate itself by specific gamma radiation then the measurement of the decay curve shape can be used. Several gamma spectra are to be repeatedly counted for every sample in this case. The accurate σ_{B^m} data taken from other experiments or evaluations would be also sufficient for solution of the problem.

At the present stage let us assume that the mother-to-daughter cross section ratio is deduced already and try to consider in more detail the sample activity determination. The immediate result of gamma counting is the gamma peak area, $P(E_{\gamma})$. It is related to the amount of **B**-nuclei as:

$$P(E_{\gamma}) = Y(E_{\gamma}) \cdot \varepsilon(E_{\gamma}) \cdot \lambda_{B} \cdot \int_{t_{1}^{cnt}}^{t_{2}^{cnt}} N_{B}(t^{cnt} - t_{00}) \cdot dt^{cnt}$$

$$(2.5.23)$$

where t_1^{cnt} , t_2^{cnt} are the times of the beginning and the end of gamma counting;

 t^{cnt} is the time of gamma counting $t_1^{cnt} \le t^{cnt} \le t_2^{cnt}$;

 $Y(E_{\gamma})$ is the yield of the relevant gamma line with energy E_{λ} ;

 $\varepsilon(E_{\gamma})$ is the absolute peak efficiency of the gamma detector.

Formula (2.5.23) is written as if all the **B**-nuclei were produced at the time moment t_{00} . In reality, the sample irradiation was extended from t_1^{irr} to t_2^{irr} , and the irradiation intensity could vary arbitrary as $\phi_n(t^{irr})$. To take this into account we can write:

$$P(E_{\gamma}) = \frac{Y(E_{\gamma}) \cdot \varepsilon(E_{\gamma}) \cdot \lambda_{B} \cdot \int_{t_{1}^{irr}}^{t_{2}^{irr}} dt^{irr} \int_{t_{1}^{irr}}^{t_{2}^{ent}} \phi_{n}(t^{irr}) \cdot N_{B}(t^{ent} - t^{irr}) \cdot dt^{ent}}{\int_{t_{1}^{irr}}^{t_{2}^{ent}} \phi_{n}(t^{irr}) \cdot dt^{irr}}$$
(2.5.24)

Then, gathering the previous formulas we can obtain finally:

$$P_{k}(E_{\gamma}) = Y(E_{\gamma}) \cdot \varepsilon(E_{\gamma}) \cdot (t_{2k}^{cnt} - t_{1k}^{cnt}) \cdot \sigma_{B} \cdot N_{A} \cdot \Phi_{n} \cdot \lambda_{B} \cdot F_{k}(B, B^{m}, \phi)$$
(2.5.25)

where: $P_k(E_{\gamma})$ is the E_{γ} peak area determined at the k^{th} gamma counting;

 $Y(E_{y})$ is the yield (intensity) of the relevant gamma line with energy E_{y} ;

 $\varepsilon(E_{\gamma})$ is the absolute peak efficiency of the gamma detector;

 $(t_{2k}^{cnt} - t_{1k}^{cnt})$ is the exposition of the k^{th} gamma counting;

 $\sigma_{\scriptscriptstyle B}$ is the cross section of the reaction **A**(n, x)**B**;

 N_A is the number of **A**-nuclei in the sample;

 Φ_n is the neutron fluence accumulated by the sample;

 λ_{B} is the decay constant of the **B**-nuclei;

 $F_k(B, B^m, \phi)$ is the term responsible for non-ideality of the experiment. It is defined as:

$$F_{k}(B,B^{m},\phi) = \frac{\int_{t_{1}^{irr}}^{t_{2}^{irr}} dt^{irr} \int_{t_{1}^{ont}}^{t_{2}^{ont}} \phi_{n}(t^{irr}) \cdot e^{-\lambda_{B} \cdot (t_{k}^{ont} - t^{irr})} \cdot \left[1 + \frac{\sigma_{B^{m}}}{\sigma_{B}} \cdot \alpha \cdot \frac{\lambda_{B^{m}}}{\lambda_{B^{m}} - \lambda_{B}} \cdot \left(1 - e^{-(\lambda_{B^{m}} - \lambda_{B})(t_{k}^{ont} - t^{irr})}\right)\right] dt^{cnt}}{(t_{2k}^{ont} - t_{1k}^{ont}) \cdot \int_{t_{1}^{irr}}^{t_{2}^{ont}} \phi_{n}(t^{irr}) \cdot dt^{irr}}$$
(2.5.26)

The total algorithm of the activation cross-section determination described in the chapter (2.5) has been implemented in a computer code that was used at the KRI experiments. It should be noted that the algorithm described above is a generalization and includes many effects usually regarded as corrections. This allows data processing in a standard way that enhances the productivity of set-up and the reliability of results. Additionally, careful calculation of geometrical factors and detailed measurement of detector characteristics allow experiments to be undertaken with "close" geometries leading to optimization of the use of neutron flux generated by the Neutron Generator.

2.6. Cross section uncertainties

Uncertainties of many parameters contribute to the total uncertainty of the measured crosssection. The most important of them are as follows:

- decay data -(0.01% - 35%);

- area of γ -peak (1% 40%);
- efficiency of γ -detector (1.2% 1.5%);
- corrections for self-absorption and γ -cascade summing (0.01% 8.5%);
- reference cross sections used for neutron fluence determination (0.4% 0.8%);
- sample mass -(0.01% 0.1%);
- isotopic abundance -(0.01% 7.5%).

The contribution of uncertainties of other parameters is small and can be neglected. The total relative uncertainty of an individual cross section can be calculated in the standard way assuming that the uncertainty contributors are independent.

In some cases, the Monte-Carlo method was used for the uncertainty calculation. For this, the input parameters of the computer code that followed the formula (2.5.25) were considered as random variables having normal distributions with dispersions correspondent to uncertainties of the formula (2.5.25) parameters. Calculations were repeated for every new randomized set of the input parameters many times. The dispersion of the calculation results was taken as the output parameter uncertainty.

The Monte-Carlo procedure was used when the uncertainty contribution of the correction factor $F_k(B, B^m, \phi)$ defined by the formula (2.5.26) could not be neglected. Also, this method was used at calculations of isomeric ratios (sub-paragraph 2.7.2).

The Tables in Chapter 3 that contain the experimental data where a separate column was reserved for the reference data uncertainties ($\Delta \sigma_{ref}$) that are bound with the outer databases and are mainly independent on the experiment quality. The reference uncertainties were defined as:

$$\Delta \sigma_{ref} = \sqrt{\left(\Delta CS_{ref}\right)^2 + \left(\Delta I_{abu}\right)^2 + \left(\Delta Y_{\gamma}\right)^2 + \left(\alpha_d \cdot \Delta T_{\frac{1}{2}}\right)^2} \tag{2.6.1}$$

NOTE! All the uncertainties in (2.6.1) are <u>the relative uncertainties!</u>

where ΔCS_{ref} is the relative uncertainty of the reference cross section;

 ΔI_{abu} is the relative uncertainty of the target isotope abundance;

- ΔY_{γ} is the minimal relative uncertainty of the used gamma intensities;
- $\Delta T_{1/2}$ is the relative uncertainty of the half life of the reaction product;
- α_d is the correction factor responsible for the decay during irradiation, cooling and gamma counting.

The last parameter α_d is not easy to calculate exactly because the uncertainty arising during decays is involved in the factor $F_k(B, B^m, \phi)$ in a complicated form. The α_d values used in calculation of the reference data uncertainties are given in the text explicitly. The values I_{abu} , Y_{γ} and $T_{1/2}$ used for ΔCS_{ref} calculation are marked in the "decay data used" Tables by the bold font.

2.7. Examples of special measurements

As previously noted this paper uses one page for one reaction data presentation. Sometimes, the space reserved is not sufficient for a clear description of the experiment peculiarities. Some specific methods used in the nonstandard or problematic situations are in detail considered in the present chapter.

The main method of the activation cross section determination used in our work is tightly bound with the measurement of the induced gamma activity. However, some reaction products do not reveal themselves by emitting the gamma-rays that could be registered by the conventional gamma detectors. Nevertheless, the corresponding cross sections can be measured sometimes. In the next two sub-paragraphs, 2.7.1 and 2.7.2, methods are described that were applied to cross-section determination for the reactions that do not produce the gamma emitting nuclides.

Another difficulty in the reaction cross-section determination can emerge in the situation where the products of the different reactions have the common decay paths. An example of resolving a similar problem is considered in sub-paragraph 2.7.3.

Modification of the set-up made for the adaptation of it to the measurement of short-lived activities is described in subparagraph 2.7.4.

Additional details of the pioneer cross-section measurement carried out with the 241 Am, the material of high radio-activity, are presented in sub-paragraph <u>2.7.5.</u>

2.7.1. Determination of isomeric ratios using analysis of the decay curve shape

It was shown in the chapter 2.5 that two generations of the produced radioactive nuclei are always considered at the cross-section calculation in our work. The time behavior of the induced radioactivity is described then by:

$$N_{B}(t) = N_{B}(t_{00}) \cdot e^{-\lambda_{B} \cdot t} \cdot \left[1 + \frac{\sigma_{B^{m}}}{\sigma_{B}} \cdot \alpha \cdot \frac{\lambda_{B^{m}}}{\lambda_{B^{m}} - \lambda_{B}} \cdot \left(1 - e^{-(\lambda_{B^{m}} - \lambda_{B}) \cdot t}\right)\right]$$
(2.5.22)

where the formula terms denoted by "B" are the parameters related to the assigned nuclide produced at neutron irradiation immediately, and the terms denoted by " B^{m} " are related to the "mother" nuclide that produces the nuclide "B" in the decay process. It is significant that the "mother" nuclide is also produced at neutron irradiation.

The unknown ratio
$$\frac{\sigma_{B^m}}{\sigma_{B}}$$
 can be obtained from

the analysis of the decay curve shape described by the formula (2.5.22). In Fig. 2.7.1.1, it is illustrated the work of the computer code that was written for determination of this ratio, or more exactly, the value that is more convenient to use, i.e. the Isomeric Ratio (IR) that is defined as:

$$IR = \frac{\sigma_{B^m}}{\sigma_B + \sigma_{B^m}}$$
(2.7.1.1)



Interrelations between the cross-section ratio used in (2.5.22) and the isomeric ratio (2.7.1.1) are easy and everybody can write these on one's own.

It is worth emphasizing two important advantages of the IR determination described here. <u>Firstly</u>, high accuracy of IR can be obtained since only the relative activity change is analyzed. Consequently, many uncertainty sources such as the neutron monitor cross sections, sample mass, isotope abundance, reference gamma intensity, gamma self absorption, gamma spectrometer efficiency, etc. are excluded. <u>Secondly</u>, excitation cross-sections can be deduced for states which do not emit measurable radiation.

Fig. 2.7.1.1 relates to the measurement of 58 Ni(n, p) 58m,g Co cross sections. The metastable state 58m Co populates the ground state 58g Co via the isomeric transition which is strongly converted and cannot be registered by the gamma detector. The 58m Co half-life is 9.10 h. The 58g Co half-life is 70.86 d, and it emits gamma rays 810.8 keV with the probability 99.45%. The gamma-line is observable well. The change of the 810.8 keV count rate was used for analysis.



Fig. 2.7.1.2. Distribution of the IR values calculated with different input sets varied within own uncertainties. IR=0.547(5)

The best isomeric ratio was deduced in a computer code by the least squares method. For the IR uncertainty determination, the Monte-Carlo method was applied where the input data were varied inside own uncertainties that were suggested to have the normal distribution. In other words, the IR value was repeatedly calculated thousands of times with every new input parameter set that was generated from the original one by random Gauss scattering within own uncertainties of every input parameter such as the gamma peak area, time of gamma counting, half-lives of the metastable and ground states, neutron irradiation intensity, time of irradiation, etc. Then, the set of the IR values obtained was treated in a conventional way for determination of the average IR value and its dispersion. The work of the computer code is illustrated by Fig. 2.7.1.2.

The described method of the isomeric ratio determination was developed for studying the ${}^{58}\text{Ni}(n, p){}^{58\text{m,g}}\text{Co}$ and ${}^{59}\text{Co}(n, 2n){}^{58\text{m,g}}\text{Co}$ reactions but was extended later to other similar situations.

2.7.2. Isomeric ratio measurement using β -particles registered by γ -detector

It was shown in the previous sub-paragraph that the isomeric ratio can be obtained from the analysis of the decay curve shape. The method has the important advantage because it does not require the complicated absolute calibrations, calculation of numerous corrections but needs only the high registration stability of the parameter the count rate change of which is measured. This parameter could be, for example, the β -particle intensity.

Really, the β -particles with the energy more than approximately 0.5 MeV go free through the thin entrance window and are directly registered by the HPGe detector we used. The problem was to separate the signals produced by original β -particles from those generated by Compton- and other electrons in the gamma detector. The reaction ⁹³Nb(n, α)^{90m,g}Y proved to have the appropriate properties for the considered method application.

<u>Firstly</u>, there are only few reactions that produce considerable radiation in the niobium samples irradiated by neutrons. These reactions are ${}^{93}Nb(n, 2n){}^{92m}Nb$ (T_{1/2} = 10.15 d), ${}^{93}Nb(n, \alpha){}^{90m}Y$ (T_{1/2} = 0.133 d) and ${}^{93}Nb(n, \alpha){}^{90g}Y$ (T_{1/2} = 2.67 d). <u>Secondly</u>, the β -particles of the

 90g Y-decay have the high end-point energy of 2280 keV. <u>Thirdly</u>, the main part of gamma radiation emitted by irradiated niobium has the lower energy (E_{γ} < 480 keV for ^{90m}Y, E_{γ} < 950 keV for ^{92m}Nb and no gammas for ^{90g}Y).

Hence, the detector pulse height region (950 keV<E<1460 keV) is favorable for determination of the β -particle contribution. The region is limited from the left by the most intensive gamma-rays of the reaction products, and from the right by the prominent gamma-rays of the background.

Unfortunately, the chosen range of interest does not contain exclusively the unchangeably background and the changeably β -particle contribution. The ^{92m}Nb 1847.5 keV weak gamma-rays that have the intensity 0.85% per decay are admixed to this region also. It was necessary to take into account their changeable contribution as well.

For this, every niobium sample irradiated by neutrons was counted many times during several days after irradiation. A special computer code was created that divided the measured pulse height spectra into the components that have different time behavior. Three constituents were searched for: the background (unchanged in time), production of the ^{92m}Nb decay ($T_{1/2} = 10.15$ d) and production of the ^{90g}Y decay ($T_{1/2} = 2.671$ d). The dividing procedure was applied to the spectra with cooling time more than 2 days.

The three separated components of the total pulse height spectra obtained in such a way are shown in Fig. 2.7.2.1. The component intensities are normed to the two day cooling time. Also, the detector response functions emulated by MCNP are presented in Fig. 2.7.2.1. These were calculated for the β -particles of 90g Y and for the 1847.5 keV gamma rays of 92m Nb.



Fig. 2.7.2.1. Three pulse-height distribution components that have the different time behavior. Component points were extracted from the spectra counted for the Nb sample in the time interval of 2 – 10 days after irradiation. The solid lines are the corresponding response functions calculated by MCNP.

As it is seen in Fig. 2.7.2.1, the components extracted from the experimental spectra and the response functions calculated by MCNP are in a good agreement. Even the single-escape peak at $E_{\gamma} = 1336$ keV is equally reproduced in both amplitude distributions of the ^{92m}Nb constituent. A close match between the pulse height distributions extracted from the experimental spectra using a priory information on the expected time behavior of the spectra components and the detector

response functions to the corresponding radiations calculated by MCNP can be considered as an additional evidence for the correctness of the separation procedure we used.

So, if from the total number of pulses counted in the region 960 keV – 1450 keV are subtracted the number of background events and the number of events produced in ^{92m}Nb decay then the rest amount should be proportional to the number of β -particles emitted in the ^{90g}Y-decay.

The background was measured with high accuracy earlier (Fig. 2.4.1 on p.10). Long gamma counting made for every sample after 12 - 16 days after irradiation (when 90g Y decayed) allowed the experimental determination of the 92m Nb contribution in the region of interest. For further applications, it was linked to the intensity of the main gamma peak of 92m Nb with the energy of 934.4 keV. The result was also confirmed by MCNP calculations.

Fig. 2.7.2.2 demonstrates an example of the time behavior of the ^{90g}Y beta radiation intensity measured by the method described above (red points). The red solid line corresponds to the calculated decay curve of ^{90g}Y ($T_{1/2} = 64.10$ h) that have "a mother", ^{90m}Y ($T_{1/2} = 3.244$ h) produced with the probability that corresponds to the isomeric ratio value *IR*=0.43.

Simultaneously, areas of several gamma peaks were determined in the spectra measured by the same detector. For completeness, the changing of count rates of the 934.4 keV peak of ^{92m}Nb, the 479.5 keV peak of ^{90m}Y and the 1461.0 keV background peak are given in Fig. 2.7.2.2 also. Solid lines show the corresponding decay curves.



Fig. 2.7.2.2. Measured and calculated decay curves obtained for the Nb sample irradiated by neutrons of 14.86 MeV.

The isomeric ratio for reaction ${}^{93}Nb(n, \alpha){}^{90g,m}Y$ was determined as the best fit of the curve produced by formula (2.5.22) (solid red line) to the experimental points (red circles) that are the relative intensities of β -particles counted by the HPGe detector. For more detail see previous sub-paragraph, where isomeric ratio determination using the analysis of the decay curve is described.

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The experimental results of the ${}^{93}Nb(n, \alpha){}^{90g,m}Y$ cross sections are presented in p. 94-95.

2.7.3. Cross section separation for reactions that result in the common product

The problem is considered via the example of ${}^{50}V(n, \alpha){}^{47}Sc$ and ${}^{51}V(n, n'\alpha){}^{47}Sc$ reactions.

A very low abundance of ⁵⁰V (0.25%) in natural vanadium made it necessary to purchase an enriched sample. The sample used in the work had an enrichment of 17% for ⁵⁰V. It was quite enough for reliable measuring the ⁵⁰V(n, n' α)⁴⁶Sc cross section but for the ⁵⁰V(n, α)⁴⁷Sc cross section the contribution of the ⁵¹V(n, n' α)⁴⁷Sc reaction could not be ignored. To solve the problem, a sample of natural abundance was included in the assembly, back-to-back to the enriched sample. The summarized cross-section of the ⁴⁷Sc production in vanadium (σ_V) was exclusively measured for both samples. The individual values of the cross sections ⁵⁰V(n, α)⁴⁷Sc (σ_{50}) and ⁵¹V(n, n' α)⁴⁷Sc (σ_{51}) were determined as a solution of the system of two equations with two variables:

$$\sigma_{50} \cdot v_{50}^{1} + \sigma_{51} \cdot v_{51}^{1} = \sigma_{V}^{1}$$

$$\sigma_{50} \cdot v_{50}^{2} + \sigma_{51} \cdot v_{51}^{2} = \sigma_{V}^{2}$$
(2.7.3.2)

where

 v_{50}^1 and v_{51}^1 – are shares of isotopes 50 V and 51 V in vanadium of the first sample; v_{50}^2 and v_{51}^2 – are shares of isotopes 50 V and 51 V in vanadium of the second sample; σ_V^1 is the summarized cross section (V(n, x)⁴⁷Sc) obtained for the first sample;

 σ_v^2 is the summarized cross section (V(n, x)⁴⁷Sc) obtained for the second sample.

The solution of the equation system is well known:

$$\sigma_{50} = \frac{\sigma_V^1 \cdot v_{51}^2 - \sigma_V^2 \cdot v_{51}^1}{v_{50}^1 \cdot v_{51}^2 - v_{50}^2 \cdot v_{51}^1}$$
(2.7.3.3)
$$\sigma_{51} = \frac{\sigma_V^2 \cdot v_{50}^1 - \sigma_V^1 \cdot v_{50}^2}{v_{50}^1 \cdot v_{51}^2 - v_{50}^2 \cdot v_{51}^1}$$
(2.7.3.4)

If we recollect that $v_{50}^1 + v_{51}^1 = 1$ and $v_{50}^2 + v_{51}^2 = 1$ then the view of the equation system solution can be simplified:

$$\sigma_{50} = \frac{\sigma_{V}^{1} \cdot v_{51}^{2} - \sigma_{V}^{2} \cdot v_{51}^{1}}{v_{51}^{2} - v_{51}^{1}}$$

$$\sigma_{51} = \frac{\sigma_{V}^{2} \cdot v_{50}^{1} - \sigma_{V}^{1} \cdot v_{50}^{2}}{v_{50}^{1} - v_{50}^{2}}$$
(2.7.3.6)

The experimental results of the ${}^{50}V(n, \alpha){}^{47}Sc$ and ${}^{51}V(n, n'\alpha){}^{47}Sc$ cross sections are presented in p. 43-44.

2.7.4. Measurement of short lived reaction products

To extend the region of nuclear reaction cross-sections measured at KRI to the reactions that generate short-lived nuclei, the set-up was modified. A system of quick sample transportation from the measuring room to the neutron target and back was designed, made, and tested in experimental conditions. The time needed for transportation of the sample to the neutron target or to the gamma detector was about 2 s. The sample arrival and departure moments were determined by means of two sensor units which were connected with a computer via a CAMAC module. Time uncertainties of sample irradiation and gamma counting were estimated to be less than 0.5 s.

2.7.5. Measurement of ²⁴¹Am cross sections

Conversion and utilization of long-lived radioactive products and, especially, actinide elements are among the most important problems of nuclear industry. Given the significance of this issue, experimental data on nuclear cross-sections are scanty or often nonexistent in this field.

In this chapter, one of the first successful attempts for measuring neutron cross-sections on ²⁴¹Am is described. These measurements were very difficult from the experimental point of view. For example, the natural gamma activity of a ²⁴¹Am sample used was about 10⁸ Bq but the induced gamma activity of the most probable (n, 2n) reaction product was expected to be of order of 100 Bq. Other reactions that were planned to measure were yet less probable. Additionally, some relevant gamma peaks could be masked by the numerous gamma peaks generated by fission fragments.

The first problem to resolve was the suppression of intensive natural gamma radiation of americium samples. For this purpose, a leaden container with the 5 mm walls was used from the beginning. This was later changed to a cadmium screen because the gamma absorption in cadmium dependes more steeply on gamma energy near 100 keV, and this provides a higher registration efficiency for the cadmium absorber in the region of 100–300 keV at the equivalent suppression of 60 keV gamma radiation.

The second problem was the container had to be sealed absolutely reliably for neutron irradiation and gamma counting but could be easily unpacked for radiochemical cleaning of the samples in order to provide repeated use of a very expensive material and its final return to a host laboratory. The container and the near-detector arrangement are shown in Fig. 2.7.5.1.



Fig. 2.7.5.1. Geometry for gamma counting. The thickness of cadmium absorber was 3.5 mm. The lead tube was used to shield the gamma-detector from scattered gamma-radiation.

The third problem was detector calibration. To determine the efficiency for gamma counting with the lead container used at the beginning of experiment, we carried out measurements for a 92m Nb sample with a well determined activity which was placed inside the container in the position of the geometric center of the Am sample. Since the gamma radiation energies of 92m Nb and 240 Am, the product of 241 Am(n, 2n) 240 Am reaction, are near, any corrections for the efficiency registration difference were small. In particular, the difference in absorption of gamma rays with energies of 934.2 keV (92m Nb) and 987.8 keV (240 Am) by the lead container bottom was -1.8%; the difference in gamma counting efficiency without the container was +5.3%. The efficiency difference related to the form of samples and self-absorption was -0.5%. Thus, the total difference in gamma counting efficiency of 934.2 keV (92m Nb) and 987.8 keV (240 Am) was only 4.0%.

The dependence of gamma radiation attenuation values on the cadmium absorber thickness in the geometry presented in Fig. 2.7.5.1 was determined experimentally using a kit of three standard gamma sources (²⁴¹Am, ¹⁰⁹Cd and ¹⁵²Eu) which were placed in the position of the studied americium sample. The absorption values measured experimentally were compared with those calculated by the simple formula:
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$$C = e^{-\mu \cdot d}$$
 (2.7.5.1)

where C is the attenuation value;

 μ is the attenuation coefficient, [mg/cm²]⁻¹

d is the absorber thickness, $[mg/cm^2]$

The comparison results are shown in Fig. 2.7.5.2 where it is seen that the simple formula (2.7.5.1) describes the real gamma ray attenuation rather well at energies above 200 keV that cover the whole region of interest in the experiment on the neutron induced activation cross-section ²⁴¹Am. measurement for The deficit of the calculated attenuation values at lower gamma energies appears to be associated with the oblique tracks that become important in this case.



Fig. 2.7.5.2. Ratio of the measured gamma ray attenuation to the calculated one in dependence on gamma energy. The cadmium absorber thickness is 3.5 mm.

The main part of the natural gamma radiation

of ²⁴¹Am that has the energy 60 keV and less, was reliably suppressed by the constructions described above. However, some gamma-rays associated the with natural radioactive decay having higher energies and the intensities of order of 10^{-6} per decay were clearly observed in the gamma spectra counted with a non-irradiated americium sample. An example of such spectrum is shown in Fig. 2.7.5.3.



Fig. 2.7.5.3. Gamma spectrum of a non-irradiated Am-sample counted by the HPGe detector shielded from low-energy gamma radiation with cadmium absorber of 3.5 mm thickness.

The additional background gamma radiation shown in Fig. 2.7.5.3 has not only a negative but also a positive role because energies and intensities of the gamma rays that accompany the ²⁴¹Am natural decay are well known. The inclusion of the corresponding information into the process of the gamma detector efficiency calibration made the data on efficiency more accurate and reliable.

Besides, the absolute intensity values of the ²⁴¹Am natural gammas can be used for the determination of the sample mass.

The samples were prepared from an aqueous solution or from americium nitrate. The dimensions of the volume filled with the solution or americium nitrate varied between 8 and 16 mm in diameter and between 1.0 and 3.5 mm in height. Two standard niobium foils used for the neutron fluence determination, were fastened to the front and back surfaces of the inner container (Fig. 2.7.5.1).

The samples were of high purity. According to the certificate, the total concentration of alien fissile elements in the samples did not exceed 0.1%. However, for the present experiment, isotopic purity requirements were much higher. The stringent limitations concerned the possible admixture of the other americium isotopes, ^{242m}Am and ²⁴³Am, because the products of their natural decay emit the same gamma rays as ²³⁹Am and ²³⁸Np, the products of ²⁴¹Am(n, 3n)²³⁹Am and ²⁴¹Am(n, α)²³⁸Np reactions. The available experimental conditions (neutron flux, duration of irradiation and gamma counting) determined an upper limit on permissible admixture of other americium isotopes as 10⁻⁷ g/g. To make sure that the samples meet such conditions, a long gamma counting was carried out with the freshly prepared sample. No traces of other americium isotopes were revealed in the accumulated gamma spectra. The upper limit of possible content of ^{242m}Am and ²⁴³Am isotopes in the samples was estimated as less than 10⁻⁸ g/g.

For determination of neutron induced activation cross-sections, six irradiations were carried out with nine Am-samples. Typically, the irradiations lasted from 1 to 17 hours, and the total fluence accumulated by samples varied from $(1 \text{ to } 10)*10^{13} \text{ n/cm}^2$.

The gamma spectra of the irradiated samples were found to be very complex. They contained hundreds of gamma peaks, most of which belonged to the fission fragments. In order to increase the peak identification reliability, the peak half-lives were also analyzed. For this, every irradiated sample was counted several times during 200 h after irradiation.

For the ²⁴¹Am(n, 2n)²⁴⁰Am reaction, two gamma peaks were to be observed: E_{v1}=987.8 keV and $I_{\gamma 1} = 72.2\%$, and $E_{\gamma 2}$ =888.9 keV and $I_{\gamma 2}$ =24.7%. These peaks were revealed in the spectra, where their intensity ratio corresponded to the values given above, and their half lives were equal to 50.8 h which coincides with the half life of ²⁴⁰Am. (Fig. 2.7.5.4). Therefore, the 241 Am(n, 2n) 240 Am cross sections could be reliably calculated using the standard procedures without any complications. The data obtained is presented on p. 273 of this paper.

For the ²⁴¹Am(n, 3n)²³⁹Am cross section measurement, the experimental conditions were more sophisticated than for the reaction considered above. There were many factors that contributed to making this experiment more difficult, such as the smaller cross section, the smaller intensities of the ²³⁹Am gamma radiation; the lower energies of the gamma rays emitted, the higher background level in this region of



Fig. 2.7.5.4. Decay curves for the gamma rays 888.9 keV and 987.8 keV that were identified as the ²⁴⁰Am decay gammas.

gamma spectra. In addition, the interference probability with some gamma rays of the ²⁴¹Am natural decay or with some gammas of the numerous fission fragments is higher. These factors arose in reality during the experiment.



Fig. 2.7.5.5. Decay curves for the gamma rays with energies close to 277.6 keV. The magenta solid line corresponds to the ²³⁹Am decay gammas.

The ²⁴¹Am(n, α)²³⁸Np cross-section was most difficult for measuring because of its small magnitude that was expected to be of order of 1 mb. Remember that the fission cross section at neutron energy 14 MeV is about 2500 mb. Using tables of evaluated cumulative fission fragment yields one can estimate that more than 300 fission fragments are produced with the cross-section of 10 mb or higher at neutron irradiation of ²⁴¹Am. Really, several hundreds of gamma peaks were revealed during processing of gamma spectra of irradiated americium samples. Most of them were identified as fission fragment gamma rays.

The best region for observation of the ²³⁸Np decay gamma radiation proved to be the region around $E_{\gamma} = 1028.5$ keV that is shown in Fig. 2.7.5.6. There are not observed any intensive gamma peaks that could mask a very weak radiation of the ²⁴¹Am(n, α)²³⁸Np reaction. In Fig. 2.7.5.6, the range (1028.5±1.5) keV where

Three candidates were selected for ²³⁹Am radiation identification:

 $\begin{array}{lll} E_{\gamma 1}{=}226.4 \ keV \ and \ I_{\gamma 1}{=}3.3\%; \\ E_{\gamma 2}{=}228.2 \ keV \ and \ I_{\gamma 2}{=}11.3\%; \\ E_{\gamma 3}{=}277.6 \ keV \ and \ I_{\gamma 3}{=}15.0\%. \end{array}$

The first two candidates proved to be totally masked by very intensive gamma radiation of the ¹³²Te with E_{γ} =228.2 keV and $T_{1/2}$ =3.20 d.

The third candidate was also not simple. The decay curve for gamma radiation with the appropriate energy could not be approximated by the straight line in the logarithmic scale as was done for gamma rays of the 241 Am(n, 2n) 240 Am reaction. The studied decay curve consists of three contributions (shown in Fig. 2.7.5.5). Happily, half-lives of the contributors are very different, and the component related to the $^{241}Am(n, 3n)^{239}Am$ reaction was determined without significant additional uncertainty. The cross-section data obtained are presented on p.272.



Fig. 2.7.5.6. Fragment of the gamma spectrum in the vicinity of $E_{\gamma} = 1028.5$ keV that was expected in the ²³⁸Np decay. Gamma counting was started at 60 h after irradiation and was carried out during 100 h

one of the biggest gamma peaks of ²³⁸Np was to be observed is marked out by the dashed lines. Since no peak can be detected there the region between the dashed lines was interpreted as the background. The possible 1028.5 keV peak area was set as being less than 3 background statistical uncertainties that resulted in the upper limit of the ²⁴¹Am(n, α)²³⁸Np cross section presented on p.271.

A significant volume of information about the fission fragmens was obtained in this experiment as a byproduct of the irradiations. Since we had about two hundred decay curves for different gamma peaks obtained during eight days after irradiation, we were able to identify many gamma peaks belonging to fission fragments. Some examples of these decay curves are presented in Fig. 2.7.5.7 together with the identification results.



Fig. 2.7.5.7. Several decay curves of nuclides identified as fission fragments of ²⁴¹Am irradiated by neutrons.

It can be noted that some curves have a shape which is inherent to nuclide chains. These figures are correspondingly marked.

About twenty fission fragment yields were also obtained in the data processing. However, the fission data was considered as a byproduct of the main experiment on the ²⁴¹Am(n, 2n)²⁴⁰Am, ²⁴¹Am(n, 3n)²⁴⁰Am and ²⁴¹Am(n, α)²³⁸Np cross section measurement and therefore is not included in the present paper.

3. EXPERIMENTAL RESULTS

The experimental data obtained at KRI Neutron Generator NG-400 in neutron energy region 13.4 - 14.9 MeV are presented in the following pages. As discussed earlier a format showing "one reaction on one page" was chosen for the data presentation. The page structure is as follows:

Reaction name is in the page header.

A Table with experimental results is placed in the top left of the page. The Table consists of four columns: the neutron energy (E_n [MeV]), the reaction cross-section (σ [mb]), the cross-section total uncertainty ($\pm \Delta \sigma_{total}$ [%]) and the cross-section reference uncertainty ($\pm \Delta \sigma_{ref}$ [%]). The last column data is the uncertainty contributed by the reference data used for the cross-section calculation. A rule of $\Delta \sigma_{ref}$ calculation is previously described on p. 19. The correction factor α_d used in $\Delta \sigma_{ref}$ calculation is given in the last line of the Table. This line contains also an indication of the reference cross-section used.

To the right of the Table, a Figure is allocated where the present data are compared with results of other experiments and with available evaluations.

Under the Table and Figure, necessary comments and explanations are given to the experimental results presented.

Reference decay data used for the cross-section calculation is in the table at the page bottom. This includes the half-life of the reaction product, the energy and intensity of the gamma-lines used during data processing. The target isotop abundance used is given in this table too.

In addition, a picture showing the decay scheme of the nuclides related to studied reaction is placed near the table with the reference decay data. The picture is a copy fragment of relevant schemes presented in the Handbook "Tables of Isotopes" by Richard B. Firestone issued in 1996 [25]. Note that the half-life values presented in the picture may differ from the half-life values presented in the table with the reference decay data. The difference reflects the progress in Nuclear Structure Data that has happened over the last twenty years. It should be underlined that the latest decay data given in the table at the page bottom were used in the process of the cross section calculation.

23 Na(n, 2n) 22 Na



	³⁺ 2.6019 y EC ²² Na
^{0⁺} 22Ne	

Decay data used for 23 Na(n,2n) 22 Na.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²³ Na	100	NaCl	²² Na	2.6018 y 22	1274.5	99.940 <i>14</i>

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26 Mg(n, α) 23 Ne



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$^{27}\mathrm{Al}(\mathrm{n},\mathrm{p})^{27}\mathrm{Mg}$



 $^{27}\mathrm{Al}(\mathrm{n},2\mathrm{n})^{26\mathrm{g}}\mathrm{Al}$



The aluminum samples were irradiated at KRI, and the accumulated neutron fluence was determined using KRI methods. A very long half-life of ²⁶Al made the gamma counting of the irradiated aluminum samples ineffective. It was not done. The gamma-ray energy and intensity given in the table at the page bottom are presented for the sake of information completeness but were not used in this experiment.

The amount of ²⁶Al produced during the irradiations was measured via accelerator mass spectrometry (AMS) with the Vienna Environmental Research Accelerator (VERA) in the frame of IRK-KRI collaboration.

More exhaustively, the experiment on the ${}^{27}Al(n,2n){}^{26g}Al$ cross section measurement is described in [5].



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
²⁷ A1	100	Al-metal	²⁶ Al	$7.17 \cdot 10^5$ y 24	1808.7	99.76 4

 $^{27}Al(n,2n)^{26g}Al$.

Decay data used for

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 $^{28}\mathrm{Si}(\mathrm{n,\,p})^{28}\mathrm{Al}$



Decay	data used for	28 Si(n,p) 28 Al.
2		

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²⁸ Si	92.223 19	Si-crystal	²⁸ Al	2.245 m 2	1779.0	100

$^{29}Si(n, 2p)^{28}Mg$



No peaks were revealed in gamma spectra of irradiated silicon samples that could be identified as peaks belonging to 28 Mg. Therefore, only the upper limit of the 29 Si(n,2p) 28 Mg cross section was obtained that was set to be equal to three background uncertainties calculated for regions of the expected gamma peaks.

Decay	y data used for	²⁹ Si(n,2p)) ²⁸ Mg.	0+ 20.91 h 28 12 Μg β	<u>3+</u> 2.2414 28 13 β	m ▲
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					400.6	35.9 10
²⁹ c :	4 (95 0	Ci omustal	285	20.015 h o	941.7	36.3 10
²⁹ Si	4.685 8	SI-Crystal	IVIg	20.915 N 9	1342.2	54.0 16
					1779.0	100

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30 Si(n, α) 27 Mg

				1/2	2+ 9.458 m 27Mg β-	
Γ	Decay data used	1 for ³⁰ Si(1	$(n,\alpha)^{27}$ Mg.			5/2+ 27AI 13
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
³⁰ Si	3.092 11	Si-crystal	²⁷ Mg	9.458 m 12	843.8	71.80 2

1014.5

 $28.20\ 2$

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

 39 K(n, 2n) 38g K



	EC 0+ EC	130.4 0 38 19	923.9 ms 7.636 m
0+ 38Ar			

Decay data used for ${}^{39}K(n,2n)^{38g}K$.				38 Ar		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
³⁹ K	93.2581 44	KNO ₃	^{38g} K	7.636 m 18	2167.5	99.858 13

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



 41 K(n, α)³⁸Cl



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁴¹ K	6.7302 44	KNO ₃	³⁸ Cl	37.24 m 5	1642.4	33.3 7
					2167.5	44.4 9

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



 $^{41}\mathrm{K}(\mathrm{n,\,p})^{41}\mathrm{Ar}$



A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

$^{50}V(n,n\alpha)^{46}Sc$



Two samples enriched with V-50 were used. The total mass of V-50 in the samples was 6.94 mg.

Ι	Decay data used	for ⁵⁰ V(1	n,nα) ⁴⁶ Sc.		$\begin{array}{c} 1- & 142.528 \\ 1+ & 0 \\ 46 \\ 21 \\ \text{Sc} \\ \beta- \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ $	8.75 s 3.79 d <u>0+</u> <u>46</u> Ti
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵⁰ V	*17.0 1	V_2O_5	⁴⁶ Sc	83.79 d 4	889.3	99.984 <i>1</i>

* - refer to samples enriched with V-50

50 V(n, α) 47 Sc



The ⁵⁰V abundance in the natural vanadium is only 0.25%. Samples enriched with ⁵⁰V to 17% were purchased for this experiment. Gamma counting of the irradiated samples provided the determination of amount of ⁴⁷Sc nuclei but could not divide the contributions of ⁵⁰V(n, α)⁴⁷Sc and ⁵¹V(n,n' α)⁴⁷Sc reactions both of which were responsible for ⁴⁷Sc production. To resolve the problem, a second sample with a natural mixture of two vanadium isotopes was irradiated simultaneously. In more detail, the separation method of the ⁵⁰V(n, α)⁴⁷Sc and the ⁵¹V(n,n' α)⁴⁷Sc cross-sections is described on p.25.

	Decay data use	ed for 50	$V(\mathbf{n}, \alpha)^{47} \mathbf{Sc}.$		$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	5/2- 47 22
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵⁰ V	* 17.0 <i>1</i> (sep) 0.250 <i>4</i> (nat)	V ₂ O ₅ V-metal	⁴⁷ Sc	3.3492 d 6	159.4	68.3 4

-refer to samples enriched with V-50

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

51 V(n, n α) 47 Sc



The contribution of the ${}^{50}V(n,\alpha){}^{47}Sc$ reaction considered in the previous page was deduced by the method described on p.25.

	Decay data use	ed for 51,	V(n,na) ⁴⁷ Sc.		^{7/2-} 3.345 d 47 21SC β	<u>5/2-</u> 47 22
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵¹ V	99.750 4 (nat)	V-metal	⁴⁷ Sc	3.3492 d 6	159.4	68.3.4
	83.0 1 (sep)	V_2O_5			10,000	

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



 $^{51}V(n, \alpha)^{48}Sc$

					6+ 43.67 h 48Sc β-	
	Decay data us	sed for 51	/(n ,α) ⁴⁸ Sc.			<u>₀+</u> 48 22
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵¹ V	99.750 <i>4</i>	V-metal	⁴⁸ Sc	43.67 h 9	983.5 1037.5	100.1 <i>6</i> 97.6 7

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



 $^{55}Mn(n,\alpha)^{52}V$



	Decay data u		⁵² 24Cr			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵⁵ Mn	100	KMnO ₄	⁵² V	3.743 m 5	1434.1	100.0 14

46

${}^{55}Mn(n, 2n){}^{54}Mn$

$^{55}Mn(n,2n)^{54}Mn$						
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]			
13.56	622	4.04	0.597	2		
13.74	636	3.97	0.595	2		
13.96	661	3.97	0.567			
14.19	709	3.65	0.550	000		
14.42	740	3.63	0.554	C,		
14.61	764	3.37	0.569			
14.78	786	3.42	0.569			
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.1$						





0+

Decay data used for

sed for ${}^{55}Mn(n,2n){}^{54}Mn$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁵⁵ Mn	100	KMnO ₄	⁵⁴ Mn	312.2 d 2	834.8	99.976 1

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



54 Fe(n, p) 54 Mn



Decay data used for	ta used for
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁵⁴ Fe	5.845 35	Fe-metal	⁵⁴ Mn	312.2 d 2	834.8	99.976 1

⁵⁴Fe(n,p)⁵⁴Mn.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁵⁶Fe(n, p)⁵⁶Mn



ucleus	[%]	form	product	T _{1/2}	E _γ [keV]	Υ _γ [%]
					846.8	99.85 <i>3</i>
⁵⁶ Fe	91.754 36	Fe-metal	⁵⁶ Mn	2.5789 h 1	1810.7	26.9 4
					2113.1	14.2 3

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁵⁷Fe(n, p)⁵⁷Mn

The set-up for measurement of short-lived activities was used. (p.25).

					5/2- 87.2 s 57Mn β-	
	Decay data use	d ⁵⁷ Fe	(n,p) ⁵⁷ Mn.			^{1/2-} 57 Fe
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵⁷ Fe	2.119 10	Fe-metal	⁵⁷ Mn	85.4 s 18	122.1	14 5

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



59 Co(n, α) 56 Mn



Target	Abundance	Chemical	Reaction	T _{1/2}	E _γ	Υ _γ
Nucleus	[%]	form	product		[keV]	[%]
⁵⁹ Co	100	Co-metal	⁵⁶ Mn	2.5789 h 1	846.8 1810.7 2113.1	99.85 3 26.9 4 14.2 3

Decay	v data used for	59 Co(n. α) 56 Mm
D CCu		

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁵⁹Co(n, p)⁵⁹Fe



Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ	
nucleus	[%]	form	product		[keV]	[%]	
⁵⁹ Co	100	Co-metal	⁵⁹ Fe	44.495 d <i>9</i>	1099.2 1291.6	56.5 <i>18</i> 43.2 <i>14</i>	

⁵⁹Co(n,p)⁵⁹Fe.

⁵⁹Co(n, 2n)⁵⁸Co



70	82 d EC 27CO 82 d EC 27CO 9.15 h
<u>₀+</u> 58/76	-

Decay data used f	for
-------------------	-----

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ	
nucleus	[%]	form	product		[keV]	[%]	
⁵⁹ Co	100	Co-metal	⁵⁸ Co	70.86 d 6	810.8	99.450 10	

⁵⁹Co(n,2n)⁵⁸Co.



⁵⁹Co(n, 2n)^{58m}Co

The ⁵⁹Co(n,2n)^{58m}Co cross section was not measured directly but is the product of two other values measured immediately, the ⁵⁹Co(n,2n)⁵⁸Co cross section presented in the previous page and the isomeric ratio, i.e. $\sigma_m/(\sigma_m + \sigma_g)$, that was determined from the analysis of the 810.8 keV gamma decay curve. The method used for the isomeric ratio determination is described in detail earlier (pp. 20-21). The obtained isomeric ratio values were approximated by the straight line:

$$IR = a \cdot (E_n - E_0) + b$$

where $a = -0.005 \pm 0.009 \ MeV^1$
 $b = 0.654 \pm 0.012$
 $E_0 = 14.08 \ MeV$

The straight line parameters were determined by the least squares methods.





Decay data used for ${}^{59}Co(n,2n){}^{58m}Co$.

Target nucleus	Reaction product	T _{1/2}	Ε _γ [keV]	
⁵⁹ C a	^{58m} Co	9.10 h <i>9</i>	_	
Co	^{58g} Co	70.86 d 6	810.8	

${}^{58}Ni[(n, np) + (n, d)]{}^{57}Co$

58Ni(n x) ⁵⁷ Co				⁵⁸ Ni(n,x) ⁵⁷ Co
		,x) C0		
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]	
13.47	512	5.54	0.663	
13.56	540	4.89	0.626	
13.64	551	3.82	0.626	
13.74	574	2.59	0.624	CENDL-3.1 Present Data 10 '04 / Semkova+
13.88	565	5.57	0.596	Ø 300 - ★ 96 K.T.Osman+ Ø 94 Nusaki+ • 92 Melgaid+ • 92 Melgaid+
13.96	593	2.85	0.597	0 200 - 200
14.05	600	4.15	0.582	▲ '82 M.Viennot+ ◆ 77 Huang Janzhou+ ◆ 75 H.Weiget+
14.19	638	2.78	0.581	100 - (69 R.C.Barali+
14.28	638	4.02	0.579	
14.42	656	4.91	0.585	Neutron Energy [MeV]
14.47	649	4.74	0.585	
14.61	676	2.91	0.600	
14.68	660	4.99	0.600	
14.78	683	3.98	0.600	
14.86	689	3.92	0.622	
Ref. CS	is ⁹³ Nb(n,	2n) ^{92m} Nb;	$\alpha_d = 1.2$	

The data was corrected for the contribution of the ${}^{58}Ni(n,2n){}^{57}Ni$ cross section.

.....

		3/2- 35.60 h EC 28Ni
	7/2- 271.79 d ►C 27CO	
^{1/2-} 57 26 Fe		

Decay data used for		⁵⁸ Ni(n,x) ⁵⁷ Co.		⁵⁷ ₂₆ Fe		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵⁸ Ni	68 0760 00	Ni metal 5^{7} Co 271 74 d 6		122.1	85.60 17	
111	00.0709 89	INI-IIICtai	CO	271.74 u 0	136.5	10.68 8

⁵⁸Ni(n, p)⁵⁸Co

		500	'			
E _n [MeV]	σ [mb]	±Δ O _{tOtal} [%]	±Δ σ ref [%]		400	
13.47	424	6.54	0.640	[q		
13.56	418	3.19	0.601	<u>س</u>] ر	300	
13.64	399	4.92	0.602	ection		● P ◆ '1 ■ '0
13.74	388	4.04	0.600	s Se	200	- 0 0 □ '9 ◆ '9
13.88	379	6.57	0.570	Cros	200	● '9 ▼ '9 ▶ '8
13.96	364	4.61	0.571	0		• ⊗ '8 ⊕ '8 ⊕ '8
14.05	355	4.96	0.556		100	• 7 • 7 • 7 • 7
14.19	332	4.10	0.555			- ∢ '6 ▲ '6
14.28	342	5.74	0.553		0 1:	3.0 ·
14.42	316	4.33	0.559			
14.47	317	4.90	0.559			
14.61	294	4.67	0.574			
14.68	295	5.99	0.574			
14.78	279	4.22	0.574			
14.86	278	6.11	0.598			
Ref. CS	is ⁹³ Nb(n,2	$2n)^{92m}Nb;$	$\alpha_d = 1.2$			



	70.82 d E	5+ 2+ C	• 9.15 h •
<u>₀+</u> 58F	e		

⁵⁸Ni(n,p)⁵⁸Co.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁵⁸ Ni	68.0769 <i>89</i>	Ni-metal	⁵⁸ Co	70.86 d 6	810.8	99.450 10



⁵⁸Ni(n, p)^{58m}Co

The ⁵⁸Ni(n, p)^{58m}Co cross section was not measured directly but is the product of two other values measured immediately, the ⁵⁸Ni(n, p)⁵⁸Co cross section presented in the previous page and the isomeric ratio, i.e. $\sigma_m/(\sigma_m + \sigma_g)$, that was determined from the analysis of the 810.8 keV gamma decay curve. The method used for the isomeric ratio determination is described in detail earlier (pp. 20-21). The obtained isomeric ratio values were approximated by the straight line:

$$IR = a \cdot (E_n - E_0) + b$$

where $a = 0.0191 \pm 0.0037 \, MeV^1$
 $b = 0.550 \pm 0.005$
 $E_0 = 14.07 \, MeV$
The straight line parame

The straight line parameters were determined by the least squares methods.





Target nucleus	Reaction product	T _{1/2}	Ε _γ [keV]
⁵⁸ NI;	^{58m} Co	9.10 h 9	_
111	^{58g} Co	70.86 d 6	810.8

⁵⁸Ni(n, 2n)⁵⁷Ni

⁵⁸ Ni(n,2n) ⁵⁷ Ni				
E _n [MeV]	σ [mb]	±Δ σ _{total} [%]	±∆ 0 _{ref} [%]	
13.47	11.5	4.17	3.01	
13.56	12.6	4.56	3.00	
13.64	14.6	3.82	3.00	
13.74	15.4	4.35	3.00	
13.88	19.7	4.02	3.00	
13.96	20.2	4.12	3.00	
14.05	23.5	3.81	2.99	
14.19	25.1	5.89	2.99	
14.28	27.8	4.30	2.99	
14.42	29.8	4.59	3.00	
14.47	32.9	3.63	3.00	
14.61	33.4	3.65	3.00	
14.68	36.0	3.81	3.00	
14.78	35.9	3.81	3.00	
14.86	39.5	3.71	3.00	
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.2$				



		3/2- 35.60 h EC 28Ni
	7/2- 271.79 d EC 57 C0	
^{1/2−} 57 Fe		

Decay data used for ${}^{58}Ni(n,2n){}^{57}Ni$.			26			
Target Nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁵⁸ Ni	68.0769 89	Ni-metal	⁵⁷ Ni	35.60 h 6	127.2 1377.6 1919.5	16.7 5 81.7 24 12.3 4



⁶⁰Ni(n, p)⁶⁰Co



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
58NT:	26 2221 77	Ni motol	⁶⁰ C a	1025 28 4 14	1173.2	99.85 <i>3</i>
INI	20.2231 //	INI-metai	Co	1925.28 u 14	1332.5	99.9826 6

Decay data used for 60 Ni(n,p) 60 Co.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



63 Cu(n, α) 60 Co



Decay data used for	$^{63}Cu(n,\alpha)^{60}Co.$
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁶³ Cu	69.17 3	Cu-metal	⁶⁰ Co	1925.28 d 14	1173.2 1332.5	99.85 <i>3</i> 99.9826 6

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁶⁵Cu(n, p)⁶⁵Ni

5/2- 2.5172 65 28Ni	β
	^{3/2-} 65 29

Decay data used for ${}^{\circ}Cu(n,p){}^{\circ}Ni$.					29 ^{CU}	
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					366.3	4.81 6
⁶⁵ Cu	30.83 3	Cu-metal	⁶⁵ Ni	2.5175 h 5	1115.5	15.43 13
					1481.8	23.59 14

${}^{65}Cu(n, 2n){}^{64}Cu$



	$\underset{c_{1.0\%}}{\overset{1^+}{\underset{29}{\overset{12.700}{\overset{1}{\overset{1}{\overset{1}{}}}}}}}_{39.0\%}}{\overset{1^+}{\underset{39.0\%}{\overset{12.700}{\overset{1}{\overset{1}{}}}}}$	0⁺ 64 30Zn
⁶⁴ 28Ni		

Decay	data	used	for
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁶⁵ Cu	30.83 <i>3</i>	Cu-metal	⁶⁴ Cu	12.701 h 2	1345.8	0.475 11

 $^{65}Cu(n,2n)^{64}Cu.$
$^{64}Zn(n, 2n)^{63}Zn$





-	seedy adda abo					
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁶⁴ Zn	48.27 32	Zn-metal	⁵⁴ Mn	38.47 m 5	669.6 962.1	8.2 3 6.5 4

Decay data used for 64 Zn(n.2n) 63 Zn.

⁶⁸Zn(n, p)^{68m}Cu

	687(0				⁶⁸ Zn	(n,p) ^e	^{38m} Cu		
	Zn(n,	p) Cu				TENDL2	014 .0	I					
E _n [MeV]	0 [mb]	$\pm \Delta \sigma_{total}$ [%]	±Δ σ _{ref} [%]	-	7 –	JEFF3.2 A.A.Filat B.Lalrem Y.Kasug D.Kielan S.K.Gho A.Ercan	en Irua ai					-	
13.47	3.42	5.61	3.46	_		 M.Vienno R.Pepelr C.V.Srin 	ot nik iva					-	_
13.66	3.58	5.92	3.41	[up]	5 -	 J.L.Barre M.Vienne V.K.Tikk 	eir ot u				I		∮∮
13.88	3.77	7.23	3.40	ction	4					Ţ	, F	1	
14.04	4.32	5.47	3.40	Sec				÷	1				
14.26	4.59	9.69	3.39	coss				-		Ī			
14.44	4.54	6.11	3.39		2		/	-	4		a		
14.63	4.60	6.93	3.39	1	1								
14.81	4.38	6.30	3.39		ł								
Ref. (CS is ²⁷ Al(1	n, α) ²⁴ Na;	$\alpha_d = 1.8$		0 L 13.0	13.2	13	.4	13.6 1 Neu	3.8 tron [14.0 Energ	 14.2 y [N	 1eV]



4

14.8

15.0

1

14.6

4

14.4

Decay data used for	68 Zn(n,p) 68m Cu.
---------------------	------------------------------

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁶⁸ Zn	19.02 12	Zn-metal	^{68m} Cu	3.75 m 5	526.4	74.8 17

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁷⁴Ge(n, p)⁷⁴Ga



8.12 m 12

595.9

91.8 2

6	5
O	
	-

Target

nucleus

⁷⁴Ge

36.72 15

Ge-crystal

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



$^{74}Se(n, p)^{74}As$

	$\begin{array}{c} 2^{-} & 17.77 \text{ d} \\ \hline & & \\ EC \\ 66\% \end{array} \begin{array}{c} 74 \text{ As} \\ 33 \text{ As} \\ 34\% \end{array}$	0⁺ 74Se 34Se
^{0⁺} 32 Ge		

Υ_γ [%]

59 3

Deca	y data used for	Se(n, p) AS .	JL	52		
Target Nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]		
⁷⁴ Se	0.89 4	Se-crystal	⁷⁴ As	17.77 d 2	595.8		

becay data used for 74 Se(n, p) 74 A

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

⁷⁴Se(n, 2n)^{73m}Se



	EC 73 EC 34 EC 34
0.499 s <u>9/2⁺ ¥ 0</u>	3/2- 80.30 d
73 Ge	€C 73AS

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
	0.89 4	Se-crystal	^{73m} Se	39.8 m 13	84.0	2.03 19
⁷⁴ Se					253.7	2.36 19
					393.4	1.63 <i>13</i>

Decay data used for 74 Se(n, 2n) 73m Se.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

74 Se(n, 2n) 73g Se



	EC 34 ^{27.4%} 3/2- √ ^{√^{S[®]} 25.71} 39.8 m ≠ 9/2+ ▼ 0 FC 34 ^{Se} 7.15 h
0.499 s <u>1/2-∜ 66.716</u> 9/2+♥ 0 73 32 Ge	^{3/2-} 80.30 d EC 73As

						/
Target Nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁷⁴ Se	0.89 4	Se-crystal	^{73g} Se	7.15 h 8	361.2	97.0 10

Decay data used for 74 Se(n, 2n) 73g Se.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

⁷⁴Se(n, 2n)⁷³Se



This is the sum of two cross sections measured independently: 74 Se(n, 2n) 73m Se and 74 Se(n, 2n) 73g Se.



A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁷⁶Se(n, p)⁷⁶As

0+ 76 32Ge	2- 26.32 h EC 33AS β-	
		0⁺ 76 34Se

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁷⁶ Se	9.37 29	Se-crystal	⁷⁶ As	26.24 h 9	559.1 657.0	45 2 6.2 4

⁷⁶Se(n, p)⁷⁶As.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

⁷⁶Se(n, 2n)⁷⁵Se

76 Se(n, 2n) 75 Se					
E _n [MeV]	σ [mb]	±Δ O _{tOtal} [%]	±∆ 0 ref [%]		
13.49	625	6.21	3.27		
13.89	807	5.40	3.21		
14.05	736	12.1	3.20		
14.26	870	6.69	3.20		
14.44	844	10.6	3.19		
14.81	988	6.35	3.20		
Ref. C	CS is $^{27}Al(1)$	n, α) ²⁴ Na;	$\alpha_d = 1.2$		



Decay data used for ⁷⁶ Se(n, 2n) ⁷⁵ Se.			$16.79 \text{ ms} \xrightarrow{9/2^+ 303.9255}_{33}$			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					121.1	17.20 12
⁷⁶ Se	9.37 29	Se-crystal	⁷⁵ Se	119.78 d 5	136.0	58.5 4
					279.5	25.02 18

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

⁷⁷Se(n, p)⁷⁷As



<u>3/2</u> - 38.83 h	
77As ϡ-	$\frac{7/2^{+}}{161.9200}$ 17.36 s
	<u>1/2-</u> 0 77 Se

Decay data used for $Se(n, p)$ As.						
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁷⁷ Se	7.63 16	Se-crystal	⁷⁷ As	38.79 h <i>5</i>	239.0 249.8 520.7	1.59 24 0.39 6 0.56 9

Decay data used for 77 Se(n, p) 77 As.

78 Se(n, α) 75 Ge





Decay data used for ⁷⁸ Se

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁷⁸ Se	23.77 28	Se-crystal	⁷⁵ Ge	82.78 m 6	198.6	1.19 12

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



$^{78}\mathrm{Se}(\mathrm{n,\,p})^{78}\mathrm{As}$



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					613.8	54 6
⁷⁸ Se	23.77 28	Se-crystal	⁷⁸ As	90.7 m 2	694.9	16.7 22
					1308.7	13.0 18

Decay data used for 78 Se(n, p) 78 As.

$^{80}Se(n,\,\alpha)^{77(g+0.19m)}Ge$

80 Se(n, α) $^{77(g+0.19m)}$ Ge						
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±Δ σ _{ref} [%]			
13.49	2.07	5.80	1.48			
13.89	2.49	5.64	1.34			
14.05	2.26	8.80	1.33			
14.26	2.76	5.76	1.31			
14.44	2.71	7.42	1.30			
14.81	3.05	6.19	1.32			
Ref. 0	Ref. CS is ²⁷ Al(n, α) ²⁴ Na; $\alpha_d = 1.2$					





Decay	alla used for	Se(II, U)				
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					264.5	53.3 5
⁸⁰ Se	49.61 <i>41</i>	Se-crystal	⁷⁷ Ge	11 011 1 0	1.211 h <i>3</i> 367.5 416.4	14.5 7
				11.211 n 3		22.7 11
					557.9	16.8 10

Decay data used for 80 Se(n, α)⁷⁷Ge



⁸⁰Se(n, p)⁸⁰As



Decay data used for 80 Se(n, p) 80 As.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸⁰ Se	49.61 <i>41</i>	Se-crystal	⁸⁰ As	15.2 s 2	666.2	42 5

⁸⁰Se(n, 2n)^{79m}Se

80 Se(n, 2n) 79m Se						
$\begin{bmatrix} E_n & \sigma & \pm \Delta \sigma_{total} \\ [MeV] & [mb] & [\%] & [\%] \end{bmatrix} \xrightarrow{\pm \Delta \sigma_{re}}$						
14.05	213	7.30	1.02			
14.64	213	11.4	0.99			
14.84	242	9.77	1.01			
Ref. CS is ²⁷ Al(n, α) ²⁴ Na; $\alpha_d = 1.5$						



$3.91 \text{ m} \frac{1/2^{-4}}{7/2^{+4}} = \frac{95.73}{0}$	
<6.5×10 ⁴ y 79/34Se β-	3/2- 0
	⁷⁹ 35Br

Decay data	used	for
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 80 Se(n, 2n) 79m Se.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸⁰ Se	49.61 <i>41</i>	Se-crystal	^{79m} Se	3.92 m <i>1</i>	95.7	9.6192 11

⁸²Se(n, 2n)^{81m}Se





Decav	data	used
Doou	anna	

ed for 82 Se(n, 2n) 81m Se.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸² Se	8.73 22	Se-crystal	^{81m} Se	57.28 m 2	103.0	12.8 3

82 Se(n, 2n) 81g Se





Decay data used for 82 Se(n, 2n) 81g Se.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸² Se	8.73 22	Se-crystal	^{81g} Se	18.45 m <i>12</i>	290.0	0.56 6

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



⁸²Se(n, 2n)⁸¹Se

This is the sum of two cross sections measured independently: 82 Se(n, 2n) 81m Se and 82 Se(n, 2n) 81g Se.





$^{88}Sr(n,\alpha)^{85m}Kr$



	Decay da	ta used for	⁸⁸ Sr(n,	α) ^{85m} Kr.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸⁸ Sr	*99.94 3	SrCO ₃	^{85m} Kr	4.480 h 8	151.2	75.2 10

*- refer to samples with separated isotopes of Strontium

${}^{88}Sr(n, 2n){}^{87m}Sr$



4.7	3×10^{10} y $\frac{3/2^{-}}{37}$ B7 B β^{-}	³² 2.803 h
	9/2+ 87 c -	0

Decay data used for 88 Sr(n, 2n) 87m Sr				38 ⁷ Sr			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]	
⁸⁸ Sr	* 99.94 3	SrCO ₃	^{87m} Sr	2.815 h 12	388.5	82.19 22	

- refer to samples with separated isotopes of Strontium



89 Y(n, α) 86m Rb



Decay data used for	89 Y(n, α	86m Rb.
Decay data used 101	I (II, U	, 110

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸⁹ Y	100	Y-metal	^{86m} Rb	1.017 m <i>3</i>	556.1	98.20 6

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



89 Y(n, α) 86 Rb



Decay	data	used	for	
Ducay	uata	uscu	101	

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
Nucleus	[%]	form	product		[keV]	[%]
⁸⁹ Y	100	Y-metal	^{86m} Rb	18.642 d 18	1077.0	8.64 4

 89 Y(n, α) 86 Rb.



${}^{89}Y(n, 2n){}^{88}Y$



Decay data used for	89 Y(n, 2n) 88 Y.
	- (,)

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁸⁹ Y	100	Y-metal	⁸⁸ Y	106.626 d 21	898.0 1836.1	93.7 3 99.2 3



90 Zr(n, α) 87m Sr

4.75×10 ¹⁰ y ^{3/2−} 87 Rb	0.30% EC β_ B_ B_ 0.30% 1/2-√9 ⁸⁹ 388.532 2.803 h
57	9/2+ 0 87 Sr

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
Nucleus	[%]	form	product		[keV]	[%]
⁹⁰ Zr	51.45 40	Zr-metal	^{87m} Sr	2.815 h 12	388.5	82.3 5

Decay data used for 90 Zr(n, α) 87m Sr.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



 ${}^{90}Zr(n, p){}^{90m}Y$



Decay data used for	90 Zr(n, p) 90m Y.
---------------------	------------------------------

40 Zr-metal	^{90m} Y	3.19 h 6	202.5	97.3 <i>4</i>
	40 Zr-metal	40 Zr-metal ^{90m} Y	<i>40</i> Zr-metal ^{90m} Y 3.19 h <i>6</i>	40 Zr-metal 90mY 3.19 h 6 202.5 479.5



 90 Zr[(n, np) + (n, d)]^{89m}Y

The set-up modified for short-lived activity measurement was used (p.25).



Decay data used for

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
Nucleus	[%]	form	product		[keV]	[%]
⁹⁰ Zr	51.45 40	Zr-metal	^{90m} Y	15.663 s 5	909.0	99.16 3

 90 Zr(n, x) 89m Y

⁹⁰Zr(n,2n)^{89m}Zr 90 Zr(n, 2n) 89m Zr 200 TENDL-2014 ENDF/B-VIL1 EAF-2010 JEFF-3.2 2010 V Semkova+ 2008 Tran Due Thie 2008 Transformation 1975 R. Asliga+ 1972 V Kanda 1975 R. Rieder-1963 S.K. Mangal+ 1963 S.K. Mangal+ 1963 S.K. Mangal+ 180 E_n $\pm \Delta \sigma_{total}$ σ $\pm \Delta \sigma_{ref}$ • 160 [%] [MeV] [mb] [%] 140 13.47 30.6 1.186 2.63 Cross Section [mb] 40.1 3.04 1.024 13.66 1.007 13.88 62.9 2.53 0.991 14.04 77.2 2.68 0.970 14.26 91.8 2.88 14.44 116.4 2.43 0.956 40 14.63 124.6 2.51 0.959 20 14.81 141.1 2.42 0.979 0 13.4 13.6 13.8 14.0 14.6 14.8 15.0 13.2 14.2 14.4 13.0 Ref. CS is 27 Al(n, α) 24 Na; $\alpha_d = 1.5$ Neutron Energy [MeV]

90 Zr(n, 2n) 89m Zr

<u>9/2+∜ 908.96</u>	$EC_{ec}^{9/2+\sqrt{0}} \frac{9/2+\sqrt{0}}{40} $ 4.18 m FC_{ec}^{9/2+\sqrt{0}} \frac{9/2+\sqrt{0}}{40} 78.41 h 16.06 s
<u>1/2</u> − 0 89¥	

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁰ Zr	51.45 40 Zr-met	7 r-metal	^{89m} Zr	4 161 m 10	587.8	89.62 17
		Zi-metai		4.101 III <i>10</i>	1507.4	6.06 18

data used for	90 Zr(n, 2n) 89m Zr.
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Decay

90 Zr(n, 2n) 89 Zr





Decay	data	usec
Decuy	uuuu	abot

ed for 90 **Zr**(**n**, **2n**) 89 **Zr**.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹⁰ Zr	51.45 40	Zr-metal	⁸⁹ Zr	78.41 h <i>12</i>	909.1	99.04 <i>3</i>

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



 $^{92}Zr(n, p)^{92}Y$



Decay data used for

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹² Zr	1715 0	8 Zr-metal	⁹² Y	3.54 h 1	934.5	13.9 15
	17.13 0				1405.4	4.8 5

 92 Zr(n, p) 92 Y.

91

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



94 Zr(n, α) 91 Sr



${}^{96}Zr(n, 2n){}^{95}Zr$





for ${}^{96}Zr(n, 2n)^{95}Zr$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹⁶ Zr	2.80 9	Zr-metal	⁹⁵ Zr	64.032 d 6	724.2 756.7	44.27 22 54.38 22

 93 Nb(n, α) 90m Y $\pm \Delta \sigma_{ref}$ E_n σ $\pm \Delta \sigma_{total}$ [MeV] [%] [mb] [%] 13.47 5.03 4.38 2.71 13.56 4.81 7.18 2.70 5.15 2.70 13.66 4.49 13.74 5.15 5.89 2.70 13.88 5.28 5.40 2.69 5.35 13.96 6.56 2.69 5.68 5.22 2.69 14.05 14.19 5.68 5.88 2.69 5.61 14.27 4.18 2.69 14.42 5.58 4.55 2.69 14.47 5.85 5.07 2.69 14.61 5.80 5.93 2.69 5.74 14.68 6.10 2.69 14.78 5.81 6.59 2.69 14.86 5.84 4.12 2.70 Ref. CS is ${}^{93}Nb(n,2n)^{92m}Nb$; $\alpha_d = 1.4$



Neutron Energy [MeV]



15.0

Decay data used for Nb(n	, α) ^{20m} Y.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹³ Nb	100	Nb-metal	^{90m} Y	3.19 h 6	202.5 472.5	97.3 <i>4</i> 90.74 <i>5</i>

93 Nb(n, α) 90m Y

93 Nb(n, α) 90 Y						
E _n	σ	$\pm \Delta \sigma_{total}$				
[MeV]	[mb]	[%]				
13.47	12.2	7.45				
13.56	11.6	9.14				
13.66	12.4	6.96				
13.74	12.3	7.74				
13.88	12.6	7.15				
13.96	12.7	8.00				
14.05	13.5	6.89				
14.19	13.4	7.43				
14.27	13.2	6.25				
14.42	13.1	6.74				
14.47	13.7	7.20				
14.61	13.5	8.14				
14.68	13.3	8.45				
14.78	13.5	9.09				
14.86	13.5	7.79				
Ref.	Ref. CS is ${}^{93}Nb(n,2n)^{92m}Nb$					

 93 Nb(n, α) 90 Y

> g 8

6

⁹³Nb(n,α)⁹⁰Y





ENDF/B-VI EAF-2010 JENDL-4.0 CENDL-3.1 ROSFOND

Present Data 1988 R.Woelfle 1982 R.Fischer 1972 M.Bormar 1964 P.Kulisic+

14.8

15.0

1961 B.F 1960 H./

14.6

14.8 15.0

14.6

 93 Nb(n, α) 90 Y Decay data used for

Target nucleus	Reaction product	T _{1/2}	<mark>Ε_{β max}</mark> [keV]
⁹³ Nb	^{90m} Y	3.19 h 6	_
	^{90g} Y	64.00 h 21	2278.7

measured directly but is the product of two other values measured immediately, the ${}^{93}Nb(n, \alpha){}^{90m}Y$ cross section presented in the previous page and the isomeric ratio IR, i.e. $\sigma_m/(\sigma_m + \sigma_g)$, that was determined from the analysis of the time behavior spectrum of the continuous corresponded to the 90 Y β -particles that were registered by the HPGe detector with thin entrance window. The experiment is described in detail earlier (pp. 21 - 24). The obtained isomeric ratio values were approximated by the straight line in the used neutron energy interval:

 $IR(E_n) = a \cdot (E_n - E_0) + b$ where $a = 0.013 \pm 0.027 \, MeV^{1}$ $b = 0.422 \pm 0.019$ $E_0 = 14.07 \, MeV$

The straight line parameters were determined by the least squares methods.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



$^{92}Mo(n,\alpha)^{89m}Zr$



Decay data used	Decay data used	for
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Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹² Mo	14 77 31	Mo-metal	^{89m} 7r	4 161 m 10	587.8	89.62 17
MO	14.77 51		ZI	4.101 M 10	1507.4	6.06 18

⁹²Mo(n, α)^{89m}Zr.



$^{92}Mo(n, \alpha)^{89}Zr$



Decay data used	l
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for ${}^{92}Mo(n, \alpha)^{89}Zr$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹² Mo	14.77 <i>31</i>	Mo-metal	⁸⁹ Zr	78.41 h <i>12</i>	909.1	99.04 <i>3</i>

⁹²Mo(n, 2n)^{91m}Mo





Decay data used for $IVIO(II, 2II)$ $IVIO.$								
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]		
					652.9	48.2 21		
⁹² Mo	14 . 77 <i>31</i>	Mo-metal	^{91m} Mo	64.6 s 6	1208.1	18.6 <i>14</i>		
					1508.0	24.2 18		

Decay data used for ${}^{92}Mo(n 2n)^{91m}Mo$


$^{92}Mo[(n, np) + (n, d)]^{91m}Nb$

The intensity of the 104.6 keV γ -transition that discharges the metastable state $\frac{1}{2}$ of the ⁹¹Nb was measured in the experiment. The ^{91m}Nb half-life is 60.86 d. Additionally, this metastable state is popupated via EC decay of the ^{91m}Mo state with the probability of (0.50 ± 0.16) and half-life of 64.6 s. The ^{91m}Mo excited state is produced in the ⁹²Mo(n, 2n)^{91m}Mo reaction which cross section was measured independently. The corresponding data are presented in the previous page. The cross section data for the ⁹²Mo[(n,np)+(n,d)]^{91m}Nb presented in this page were obtained with the correction for the contribution of the ⁹²Mo(n, 2n)^{91m}Mo reaction into population of the metastable state ^{91m}Mo.

The decay of the 91 Mo ground state does not contribute to the 91m Nb 104.6 keV state excitation.



Decay data used for	⁹² Mo(n,	$x)^{91m}Nb.$
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹² Mo	14.77 31	Mo-metal	^{91m} Nb	60.86 d 22	104.6	0.574 9



92 Mo(n, p) 92m Nb



d for ${}^{92}Mo(n, p){}^{92m}Nb$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹² Mo	14.77 <i>31</i>	Mo-metal	^{92m} Nb	10.15 d 2	934.4	99.15 <i>4</i>

⁹⁴Mo(n, 2n)^{93m}Mo

$^{94}Mo(n, 2n)^{93m}Mo$					
E _n [MeV]	σ [mb]	±∆ T _{total} [%]	±Δ σ ref [%]		
13.64	1.59	9.29	2.20		
13.87	2.24	4.06	2.19		
14.05	2.98	9.59	2.19		
14.28	3.88	8.81	2.19		
14.44	4.60	8.92	2.19		
14.64	5.86	10.1	2.19		
Ref. CS	is ⁹³ Nb(n,2	2n) ^{92m} Nb;	$\alpha_d = 1.6$		





Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					263.0	57.4 11
⁹⁴ Mo	9.23 10	Mo-metal	^{93m} Mo	6.85 h 7	684.7	99.9 8
					1477.1	99.1 <i>11</i>

Decay data used for	$^{94}Mo(n, 2n)^{93m}M$	0.
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 $^{nat}Mo(n, x)^{95}Nb$

The samples made of natural molybdenum were used. Several reactions can contribute to the production of ⁹⁵Nb:

- 95 Mo(n, p) 95 Nb is the most probable reaction, >80%;
- ${}^{96}Mo(n, x){}^{95}Nb$ the possible contribution is changing, approximately, from 1 to 15%; the evaluations are very scattered (see Figure to the right where the ratio of ${}^{96}Mo(n, x){}^{95}Nb$ cross-section to this of ${}^{95}Mo(n, p){}^{95}Nb$ is shown);
- $^{98}Mo(n, \alpha)^{95}Zr \rightarrow ^{95}Nb$ the contribution is small due to the long half-live of 95 Zr.



14.8

95 42 Mo

15.0



Decay	data	used	for
	uata	uscu	101

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁵ Mo	15.84 11	Mo motol	⁹⁵ NIL	34 001 d <	765 9	00 000 7
⁹⁶ Mo	16.67 15	Mo-metal	IND	34.991 0 0	/05.8	99.008 /

 $^{nat}Mo(n, x)^{95}Nb.$



$^{nat}Mo(n, x)^{95m}Nb$

The samples made of natural molybdenum were used. Several reactions can contribute to the production of 95m Nb:

- $^{95}Mo(n, p)^{95m}Nb$ is the most probable reaction, >80%;
- ⁹⁶Mo(n, x)^{95m}Nb − the possible contribution is changing, approximately, from 2 to 20%; the evaluations are very scattered (see Figure to the right where the ratio of ⁹⁶Mo(n, x)^{95m}Nb cross-section to this of ⁹⁵Mo(n, p)^{95m}Nb is shown);
- 98 Mo(n, α) 95 Zr → 95m Nb could be neglected since only 1.08% of 95 Zr decays to 95m Nb. Besides, the gamma counting was done in the first days after irradiation.





Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁵ Mo	15.84 11	Ma matal	95m N Th	26142	225 7	25.0.8
⁹⁶ Mo	16.67 15	Mo-metai	IND	3.01 U 3	255.1	25.0 8

 $^{nat}Mo(n, x)^{95m}Nb.$

Decay data used for



 $^{nat}Mo(n, x)^{96}Nb$

The samples made of natural molybdenum were used. Several reactions can contribute to the production of 96 Nb:

- ${}^{96}Mo(n, p){}^{96}Nb$ is the most probable reaction, >90%;
- $^{97}Mo(n, x)^{96}Nb$ the possible contribution is changing, approximately, from 1 to 8%; the evaluations are scattered (see Figure to the right where the ratio of $^{97}Mo(n, x)^{96}Nb$ cross-section to this of $^{96}Mo(n, p)^{96}Nb$ is shown).





	Decay data	used for	^{nat} Mo(n, x) ⁹⁶ N	Nb.		⁹⁶ 42Mo
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁶ Mo	16.67 15	Mo motol	⁹⁶ NIL	22.25 h 5	568.9	58.0 3
⁹⁷ Mo	9.60 14	wio-metai	ND	25.55 11 5	1200.2	19.97 10



$^{nat}Mo(n, x)^{97}Nb$					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]		
13.48	1.42	12.0	1.59		
13.64	1.44	18.0	1.57		
13.87	1.51	15.2	1.56		
14.05	1.62	3.97	1.56		
14.28	1.70	6.63	1.56		
14.45	1.74	10.9	1.56		
14.64	1.86	4.98	1.56		
14.82	2.02	5.09	1.57		
Ref. 0	CS is $^{27}Al(1)$	n, α) ²⁴ Na;	$\alpha_d = 1.5$		



The samples made of natural molybdenum were used. Several reactions can contribute to the production of 97 Nb:

- 97 Mo(n, p) 97 Nb is the most probable reaction, >75%;
- 98 Mo(n, x)⁹⁷Nb the possible contribution is changing with neutron energy, approximately, from 1 to 25%; the evaluations are very scattered (see Figure to the right where the ratio of 98 Mo(n, x)⁹⁷Nb cross-section to this of 97 Mo(n, p)⁹⁷Nb is shown);
- 97 Mo(n, p) 97 Nb is shown); 100 Mo(n, α) 97 Zr → 97 Nb – the contribution is small due to the half-life of 97 Zr which is much more than the half-life of 97 Nb.





97 42 Mo

Decay data used for	$^{nat}Mo(n, x)^{97}Nb.$
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Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁷ Mo	9.60 14	Mo-metal	Mo motol ⁹⁷ Nih	70.1 m 7	(57.0	08 23 0
⁹⁸ Mo	24.39 37		IND	/ <i>2</i> .1 III /	657.9	90.25 8

$^{98}Mo(n,\alpha)^{95}Zr$

98 Mo(n, α) 95 Zr				
E _n [MeV]	σ [mb]	±Δ σ _{total} [%]	±∆ 0 _{ref} [%]	
13.47	4.97	6.01	1.31	
13.64	5.07	6.92	1.29	
13.88	5.30	8.52	1.28	
14.05	5.94	5.36	1.27	
14.28	5.87	7.64	1.27	
14.47	6.55	5.44	1.27	
14.68	6.34	4.83	1.28	
14.86	6.69	4.95	1.29	
Ref. CS	is ⁹³ Nb(n,	2n) ^{92m} Nb;	$\alpha_d = 1.2$	



$$\frac{5/2^{+} \ 64.02 \ d}{95} \frac{1}{20} \frac{1}{20} \frac{1}{20} \frac{5.6\%}{\beta} = \frac{5.6\%}{34.975 \ d} \frac{1}{20} \frac{1}$$

Decay	data	used	for
-------	------	------	-----

for
$${}^{98}Mo(n, \alpha){}^{95}Zr$$
.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹⁸ Mo	24.19 26	Mo-metal	⁹⁵ Zr	64.032 d 6	724.2 756.7	44.27 22 54.38 22

⁹⁸Mo(n, p)^{98m}Nb

$^{98}Mo(n, p)^{98m}Nb$				
E _n [MeV]	σ [mb]	±∆ T _{total} [%]	±Δ σ ref [%]	
13.48	3.48	14.2	1.84	
13.64	4.01	15.0	1.74	
13.87	4.12	8.71	1.73	
14.05	4.14	5.41	1.72	
14.28	4.34	8.66	1.71	
14.45	4.78	10.6	1.70	
14.64	4.58	6.26	1.71	
14.82	4.75	8.71	1.72	
Ref. 0	CS is $^{27}Al(1)$	n, α) ²⁴ Na;	$\alpha_d = 1.6$	





Decay data used for ${}^{98}Mo(n, p){}^{98m}Nb$.

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁸ Mo	24.19 26	Mo-metal	^{98m} Nb	51.3 m <i>4</i>	722.6	73.8 15
					787.4	93.4 2

$^{100}Mo(n,\alpha)^{97}Zr$

100 Mo(n, α) 97 Zr					
E _n [MeV]	σ [mb]	±∆ O _{total} [%]	±∆ σ ref [%]		
13.64	2.21	5.08	2.16		
14.28	2.51	5.60	2.15		
14.44	2.74	5.48	2.15		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.7$					





Decay	data	used
2000		

l for ${}^{100}Mo(n, \alpha)^{97}Zr.$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰⁰ Mo	9.67 20	Mo-metal	⁹⁷ Zr	16.749 h <i>8</i>	743.4	93.09 16



¹⁰⁰Mo(n, 2n)⁹⁹Mo

Formally, the ¹⁰⁰Mo(n, x)⁹⁹Nb is also contributed to production of ⁹⁹Mo because the ⁹⁹Nb has a short half-life. The cross-section of the ¹⁰⁰Mo(n, x)⁹⁹Nb reaction is evaluated as 0.8 mb approximately that is much less than the cross section of the studied ¹⁰⁰Mo(n, 2n)⁹⁹Mo reaction. Therefore, the contribution of the ¹⁰⁰Mo(n, x)⁹⁹Nb was ignored.



Decay	data	used	1
Decav	uala	useu	1

for ${}^{100}Mo(n, 2n)^{99}Mo$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰⁰ Mo	9.67 20	Mo-metal	⁹⁹ Mo	65.976 h 24	181.1 739.5	6.14 <i>12</i> 12.26 22

99 Tc(n, n α) 95 Nb



In the experiment, ⁹⁵Nb ground state decay was measured. The ground state is also populated by an isomeric transition from a metastable state with the probability IT = 94.4% (see the decay scheme lower). The cooling time was ~ 20 days, and about 5.6% of the ^{95m}Nb excitation fell out of the measured ⁹⁵Nb total excitation. Let us estimate the lost value.

Isomeric ratio (or brunching ratio that is just the same) is defined as: $IR = \sigma_m/(\sigma_m + \sigma_g)$ It is estimated in evaluations as:

 $IR(JEFF) \sim 0.07;$ $IR(EAF) \sim 0.06;$ $IR(TENDL) \sim 0.04$

Assume that IR = 0.05.

The loss of the ground state population is then $0.05*0.056 \approx 0.003$ of the total value that was used for the ⁹⁹Tc(n, n\alpha)⁹⁵Nb cross-section determination. This was neglected.

In addition, the samples may contain ⁹⁸Tc with admissible level no more than 0.0008. The ⁹⁸Tc(n, α)⁹⁵Nb cross-sections are evaluated in available evaluations as (12 – 22) mb that are resulted in possible contributions of about (0.010 - 0.018) mb. The estimated contributions were not subtracted from the data presented but were added to the uncertainties.



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
⁹⁹ Tc	99.92 2	Tc-metal	To motol ⁹⁵ Nh	24 001 J <	765 8	00 808 7
⁹⁸ Tc	0.08 2		IND	34.991 a o	/05.8	99.808 /

Decay data used for 99 Tc(n, n α) 95 Nb.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



$^{99}Tc(n, \alpha)^{96}Nb$



1200.2

Yγ [%] 58.0 3

19.97 10

Decay data used for 99 Tc(n, α) 96 Nb.						
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	
⁹⁹ T a	00.02.2	To motol	⁹⁶ Nth	22.25 h 5	568.9	
10	<u> </u>	i c-metal	IND	23.35 fl 3	1000 0	

 99 Tc(n, α) 96 Nb.

⁹⁹Tc(n, 2p)^{98m}Nb



It was questionable the possibility of experimental measurement of an extremely small cross-section expected for (n,2p) reaction. Although samples were irradiated and counted in very close geometries no statistically significant peaks were revealed in the spectra. The upper limit of the cross section value was set to the 3 statistical uncertainties.



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⁹⁹Tc(n, p)⁹⁹Mo

	42	μο β ⁻ <u>1// 9//</u>		$\begin{array}{c} 0.0037\% \\ \beta^{-} 6.01 \text{ h} \\ 2.111 \times 10^5 \text{ y} \\ \beta^{-} \\ \underline{5/2^{+}} \\ 99 \\ 44 \\ Ru \end{array}$
1/:	2+ (65.94 h	06301	° .

Decay data u	sed for
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
⁹⁹ Tc	99.92 <i>2</i>	Tc-metal	⁹⁹ Mo	65.976 h 24	181.1 739.5	6.14 <i>12</i> 12.26 22

 99 Tc(n, p) 99 Mo.

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103 Rh(n, n α) 99m Tc



Decay data used for	103 Rh(n, n α) 99m Tc
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰³ Rh	100	Rh-metal	^{99m} Tc	6.0067 h 5	140.5	89 4



103 Rh(n, p) 103 Ru

$\frac{11/2}{2/31}$ $\frac{238.2}{1.69}$ 1.69 ms
39 26 d
103
ARU R
44 p-
4
7/2+ 39 756
$\frac{112}{56.12}$ 56.12 m
1/2⁻▼ 0
400
103 Dh
45
40

Decay	data	used	for
-------	------	------	-----

 103 Rh(n, p) 103 Ru.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰³ Rh	100	Rh-metal	¹⁰³ Ru	39.247 d 13	497.1	91.0 <i>12</i>



 103 Rh(n, 2n) 102m Rh

The old data were corrected for the change of the half-life used. The old value was 2.9 y, the new one is 3.742 y.



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					697.5	44.0 20
103 Rh	100	Rh-metal	^{102m} Rh	3.742 y 10	766.8	34.0 20
					1112.8	19.0 10

 103 Rh(n, 2n) 102m Rh.

Decay data used for

103 Rh(n, 2n) 102g Rh





Decer	data	usad	for
Decay	data	usea	IOL

ed for 103 Rh(n, 2n) 102g Rh.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰³ Rh	100	Rh-metal	^{102g} Rh	207.3 d 17	468.6	2.90 20

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103 Rh(n, 2n) 102 Rh

The present experimental data for 103 Rh(n, 2n) 102 Rh cross-section are the sum of the 103 Rh(n, 2n) 102m Rh and 103 Rh(n, 2n) 102g Rh cross-sections measured independently and presented on the previous pages. This was done for extending the comparative base becase most evaluations contain the data for the reaction cross-sections without splitting between ground and metastable state excitations.



102 Pd(n, 2n) 101 Pd



$\overset{(5/2^+)}{=} \begin{array}{c} 8.47 \\ h \\ EC \\ 46 \\ \hline 101 \\ EC \\ 46 \\ \hline 101 \\ 1/2^- \hline 0 \\ 101 \\ \hline 101 \\ 101 \\ \hline 101 \\ 101 \\ \hline 101 \\ 3.3 \\ y \\ \hline 101 \\ 101 \\ \hline 101 \\ 3.3 \\ y \\ \hline 101 \\ \hline $
/2+
¹⁰¹ ₄₄ Ru

Decay	data used for	102 Pd(n, 2n	$1)^{101}$ Pd.	44170		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁰² Pd	1.02 1	Pd-metal	¹⁰¹ Pd	8.47 h 6	269.7 296.3	6.4 <i>3</i> 19.2 <i>8</i>
					390.4	12.1 5

an data waad far	102 D J(2
av data used for	¹ ² Pd(n.	2n)***Pd



102 Pd[(n, np) + (n, d)]^{101m}Rh

The data was corrected for decay of the 101 Pd produced in the 102 Pd(n, 2n) 101 Pd reaction.



 102 Pd(n, x) 101m Rh.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰² Pd	1.02 1	Pd-metal	^{101m} Rh	4.34 d 1	306.9	81 4

102 Pd[(n, np) + (n, d)]^{101g}Rh



The data were corrected for decay of the 101 Pd produced in the 102 Pd(n, 2n) 101 Pd reaction.

(5/2+) 8.47 h	
[▲] 101 FC 46Pd	
92% or the second secon	92%
$EC_{1/2} = \frac{9/2+(+157.32)}{1/2} 4.34 d$	EC
101ph 3.3 y	
EC 45KN	,
2+	5/0+
2)/Z ⁻
101p.,	101 D
44 KU	44 RU

Deca	y data used for	102 Pd(n, y	$\mathbf{x})^{101g}\mathbf{Rh}.$	¹⁰¹ 44Ru		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					127.2	68 6
102 Pd	1.02 1	Pd-metal	101g Rh	3.3 y 3	198.0	73 6
					325.2	11.8 10



$^{102}Pd[(n, np) + (n, d)]^{101}Rh$

This is the sum of the ${}^{102}Pd(n, x){}^{102m}Rh + {}^{102}Pd(n, x){}^{102g}Rh$ cross sections measured independently and presented on the previous pages.



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105 Pd(n, p) 105 Rh

^{3/2+} 4.4	<u>l4 h</u> U _β –			
		1/2- √129.78 7/2+▼ (0 105 45 Rh	45 s 35.36 β-	6 h
			5	¹⁰⁵ 46Pd
		Eγ		Yγ

Decay uata useu 101	Decay	data	used	for
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 105 Pd(n, p) 105 Rh.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰⁵ Pd	22.33 8	Pd-metal	¹⁰⁵ Rh	35.36 h 6	306.1 318.9	5.1 3 19.1 6

$^{106}Pd(n,\alpha)^{103}Ru$





Decay data used for ¹⁰⁶ Pc	$l(n, \alpha)^{103}$ Ru.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰⁶ Pd	27.33 3	Pd-metal	¹⁰³ Ru	39.247 d 13	497.1	91.0 <i>12</i>



106 Pd(n, p) 106m Rh



	Decay data	used for	106 Pd(n, p) 1061	ⁿ Rh.		¹⁰⁶ ₄₆ Pd
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁰⁶ Pd			^{106m} Rh		450.8	24.2 13
	27.22.2	Dd matal		121 2	616.1	20.2 14
	21.33 3	Pu-metai		131 m 2	717.2	28.9 16
				1046.7	30.4 16	

$^{108}Pd(n,\alpha)^{105}Ru$

	¹⁰⁸ Pd(n,	α) ¹⁰⁵ Ru		¹⁰⁸ Pd(n,α) ¹⁰⁵ Ru
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]	4.0 TENDL-2014 ENDF/8-VIL1 ENDF/8-VIL1 JENDL-4.0 CENDL-31 3.5 Present Data
13.46	1.94	6.22	3.31	
13.64	2.37	13.1	3.30	
13.88	2.40	13.4	3.30	ction
14.06	2.35	6.01	3.29	
14.28	2.88	13.2	3.29	SO 1.5
14.48	3.03	7.34	3.29	1.0
14.68	3.26	9.44	3.30	0.5
14.87	3.74	12.9	3.30	
Ref. CS	is ⁹³ Nb(n,	2n) ^{92m} Nb;	$\alpha_d = 1.2$	13.0 13.2 13.4 13.6 13.8 14.0 Neutron Energy



<u>3/2+</u> 4.44 105 44	h β-	
	1/2- 129.78 7/2+♥ 105Rh 45Rh	1 45 s 35.36 h β-
		^{5/2+} 105Pd
	F	v

Decay data used for

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁰⁸ Pd	26.46 9	Pd-metal	¹⁰⁵ Ru	4.44 h 2	469.4 676.4	17.5 6 15.7 5

 108 Pd(n, α) 105 Ru.

110 Pd(n, 2n) 109 Pd



11/2-188.99	⁰ 4 696 m
5/2+▼ 109Pd	13.7012 h
46. 6	$39.6 \text{ s} \frac{7/2+88.0341}{12}$
	^{1/2-} (109 47 47

Decay data us	ed for
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used for 110 Pd(n, 2n) 109 Pd.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹¹⁰ Pd	11.72 9	Pd-metal	¹⁰⁹ Pd	13.7012 h 24	88.0	3.6 4

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 $^{107}\mathrm{Ag}(\mathrm{n},\,2\mathrm{n})^{106\mathrm{m}}\mathrm{Ag}$

EC, 6+	89.66 8.	28 d
23.96 m 1 EC	06 47 Ασ β-	0 ⁺ 0 106Cd
99.5%	<1%	40
0+		
¹⁰⁶ ₄₆ Pd		

Decay data used for	107 Ag(n, 2n) 106m Ag.
---------------------	----------------------------------

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁰⁷ Ag	51.839 8	Ag-metal	^{106m} Ag		451.0	28.2 7
				6 1 9 1 9	616.2	21.6 6
				8.28 U 2	717.3	28.9 8
					1045.8	29.6 10

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

$^{109}Ag(n, 2n)^{108m}Ag$





Decay data used for $^{109}Ag(n, 2n)^{108m}Ag$.			¹⁰⁰ 46Pd			
Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					433.9	90.5 6
¹⁰⁹ Ag * 99.4 <i>1</i>	* 99.4 1	Ag-metal	^{108m} Ag	438 y 9	614.3	89.8 19
					722.9	90.8 19

- refer to samples enriched by Silver-109

$^{106}Cd(n, x)^{105}Ag + ^{106}Cd(n, 2n)^{105}Cd$



The ¹⁰⁵Ag production cross section is presented in this page. It was measured using the ¹⁰⁵Ag γ -rays given in the Table at the page bottom. Since the ¹⁰⁵Ag half live is 41.29 d, the irradiation, cooling and gamma counting were comparatively long at more than several hours. Under these conditions not only the ¹⁰⁶Cd(n, x)¹⁰⁵Ag cross-section is measured (here (n, x) is the sum of (n, d) and (n, np) reactions) but also ¹⁰⁶Cd(n, 2n)¹⁰⁵Cd is added to the measurement result because the ¹⁰⁵Cd half-life is 55.5 min and practically all the ¹⁰⁵Cd nuclei decay into the ¹⁰⁵Ag after cooling time more than 5 hours.

The ${}^{106}Cd(n, 2n){}^{105}Cd$ was separately measured in short irradiations that were accompanied by short cooling time and short gamma countings of less than 1 h. The data obtained for the ${}^{106}Cd(n, 2n){}^{105}Cd$ cross section are presented on the next page.

To reduce errors, the ¹⁰⁶Cd(n, 2n)¹⁰⁵Cd cross section was approximated by a polynom of the second order before subtracting it from the total ¹⁰⁵Ag production cross section to determine the ¹⁰⁶Cd(n, x)¹⁰⁵Ag cross section shown two pages later.

In this page data Table, the total ¹⁰⁵Ag production cross section is presented that was measured under the conditions described above. The experimental data agree well with the corresponding cross section sums taken from evaluations. At least, this is valid for three of five available evaluations.

	0.34% EC <u>7/2+</u> 1/2-▼ 0 EC 105Ag	5/2* 55.5 m EC 48Cd 7.23 m 41.29 d
^{5/2+} 105 46Pd		

Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁰⁶ Cd	1.25 6	Cd-metal	¹⁰⁵ Ag	41.29 d 7	280.4	30.2 17
					644.5	11.1 6
					650.7	2.54 4
					1087.9	3.85 17

Decay data of ¹⁰⁵Ag

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¹⁰⁶Cd(n,2n)¹⁰⁵Cd 106 Cd(n, 2n) 105 Cd 1000 900 $\pm \Delta \sigma_{ref}$ E_n σ $\pm \Delta \sigma_{total}$ [MeV] [mb] [%] [%] 800 9.22 13.47 398 8.01 700 Cross Section [mb] 13.64 441 10.0 8.01 600 457 9.18 8.01 13.88 500 14.05 485 8.21 8.01 400 14.28 478 10.0 8.01 300 TENDL-2014 ENDF/B-VII.1 EAF-2010 JEFF-3.2 8.01 14.47 545 9.82 8.01 200 14.68 521 8.10 JENDL-4.0 Present Data 1993 C.Konr 1969 W.Lu+ 1963 L.A.Ra 14.81 543 8.14 8.01 100 14.86 559 10.4 8.01 0 13.6 13.8 14.0 14.2 13.0 13.2 13.4 14.4 14.6 14.8 15.0 Ref. CS is ${}^{93}Nb(n,2n){}^{92m}Nb;$ $\alpha_d = 1.5$ Neutron Energy [MeV]

$^{106}Cd(n, 2n)^{105}Cd$

The present data on ${}^{106}Cd(n, 2n){}^{105}Cd$ cross sections agree well with the latest experimental data [22] obtained for this reaction in JAERI. However, the old experiments carried out in sixties of the past century gave the results that are almost two times higher, and most evaluations follow the old data.

It seems desirable to pay more attention to the recent experimental data that were performed independently but are in a good agreement.

5/2⁺ 55.5 m C 48Cd

⁵⁵ 7.23 m <u>0</u> 41.29 d

Ag

Decay data used for ${}^{106}Cd(n, 2n){}^{105}Cd$.				5/2+ 103	Pd	
Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					648.5	1.57 11
¹⁰⁶ Cd	1.25 6	Cd-metal	¹⁰⁵ Cd	55.5 m <i>4</i>	1302.5	4.7 <i>3</i> 4.1 <i>3</i>
					1693.3	3.54 24



106 Cd[(n, np)+(n,d)] 105 Ag

The present data are the result of subtracting the 106 Cd(n, 2n) 105 Cd cross-sections given in the previous page from the 105 Ag production cross-section discussed two pages earlier. The data are shown in the Table and in the Figure above (olive circles in the Figure).

In the same way, the data of Konno at al. [22] were obtained (magenta triangles in the Figure). The disagreement between these two data sets is obvious.

A possible cause of the observed cross section difference may be the difference in the reference decay data used. So, the γ -ray 443.00 keV with the intensity 17.1% was taken in [22]. According to the recent evaluations [26], the intensity 10.5% is recommended for the gamma-line 443.37 keV. If the gamma intensity would be changed to the recently recommended value then the agreement between the two cross-section data sets will be much better.

In our experiment, four γ -rays were used, and the cross-section data calculated for any one of the γ -lines agrees within errors with the results for any other γ -ray. The γ -line with the energy about 443 kev was not included in the list of the analysed γ -rays.



De	ecay data used	for CS	106 Cd(n, x) 106	⁵ Ag.	^{5/2+} 46Pd	
Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁰⁶ Cd	1.25 6	Cd-metal	¹⁰⁵ Ag	41.29 d 7	280.4	30.2 17
					644.5	11.1 6
					650.7	2.54 4
					1087.9	3.85 17

$^{106}Cd(n,p)^{106m}Ag$

$^{106}Cd(n, p)^{106m}Ag$					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ σ ref [%]		
13.45	46.3	4.85	5.45		
13.88	48.0	5.13	5.44		
14.44	44.9	6.43	5.44		
14.64	44.7	8.78	5.45		
14.88	43.1	5.82	5.45		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.2$					



l 23.96 r	$ \begin{array}{c} $	^{89.66} 0 8.2 6 6 7 Αg β- <1%	^{28 d} 0 106Cd	
0 ⁺				

Decay data used for $^{106}Cd(n, p)^{106m}Ag$.			¹⁰⁶ ₄₆ Pd			
Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁰⁶ Cd	1.25 6	Cd-metal	^{106m} Ag	9 29 4 2	429.6	13.2 4
					451.0	28.2 7
				8.28 U 2	824.7	15.3 4
					1527.6	16.3 <i>13</i>

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

$^{108}Cd(n, 2n)^{107}Cd$

108 Cd(n, 2n) 107 Cd					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]		
13.45	1107	16.8	7.26		
13.88	1317	20.5	7.25		
14.44	1164	7.29	7.25	1	
14.64	1493	7.45	7.25		
14.81	1372	5.40	7.25		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.2$					



	5/2 ⁺ 6.50 h EC 48Cd
7/2+ ≮93.13 1/2- ▼ 0	44.3 s
¹⁰⁷ ₄₇ Ag	

Decay data used for	108 Cd(n, 2	$(2n)^{107}$ Cd.
---------------------	------------------	------------------

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	Form	product		[keV]	[%]
¹⁰⁸ Cd	0.89 3	Cd-metal	¹⁰⁷ Cd	6.50 h 2	93.1	4.7 3


$^{110}Cd(n,p)^{110m}Ag$

0+	$\begin{array}{c} 6^{+} & (117.59) \\ 6^{+} & (117.59) \\ 1^{+} & 0$
¹¹⁰ ₄₆ Pd	EC 47~9 β– 0.30% 99.70%

Decay data used for $Ca(n, p)$ Ag.						
Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹¹⁰ Cd	12.49 18	Cd-metal	110m A ~	249.83 d 4	657.8	95.61 9
					884.7	75.0 11
			Ag		937.5	35.0 3
					1384.3	25.1 5

Decay data used for 110 Cd(n n)^{110m}Ag

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

$^{110}Cd(n, 2n)^{109}Cd$

110 Cd(n, 2n) 109 Cd					1800	
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]		1600 1400	
13.45	1319	13.6	3.38	ā	1200	
13.88	1485	18.5	3.37	u [u	1000	•
14.44	1527	10.2	3.36	ectio	800	-
14.64	1598	8.59	3.37	ss S	000	-
Ref. CS	is ⁹³ Nb(n,	2n) ^{92m} Nb;	$\alpha_d = 1.2$	CC	400	-
					200	-



39.6 s 7/2+ 88.0341	5/2+ 462.6 d
<u>1/2</u> ↓ 0	EC 48CO
¹⁰⁹ Ag	

Decay	data	usec

d for ${}^{110}Cd(n, 2n){}^{109}Cd$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	Form	product		[keV]	[%]
¹¹⁰ Cd	12.49 18	Cd-metal	¹⁰⁹ Cd	461.4 d <i>12</i>	88.0	3.70 11



¹¹¹Cd(n, p)¹¹¹Ag

$\begin{array}{c} & & & & & & \\ & & & & & \\ \hline & & & & & \\ \hline & & & &$	64.8 s 7.45 d	
	11/2- √396.22 1/2+ ♥ 0 111 Cd 48 Cd	48.54 m

$^{111}Cd(n, p)^{111}Ag.$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	Form	product		[keV]	[%]
¹¹¹ Cd	12.80 12	Cd-metal	¹¹¹ Ag	7.45 d 1	245.4 342.1	1.24 9 6.7 3

$^{112}Cd(n,\alpha)^{109}Pd$

112 Cd(n, α) 109 Pd					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±Δ σ _{ref} [%]		
13.45	3.40	35.3	11.2		
13.88	2.34	29.3	11.2		
14.44	3.80	23.7	11.2		
14.64	3.81	26.3	11.2		
14.88	4.86	27.1	11.2		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.2$					





Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	Form	product		[keV]	[%]
¹¹² Cd	24.13 21	Cd-metal	¹⁰⁹ Pd	13.7012 h 24	88.0	3.6 4

Decay data used for $^{112}Cd(n, \alpha)^{109}Pd$.



$^{112}Cd(n, p)^{112}Ag$



3.130 h 8

617.5

43 5

Decay data used for

Cd-metal

Target

nucleus

 ^{112}Cd

24.13 21

¹¹²Ag



$^{112}Cd(n, 2n)^{111m}Cd$



Decay	data	used
Decav	uala	useu

d for
$${}^{112}Cd(n, 2n)^{111m}Cd$$
.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	Form	product		[keV]	[%]
¹¹² Cd	24.13 21	Cd-metal	^{111m} Cd	48.50 m <i>9</i>	150.8 245.4	29.1 <i>18</i> 94 7

¹¹³Cd(n, p)¹¹³Ag

	¹¹³ Cd(n,	p) ¹¹³ Ag		28	
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ 0 _{ref} [%]	24	
13.45	10.8	11.1	3.41	<u> </u>	■ 1967 Yu-We ★ 1963 V.N.Le
13.88	13.8	10.9	3.40	l mi	•
14.44	15.0	12.6	3.40		-
14.64	19.0	9.93	3.40	ອັ ₁₂	-
14.81	19.8	9.50	3.40	U S S S s	
14.88	22.3	10.3	3.40		
Ref. CS	is ⁹³ Nb(n,	$(2n)^{92m}$ Nb;	$\alpha_d = 1.2$	4	



(9/2 ⁻) 81.1 0.4 s
$\frac{1}{(5/2)^{+}}$ $\frac{100}{0}$ s
¹¹³ ₄₆ Pd μ
7/2+ < 6 36%
$\frac{7/2}{1/2}$ $\frac{43.5}{4}$ p^{-} 68.7 s
¹¹³ ₄₇ Ag _β -
× 0,000 00 86%
14.1 y $\frac{11/2^{-2}}{263.59} \beta^{-1/2^{-391.691}}$
$9.3 \times 10^{15} \text{ y} \frac{1/2 + \sqrt{-0}}{113 \text{ cd}}$
<u>92* 0</u> 113 ₁₀
49'''

Decay	data	used	for
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 113 Cd(n, p) 113 Ag.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	Form	product		[keV]	[%]
¹¹³ Cd	12.22 12	Cd-metal	¹¹³ Ag	5.37 h 5	298.6	10.0 <i>3</i>

$^{114}Cd(n,\alpha)^{111m}Pd$



¹¹⁴ Cd(n, c	ι) ^{111m} Pd.		$\begin{array}{c} & & & & & \\ \hline 11/2 & & & & \\ \hline 5/2^+ & & & & \\ \hline 111 & & & & \\ 46 Pd & & & \\ \end{array}$	$\begin{array}{c} 27\% \\ \beta^{-} 5.5 \text{ h} \\ 23.4 \text{ m} \\ \hline \\ 7/2^{+} 59.82 \\ 1/2^{-} 112^{-} 0 \\ 111 \\ 47 \\ A7 \\ \beta^{-} \end{array} \begin{array}{c} 0.7\% \\ \beta^{-} \\ 7 \\ 7 \\ \beta^{-} \end{array}$	$\frac{11/2^{-\sqrt{396.22}}}{\frac{11/2^{+}}{48}Cd}$ 48.54 m
Chemical Form	Reaction product		T _{1/2}	Ε _γ [keV]	Υ _γ [%]
Cd-metal	^{111m} Pd	5.	5 h 1	172.2	46 5

Decav	data	used for	11
Decay	uata	useu IOI	

Abundance

[%]

28.73 42

Target

nucleus

 114 Cd

$^{116}Cd(n, 2n)^{115g}Cd$

$^{116}Cd(n, 2n)^{115g}Cd$				
E _n [MeV]	ರ [mb]	$\pm \Delta \sigma_{total}$ [%]	±Δ σ ref [%]	
13.45	793	3.51	3.31	
13.56	851	3.46	3.30	
13.74	837	3.31	3.30	
13.88	814	3.56	3.30	
13.96	828	5.63	3.30	
14.19	829	3.43	3.29	
14.43	809	3.45	3.29	
14.63	807	3.85	3.30	
14.80	794	3.56	3.30	
14.88	818	4.15	3.30	
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.2$				





Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹¹⁶ Cd	7 10 18	Cd_metal	¹¹⁵ gCd	53.16 h 5	492.4	8.03 18
Cu	Cu 7.49 16 Cu-metal -Cu 55.40 ll 5	55.40 II 5	527.9	27.5 6		



$^{116}Cd(n, 2n)^{115m}Cd$



Target nucleus	Abundance [%]	Chemical Form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					484.5	0.29 10
¹¹⁶ Cd	7.49 18	Cd-metal	115m Cd	44.56 d <i>24</i>	933.8	2.0 7
					1290.6	0.9 3

Decay data used for	116 Cd(n, 2n) 115m Cd
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¹¹⁶Cd(n, 2n)¹¹⁵Cd

This is the sum of the ${}^{116}Cd(n, 2n){}^{115m}Cd$ and ${}^{116}Cd(n, 2n){}^{115g}Cd$ cross sections measured independently and presented on previous pages.

$\begin{array}{c ccccccccccccccccccccccccccccccccccc$
<u>1/2+</u> 0 115 50Sn



113 In(n, n') 113m In



Decay data used for	113 In(n, 1	n') ^{113m} In.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹¹³ In	4.29 5	In-metal	^{113m} In	99.476 m 23	391.7	64.94 <i>17</i>



$^{115}In(n,\alpha)^{112}Ag$



Decay data used for 115 In(n, α) 112 Ag.						¹¹² 48Cd
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹¹⁵ In	95.71 5	In-metal	¹¹² Ag	3.130 h 8	617.4	43 5

¹¹⁵In(n, p)^{115m}Cd



$\begin{array}{cccccccccccccccccccccccccccccccccccc$
^{1/2+} 0 115 50

Decay data used for

 115 In(n, p) 115m Cd.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹¹⁵ In	95.71 5	In-metal	^{115m} Cd	44.56 d 24	933.8	2.0 7

¹¹⁵In(n,p)^{115g}Cd 115 In(n, p) 115g Cd 6 E_n $\pm \Delta \sigma_{total}$ $\pm \Delta \sigma_{ref}$ σ 5 [MeV] [mb] [%] [%] 13.47 3.90 9.67 2.27 4 Cross Section [mb] 13.64 3.83 5.12 2.26 4.33 2.26 13.88 4.26 3 2.25 14.05 4.64 8.79 TENDL-2014 EAF-2010 ROSFOND-2010 Present Data 2006 Fei Tuo+ 1993 A.Grallert+ 1993 C.Konno+ 1990 J.Csikai+ 1988 Ke Wei+ 1987 B.E.Leshche 1983 T.B.Ryves+ 4.79 2.25 2 14.28 6.59 14.47 4.93 3.87 2.25 1 14.68 4.97 3.84 2.26 14.86 5.18 3.48 2.26 0 13.0 13.2 13.4 13.6 13.8 14.0 14.2 14.4 14.6 14.8 15.0 Ref. CS is ${}^{93}Nb(n,2n){}^{92m}Nb;$ $\alpha_d = 1.2$ Neutron Energy [MeV]



$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$^{5.0\%}_{\beta^-}$ 4.486 h 4.41×10 ¹⁴ y $^{\beta^-}$
	^{1/2+} 0 115 50

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹¹⁵ In	05 71 5	In motal	¹¹⁵ gCd	53.16 h 5	492.4	8.03 18
In	95.71 5 In-metal	m-metai	"Ca	53.46 h 5	527.9	27.5 6

Decay data used for	115 In(n, p) 115g Cd.
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¹¹⁵In(n, p)¹¹⁵Cd

This is the sum of the ${}^{115}In(n, p){}^{115m}Cd$ and ${}^{115}In(n, p){}^{115g}Cd$ cross-sections measured independently and presented on previous pages.





115 In(n, n') 115m In



Decay data used for				11	⁵ In(n, n') ¹	^{15m} In.	
								Τ

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹¹⁵ In	95.71 5	In-metal	^{115m} In	4.486 h <i>4</i>	336.2	45.8 22



115 In(n, 2n) 114m In

0+	$\begin{array}{c} \underbrace{8^{-} \underbrace{501.98}_{0.50\%}}_{\text{43.1 ms}} 43.1 \text{ ms} \\ 4.4\% \\ EC \\ \underbrace{5^{+} 5^{$
¹¹⁴ 48Cd	0⁺ 114Sn

Decay	uata used 101	111(11, 211) 111.			50
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					190.3	15.56 15
¹¹⁵ In	95.71 5	In-metal	^{114m} In	49.51 d 1	558.4	4.4 4
					725.2	4.4 4

Decay data used for 115 In(n, 2n) 114m In



112 Sn(n, 2n) 111 Sn

One of the possible reasons for the discrepancy of the present cros-sections with the data of [21] may be a difference in the reference data. For example, the 761.2 keV intensity was 0.66 in [21] but it is 1.48 in the present work, and this value agrees with intensities of other gamma lines observed in the present experiment.

$\frac{\frac{1/2^{-537.15}}{9/2^{+\frac{1}{2}}}}{2.8049 \text{ d}}$	35.3 m
$\frac{\frac{11/2}{396.22}}{\frac{112^{+}}{48.54}}$ 48.54 m	

Decay	data used	112 Sn(n, 2r	$1)^{111}$ Sn.	48Cd													
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]											
					762.0	1.48 23											
112 C m	0.07.1	Sn motol	111 C m	¹¹¹ Sm 35	25.2 m (1153.0	2.7 4										
511	0.97 1	Sn-metal	SII-IIIetai	SII-IIIetai	Sn-metai	SII-Inetai	Sn	Sn	511	511	511	511	511	Sn	55.5 III 0	1610.5	1.31 20
					1914.7	2.0 3											

lata used	¹¹² Sn(n,	$(2n)^{111}$
iala uscu	D11(11)	<u> </u>



112 Sn[(n, np)+(n, d)]¹¹¹In

The amount of ¹¹¹In nuclei was measured next day after irradiation. The cross-section of the 112 Sn[(n, np)+(n, d)]¹¹¹In reaction was determined as the difference between the total ¹¹¹In production cross-section and the 112 Sn(n, 2n)¹¹¹Sn cross-section.

The comparatively large relative uncertainty of the data obtained is the consequence of this subtraction.

^{7/2+} €C ¹¹¹ 50Sn	35.3 m
1/2- ^{537.15} 7.7 m 9/2+♥ 0 EC ¹¹¹ In 2.8049 d	
$\frac{\frac{11/2^{-5}396.22}{396.22}}{\frac{1/2^{+}}{48}}$ 48.54 m	

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
112 C n	0.07 1	Sn motol	111 I n	2 9047 J 4	171.3	90.7 9
511	0.97 1	SII-IIIetai	111	2.8047 u 4	245.3	94.1 10

Decay data used for 112 Sn(n, x) 111 In.

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI

¹¹⁴Sn(n,2n)¹¹³Sn 114 Sn(n, 2n) 113 Sn 1400 1300 $\pm \Delta \sigma_{ref}$ E_n $\pm \Delta \sigma_{total}$ 1200 σ [MeV] [mb] [%] [%] 1100 1000 13.47 1004 5.51 1.73 Cross Section [mb] 900 13.88 1073 3.65 1.61 800 700 1115 1.61 14.05 3.47 600 1.59 14.28 1183 5.64 500 - TENDL-2014 - ENDF/B-VII.1 - EAF-2010 - JENDL-4.0 - CENDL-3.1 Present Data 2005 E.Betak+ 1988 Y.Ikeda+ 1970 J.K.Temp 1.58 14.49 1200 2.93 400 300 14.70 1262 5.90 1.59 200 14.88 1268 3.05 1.60 100 Ref. CS is 27 Al(n, α) 24 Na; $\alpha_d = 1.2$ 0 13.0 13.2 13.4 13.6 13.8 14.0 14.2 14.4 14.6 14.8 15.0 Neutron Energy [MeV]

114 Sn(n, 2n) 113 Sn

$EC \xrightarrow{7/2+}{77.389} 21.4 \text{ m}$ $EC \xrightarrow{1/2+}{1/2+} 0$ 115.09 c $\frac{1}{1/2-} 391.691}{1.6582 \text{ h}}$
<u>9/2+ 0</u> 113In 49In

Decer	data	usad	for
Decay	data	usea	IOr

 114 Sn(n, 2n) 113 Sn.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹¹⁴ Sn	0.66 1	Sn-metal	¹¹³ Sn	115.09 d 3	391.7	64.97 17

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



$^{116}\!\mathrm{Sn}(\mathrm{n},\mathrm{p})^{116\mathrm{m}}\mathrm{In}$



D	ecay data used	for ¹¹⁶ Sn(1	n, p) ^{116m} In.			¹¹⁶ 50Sn
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					416.9	27.2 4
116 Sn	14.54 9	Sn-metal	116m In	54.29 m 17	1097.3	58.5 8
					1293.6	84.8 12



 117 Sn(n, p) 117m In

The contribution of $^{120}Sn(n,\alpha)^{117}Cd$ reaction to population of the ^{117m}In level via ^{117}Cd decay was estimated as $<\!1\%$ that is several times less than the data uncertainty. This contribution was not accounted for.

<u>1/2+</u> 117 50
$ \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c}$
$\frac{(11/2)^{-} 136.4}{1/2^{+} 0} \beta^{-} 3.36 \text{ h}$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹¹⁷ Sn	7.68 7	Sn-metal	^{117m} In	116.2 m 3	315.3	19.1 8

 117 Sn(n, p) 117m In.

Decay data used for



117 Sn(n, p) 117g In

The contribution of 120 Sn(n, α) 117 Cd reaction to population of the 117 In ground state via 117 Cd decay was estimated as <1% that is several times less than the data uncertainty. This contribution was not accounted for.

$\begin{array}{c} (11/2)^{-} & 136.4 \\ 1/2^{+} & 0 \\ 117Cd \\ 48Cd \\ \beta^{-} \end{array}$	3.36 h 2.49 h $1/2^{-315.302}$ β^{-1} 43 $9/2^{+}$ 0 117 β^{-}	16.2 m 3.2 m <u>11/2[∑] 314.58</u> 13.60 d ▼ 1/2 ⁺ 0 117 Sn
	1	

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹¹⁷ Sn	7.68 7	Sn-metal	^{117g} In	43.2 m <i>3</i>	552.9	100 10

 117 Sn(n, p) 117g In.

Decay data used for



117 Sn(n, p) 117 In

This is the sum of the ${}^{117}Sn(n, p){}^{117m}In$ and ${}^{117}Sn(n, p){}^{117g}In$ cross-sections measured independently and presented on the previous pages. The cross-section was not corrected for the contribution of ${}^{120}Sn(n, \alpha){}^{117}Cd$ reaction which was estimated to be less than 1% that is several times less than the data uncertainty.





118 Sn(n, 2n) 117m Sn

According to the decay schema fragment shown on this page lower, the metastable state 117m Sn (E^{*}=314.58 keV, T_{1/2}=14.00 d, I^{π}=11/2⁻) is populated not only in the 118 Sn(n, 2n)^{117m}Sn reaction but also via the processes 117 Sn(n, p)¹¹⁷In \rightarrow ^{117m}Sn and 120 Sn(n, α)¹¹⁷Cd \rightarrow ¹¹⁷In \rightarrow ^{117m}Sn. This contributon is small because most 117 In nuclei decay to the 117 Sn levels which have the lower spin values. The population of the metastable level 117m Sn with I^{π}=11/2⁻ during the 117 In decay is less than 0.5%.

The total contribution of the reactions mentioned above was estimated as less as 0.01% relatively to the 118 Sn(n, 2n) 117m Sn cross-section and was neglected.



¹¹⁸Sn(n, α)^{115g}Cd



$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
^{1/2+} 0 115 50	

Decay data	a used for
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¹¹⁸ Sn(n,	α) ^{115g} Cd.
----------------------	--------------------------------

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹¹⁸ Sn	24.22 9	Sn-metal	^{115g} Cd	53.46 h 5	492.4	8.03 18
					527.9	27.5 6



120 Sn(n, α) 117m Cd					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]		
13.47	0.100	20.0	3.40		
13.88	0.121	25.0	3.35		
14.05	0.120	25.0	3.34		
14.28	0.159	37.5	3.34		
Ref. CS is ²⁷ Al(n, α) ²⁴ Na; $\alpha_d = 1.8$					



(11/2) ⁻ 136.4 β ⁻ 3.36 h	
1/2 ⁺ 0 2.49 h	
1/2− 315.302 ß_	% = 1 1 0 0
9/2+▼ 0 ▲	116.2 m
117In	43.2 m
49''' β–	
	11/2 314.58 12 60 d
	13.60 u
	1/2+ 0
	^{11/} 50 Sn
	50

Decay data used for 120 Sn(n, α) 117m Cd.				¹¹⁷ ₅₀ Sn		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					564.4	14.7 8
¹²⁰ Sn	32.58 9	Sn-metal	^{117m} Cd	3.36 h 5	1066.0	23.1 7
					1997.3	26.2 5

$^{120}Sn(n,\alpha)^{117g}Cd$

120 Sn(n, α) 117g Cd						
E _n [MeV]	σ [mb]	$\pm\Delta \sigma_{ ext{total}}$ [%]	±∆ σ _{ref} [%]			
13.47	0.120	16.6	4.55			
13.88	0.174	13.8	4.51			
14.05	0.246	22.3	4.51			
Ref. CS is ²⁷ Al(n, α) ²⁴ Na; $\alpha_d = 1.9$						

Decay data used for



$\begin{array}{c} (\underline{11/2})^{-} & \underline{136.4} & \beta - \\ \hline \underline{1/2^{+}} & 0 \\ 117 \\ 48 \\ \end{array} \\ \begin{array}{c} & & \\ & \\ & \\ & \\ & \\ & \\ & \\ & \\ & \\ $	$\begin{array}{c} 3.36 \text{ h} \\ 2.49 \text{ h} \\ & \checkmark^{\beta} \\ \underline{1/2}^{-} 315.302 \\ \underline{9/2}^{+} \psi \\ \underline{9/2}^{+} \psi \\ \underline{117} \\ \underline{43} \\ \underline{43} \\ \underline{6} \\ $	16.2 m 3.2 m 11/2 [→] 314.58 ▼ 1/2 ⁺ 0
		¹¹⁷ 50 Sn
T _{1/2}	Ε _γ Γκονι	Υ _γ [%]

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²⁰ Sn	32.58 9	Sn-metal	^{117g} Cd	2.49 h <i>4</i>	1303.3 1576.6	18.4 6 11.2 4

 120 Sn(n, α) 117g Cd.

$^{120}Sn(n,\alpha)^{117}Cd$



This is the sum of the 120 Sn(n, α) 117m Cd and 120 Sn(n, α) 117g Cd cross sections measured independently and presented on previous pages.





$^{120}Sn(n, 2n)^{119m}Sn$



Decay data used for

 120 Sn(n, 2n) 119m Sn.

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹²⁰ Sn	32.58 9	Sn-metal	^{119m} Sn		23.875	16.50 22
				293.1 d 7	25.044	7.9289 7
					25.271	14.8392 11

124 Sn(n, 2n) 123g Sn



$\begin{array}{c} \frac{3/2^+}{11/2^-} & 24.6 \\ \hline 11/2^- & 0 \\ \hline 123Sn \\ 50Sn \\ \beta- \end{array}$	[—] 40.06 m 129.2 d
E	^{7/2+} 123Sb X.

Decay data used for	124 Sn(n, 2n) 123g Sn.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²⁴ Sn	5.79 5	Sn-metal	^{123g} Sn	129.2 d 4	1088.6	0.60 10



124 Sn(n, 2n) 123m Sn



Decay data used for	124 Sn(n, 2n) 123m Sn.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²⁴ Sn	5.79 5	Sn-metal	^{123m} Sn	40.06 m <i>1</i>	160.3	85.7 4



124 Sn(n, 2n) 123 Sn

This is the sum of the 124 Sn(n, 2n) 123m Sn and 124 Sn(n, 2n) 123g Sn cross-sections measured independently and presented on the previous pages.

3/2+ 24.6 β– 40.06 m 11/2-0 129.2 d 50 7/2 ំSb



¹²⁰Te(n, 2n)^{119m}Te

		$EC \underbrace{\frac{11/2^{-} 260.96}{1/2^{+} 0}}_{1/2^{+} 0} 4.$.70 d 16.03 h
		EC ¹¹⁹ 52 Te	
	5/2+ 38.19 FC 51Sb	<u>n</u>	
<u>11/2</u> <u>1/2</u> +	<u>89.531</u> 293.1 d		
1	¹⁹ Sn		

Decay data used for		120 Te(n, 2n) 119m Te.		¹¹⁹ 50Sn		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					153.6	66 <i>3</i>
¹²⁰ Te	0.096 2	TeO_2	^{119m} Te	4.70 d 4	270.5	28.0 4
					1212.7	66.1 3



¹²⁰Te(n, 2n)¹¹⁹gTe

EC <u>11/2⁻ 260.96</u> 12 ⁺ 4.70 d 16.03 h EC 119 EC 52
5/2+ 38.19 h ► 119Sb
$\frac{\frac{11/2^{-} 89.531}{293.1}}{\frac{1/2^{+}}{50}}$ 293.1 d

Decay	data	used	for
Decay	uuuu	uscu	101

 120 Te(n, 2n) 119g Te.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²⁰ Te	0.096 2	TeO ₂	^{119g} Te	16.05 h 5	644.0 699.9	84.1 5 10.1 5


¹²⁰Te(n, 2n)¹¹⁹Te

This is the sum of the 120 Te(n, 2n) 119m Te and 120 Te(n, 2n) 119g Te cross sections measured independently and presented on the previous pages.





¹²²Te(n, p)¹²²Sb



Decay	data	used	for
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 122 Te(n, p) 122 Sb.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²² Te	2.603 4	TeO ₂	¹²² Sb	2.7238 d 2	564.2	70.68 18

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122 Te(n, 2n) 121m Te



Decay data used for	
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for 122 Te(n, 2n) 121m Te.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²² Te	2.603 4	TeO_2	^{121m} Te	164.2 d 8	212.2	81.5 10

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122 Te(n, 2n) 121g Te





Dee	Decay data used for 122 Te(n, 2n) 121g Te.				
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	E _γ [keV]	Υ _γ [%]
¹²² T -	2 602 4	TaO	¹²¹ g T a	10 17 d /	507.6	17.7 5
Ie	2.005 4	100_2	-Te	19.17 u 4	573.1	80.4 22



¹²²Te(n, 2n)¹²¹Te

This is the sum of the 122 Te(n, 2n) 121m Te and 122 Te(n, 2n) 121g Te cross sections measured independently and presented on the previous pages.



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¹²⁴Te(n, 2n)^{123m}Te

7/2+	$\frac{\frac{11/2}{247.55}}{119.7} d$
^{7/2+} 123 51Sb	EC 52 1 C

Decay data us	sed for
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 124 Te(n, 2n) 123m Te.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²⁴ Te	4.74 14	TeO ₂	^{123m} Te	119.2 d 1	159.0	84.0 <i>4</i>

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¹²⁶Te(n,2n)^{125m}Te 126 Te(n, 2n) 125m Te 1200 $\pm \Delta \sigma_{ref}$ E_n σ $\pm \Delta \sigma_{total}$ [MeV] [mb] [%] [%] 1000 945 13.67 17.0 1.65 Cross Section [mb] 800 14.06 1214 18.2 1.63 14.28 1009 1.63 15.8 600 1107 1.63 14.46 12.6 400 940 1.64 14.67 18.0 TENDL-2014 EAF-2010 Present Data Ref. CS is ${}^{93}Nb(n,2n)^{92m}Nb;$ $\alpha_d = 1.2$ 200 . 0 L 13.2 13.4 13.6 13.8 14.0 14.2 14.4 14.6 14.8 15.0 Neutron Energy [MeV]

¹²⁶Te(n, 2n)^{125m}Te



Decay data used for $^{126}\text{Te}(n, 2n)^{125m}\text{Te}.$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹²⁶ Te	18.84 25	TeO ₂	^{125m} Te	57.40 d 15	109.3	0.280 3

¹²⁸Te(n, p)^{128g}Sb





Decay data used for	128 Te(n, p) ¹	^{128g} Sb.
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Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹²⁸ Te		TeO ₂	^{128g} Sb	9.05 h 4	314.1	61 4
	21 74 0				526.5	45 <i>3</i>
	31.74 8				743.7	100 7
					754.0	100 7

¹²⁸Te(n,2n)^{127g}Te ¹²⁸Te(n, 2n)^{127g}Te 900 800 $\pm \Delta \sigma_{ref}$ E_n $\pm \Delta \sigma_{total}$ σ [MeV] [mb] [%] [%] 700 13.48 534 7.39 3.30 600 Cross Section [mb] 13.67 548 6.30 3.29 500 7.31 543 3.28 14.06 400 3.28 14.28 554 7.13 TENDL-2014 EAF-2010 ROSFOND-2010 Present Data 2007 F.Zhou+ 2000 C.Schnabel+ 1972 S.S.Hasan+ 1970 W.Bormann+ 1970 W.D.Lu+ 1964 C.G.Bonazzo 300 3.28 14.46 567 6.94 • 200 14.67 569 6.03 3.28 14.83 586 6.76 3.29 100 Ref. CS is ${}^{93}Nb(n,2n)^{92m}Nb;$ 0 $\alpha_d = 1.7$ 14.0 14.2 13.2 13.4 13.6 13.8 14.4 14.6 14.8 13.0 15.0 Neutron Energy [MeV]

¹²⁸Te(n, 2n)^{127g}Te

$\begin{array}{c} \underbrace{\frac{11/2}{3/2^{+}} \underbrace{\overset{5}{}^{(5)}}_{88.26}}_{3/2^{+}} \underbrace{\overset{2.4\%}{\beta^{-}}}_{9.35 h} \\ \underbrace{\overset{3/2^{+}}{127}}_{52} Te \\ \end{array} $
^{5/2+} 127 53

Decay data used for	128 Te(n, 2n) 127g Te.
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Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹²⁸ Te	31.74 8	TeO ₂	^{127g} Te	9.35 h 7	360.3	0.135 14
					417.9	0.99 10



¹³⁰Te(n, 2n)^{129m}Te

3/2+ ¥ 0 129 Te β- 52 Te β-	
$\begin{array}{c c} \frac{7/2^{+} & 1.57 \times 10^{7} \text{ y}}{129} \\ \hline \beta \\ \frac{1/2^{+}}{53} \\ \beta \\ 54 \end{array}$	0

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁰ Te	34.08 62	TeO ₂	^{129m} Te	33.6 d 1	729.6	0.70 14

 130 Te(n, 2n) 129m Te.

Decay data used for



¹³⁰Te(n, 2n)^{129g}Te



Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁰ Te	34.08 62	TeO ₂	^{129g} Te	69.6 m 3	459.6 487.4	7.7 6 1.42 10

Decay data used for	130 Te(n, 2n) 129g Te.

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¹³⁰Te(n, 2n)¹²⁹Te

This is the sum of the 130 Te(n, 2n) 129m Te and 130 Te(n, 2n) 129g Te cross sections measured independently and presented on the previous pages.





138 Ba(n, α) 135m Xe

	$15.29 \text{ m} \frac{11/2-526.551}{9.14 \text{ h}} \frac{\beta}{3/2^{+}} \underbrace{\downarrow}_{54} \underbrace{\downarrow}_{\beta-} $					
^{5m} Xe.	2.3	3×10 ⁶ y ^{7/2⁺} 135 55CS	⁰ β- <u>3/2+</u> 135 56 Ba			
Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]			
^{135m} Xe	15.29 m 5	526.6	80.4 3			

Decav	data	used	fo
Doou	anca	abea	101

Г

¹³⁸ Ba(n,	α) ^{135m} Xe.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁸ Ba	*99.8 1	BaCO ₃	^{135m} Xe	15.29 m 5	526.6	80.4 <i>3</i>



$^{138}\text{Ba}(n,\alpha)^{135}\text{Xe}$

$15.29 \text{ m} \frac{11/2-526.551}{9.14 \text{ h}} \frac{\beta}{135 \times e} \beta^{3}}{54 \times e} \beta^{3}$	
2.3×10 ⁶ y $\frac{7/2^{+}}{135}$ CS β^{-}	^{3/2+} 0 135 56Ba

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁸ Ba	*99.8 1	BaCO ₃	¹³⁵ Xe	9.14 h 2	249.8 608.2	90 3 2.90 <i>13</i>

 138 Ba(n, α) 135 Xe.

*- refer to samples with separated isotopes of Barium

Decay data used for

¹³⁸Ba[(n, np)+(n, d)]¹³⁷Cs

138 Ba(n, x) 137 Cs					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ σ _{ref} [%]		
13.73	0.36	19.4	0.70		
14.73	1.21	11.5	0.69		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.01$					



$\frac{7/2^{+} \ 30.07 \ y}{137}_{55}Cs_{\beta-}$	<u>↓</u> 11/2-661.660 ↓ 2.552 m	
	^{3/2+} 0 137Ba	

Decay data used for	138 Ba(n, x) 137 Cs.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁸ Ba	71.698 42	BaCO ₃	¹³⁷ Cs	30.08 y 9	661.7	85.10 20



138 Ba(n, p) 138 Cs

Only $(81\pm2)\%$ of the ^{138m}Cs metastable state decay to the ^{138g}Cs ground state. Strictly speaking, <u>the cross section [¹³⁸Ba(n, p)^{138g}Cs + 0.81*(¹³⁸Ba(n, p)^{138m}Cs)]</u> was measured in the <u>experiment</u>.

The difference with the 138 Ba(n, p) ${}^{138(m+g)}$ Cs cross-section evaluations presented in the Figure is 0.19^{*138} Ba(n, p) 138m Cs. This value can be estimated using EAF-2010 and TENDL-2014 evaluations that contain the data for the metastable state. The difference should be approximately 0.2 - 0.3 mb.

6- < ⁶ ^(*) β ⁻ 2.91 m 3- ▼ 0 138 CS β ⁻ 33.41 m
⁰⁺ 138 56Ba

Decay data used for $^{138}Ba(n, p)^{138}Cs$.						¹³⁸ Ba
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹³⁸ Ba *99.8 <i>1</i>			¹³⁸ Cs		462.8	30.7 6
	BaCO	¹³⁸ Ca		22 11 m 19	1009.8	29.8 6
	99.0 1	DaCO ₃	CS	33.41 III 16	1435.9	76.3 16
					2218.0	15.2 3



138 Ba(n, 2n) 137m Ba

 $^{138}\text{Ba(n, np)}^{137}\text{Cs}$ contribution is neglected because of the little cross section and the very long half life.

$$\frac{7/2^{+} \ 30.07 \ y}{137 \ 55} \beta^{-} \frac{11/2^{-} \ 661.660}{11/2^{-} \ 661.660} 2.552 \ m$$

$$\frac{3/2^{+} \ 0}{137 \ 56} Ba$$

Decay	data	used	for
Ducay	uata	uscu	IUI

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁸ Ba	*99.8 1	BaCO ₃	^{137m} Ba	2.552 m <i>1</i>	661.7	89.90 14

 138 Ba(n, 2n) 137m Ba.

^{*} – refer to samples with separated isotopes of Barium

¹³⁶Ce(n, 2n)¹³⁵Ce



		$\begin{array}{c} & & \\ & & \\ & & \\ \hline & & \\ \hline & & \\ & & \\ & 1/2(^+) \checkmark & 0 \\ \hline & & & \\ & &$
11/2-268.219	^{5/2+} 19.5 h EC ¹³⁵ 57La	
₹ <u>3/2+</u> 0 135 Ba	20.7 11	

Decay	uata useu 101	CC(11, 21	I) CC.	50		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹³⁶ Ce 0.185 2		Gab	¹³⁵ Ce	¹³⁵ Ce 17.7 h <i>3</i>	518.1	13.6 6
	0 195 0				572.3	10.4 4
	0.185 2	CeO_2			577.1	5.14 19
				783.6	10.6 4	

Decay data used for	¹³⁶ Ce(n	$(2n)^{135}Ce$	

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¹³⁸Ce(n, 2n)^{137m}Ce



Decay	data	used	for
Decay	uuuu	uscu	101

 138 Ce(n, 2n) 137m Ce.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹³⁸ Ce	0.251 2	CeO ₂	^{137m} Ce	34.4 h <i>3</i>	254.3	11.1 4



 $^{140}Ce(n, p)^{140}La$

<u>3-</u> 1. 140 57	<u>6781</u> c _ a β	1 •	
		0+ 140 58	Ce

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
	88.450 <i>51</i>	CeO ₂	¹⁴⁰ La	1.67855 d 12	328.8	20.3 3
¹⁴⁰ Ca					487.0	45.5 6
Ce					815.8	23.28 19
					1596.2	95.4 <i>14</i>

 140 Ce(n, p) 140 La.

Decay data used for



¹⁴⁰Ce(n, 2n)¹³⁹Ce



Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁴⁰ Ce	88.450 51	CeO ₂	¹³⁹ Ce	137.641 d 20	165.9	79.89 2

Decay data used for $^{140}Ce(n, 2n)^{139}Ce$.



¹⁴²Ce(n, 2n)¹⁴¹Ce



Decay data used for	142 Ce(n, 2n) 141 Ce.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁴² Ce	11.114 51	CeO ₂	¹⁴¹ Ce	32.511 d 13	145.4	48.4 <i>3</i>

$^{151}Eu(n,\alpha)^{148m}Pm$

151 Eu(n, α) 148m Pm						
E _n [MeV]	σ [mb]	±Δ σ _{total} [%]	±∆ 0 _{ref} [%]			
13.47	1.76	10.7	1.23			
13.56	1.82	15.9	1.22			
13.68	1.89	10.6	1.21			
13.74	2.14	12.2	1.21			
13.88	1.88	10.6	1.20			
13.96	1.90	12.1	1.20			
14.05	1.97	9.11	1.19			
14.19	2.20	10.9	1.19			
14.26	1.89	10.0	1.19			
14.42	1.98	11.6	1.19			
14.47	2.17	8.67	1.19			
14.61	2.41	14.9	1.20			
14.68	2.28	11.3	1.20			
14.78	2.02	9.88	1.21			
14.86	2.19	12.3	1.21			
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.1$						



					41.29 d ⁶⁻ ≮ 5.370 d 1- ¥ 148 148 60 148 61	95.0% <u>137.9</u> β- Pm _{β-} 0+ 7×10 ¹⁵ γ	
I	Decay data use	d for ¹⁵¹ Eu	$n(n, \alpha)^{148m}$ Pm	l.		¹⁴⁸ ₆₂ Sm α	
	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]	
	47.81 6				550.3	94.9 12	
	*07.5 1	Eu_2O_3	^{148m} Pm	41.29 d 11	630.0	89.0 9	
	97.5 I				725 7	228 1	

725.7

32.8 4

- refer to samples with separated isotopes of Europium

Target

nucleus

¹⁵¹Eu



$^{151}Eu(n,\alpha)^{148g}Pm$



Decay data used for $^{151}Eu(n, \alpha)^{148g}Pm$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁵¹ Eu	47.81 6 *97.5 <i>1</i>	Eu ₂ O ₃	^{148g} Pm	5.368 d 7	1465.1	22.2 5



151 Eu(n, α) 148 Pm

This is the sum of the ${}^{151}Eu(n, \alpha){}^{148m}Pm$ and ${}^{151}Eu(n, \alpha){}^{148g}Pm$ cross sections measured independently and presented on the previous pages.



¹⁵¹Eu(n, 3n)¹⁴⁹Eu

$^{151}\mathrm{Eu}(n, 3n)^{149}\mathrm{Eu}$						
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ σ _{ref} [%]			
14.65	0.20	100	1.83			
14.81	1.27	19.0	1.84			
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.1$						





Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]	
151 _E .,	47.91	En O	149 E	02144	277.1	3.56 6	
Eu	47.81 0	Eu_2O_3	Eu	9 5.1 û 4	327.5	4.03 12	

Decay data used for $^{151}Eu(n, 3n)^{149}Eu$



$^{151}Eu(n, 2n)^{150m}Eu$



De	ecay data used t	for ¹⁵¹ Eu(n	¹⁵⁰ ₆₂ Sm			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁵¹ E.,	47.81 6	En O	150m	100k 7	333.9	4.0 8
Eu	*97.5 <i>1</i> Eu ₂ O ₃	Eu_2O_3	Eu	12.8 n <i>1</i>	406.5	2.8 6

¹⁵¹Eu(n,2n)^{150g}Eu $^{151}Eu(n, 2n)^{150g}Eu$ 1400 1300 $\pm \Delta \sigma_{total}$ E_n $\pm \Delta \sigma_{ref}$ σ 1200 [MeV] [mb] [%] [%] 1100 1000 13.56 1129 5.47 2.56 900 Cross Section [mb] 13.74 1174 5.01 2.55 800 1144 2.55 700 13.96 5.41 600 1095 2.55 14.19 5.74 500 2.55 14.42 1125 4.92 TENDL-2014 EAF-2010 Present Data 1996 Y.Ikeda+ 1996 J.W.Mead 1993 C.Konno+ 1974 S.M.Qairr 1972 D.R.Neth 400 • • • 14.61 1139 4.67 2.55 300 200 • 14.78 1221 4.26 2.55 100 Ref. CS is ${}^{93}Nb(n,2n)^{92m}Nb;$ $\alpha_d = 1.01$ 0 14.2 13.2 13.4 13.6 13.8 14.0 14.4 14.6 14.8 15.0 13.0 Neutron Energy [MeV]

$^{151}Eu(n, 2n)^{150g}Eu$

	12.8 h EC <u>5(~)</u> 35.8 y <u>5(~)</u> EC 63	Eu <u>0+ 1.79×10</u> ⁶ y <u>150</u> Gd α α
	11%	42.1 B_
	11%	42.1 β_
	12.8 h EC 0(⁻) 35.8 y 5(⁻) EC 63	<u>42.1</u> β- <u>0</u> Eu
,	► 150 EC 63	Eu <u>0+ 1.79×10⁶ y</u>
		¹⁵⁰ 64Gd∜ α
	^{0⁺} 150 62Sm	
	Г	V

Decay data used for	151 Eu(n, 2n) 150g Eu	u.
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Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
	A7 81 6				334.0	95.2 20
¹⁵¹ Eu	*97.5 <i>1</i> Eu ₂ O	Eu_2O_3	^{150g} Eu	36.9 y <i>9</i>	439.4	79.6 16
					584.3	52.1 14



$^{151}Eu(n, 2n)^{150}Eu$

The 151 Eu(n, 2n) 150 Eu cross-section was not measured in the experiment directly. This is the sum of the 151 Eu(n, 2n) 150m Eu and 151 Eu(n, 2n) 150g Eu cross-sections presented in the previous pages.

$12.8 h EC \xrightarrow{0(-)}{5(-)} 0$ $35.8 y \xrightarrow{5(-)}{150} 0$ $150 C G G = 0$ $150 G G = 0$ α
0⁺ 150 62Sm



¹⁵³Eu(n, α)¹⁵⁰Pm

Three γ -lines given in the table at the page bottom were used for experiments with samples with separated isotopes of Europium. If samples of a natural abundance were irradiated then only the last γ -line 1324.5 keV was used for ¹⁵³Eu(n, α)¹⁵⁰Pm cross-section determination. This was done to avoid a possible interference of the ¹⁵¹Eu(n, 2n)^{150m}Eu reaction where the first two γ -lines were also excited in a branch decay to ¹⁵⁰Sm.

$\begin{array}{c} (1^{-}) & 2.68 \text{ h} \\ \hline 150}{} Pm \beta^{-} & 12.8 \text{ h} & EC & 0^{(-)} & 42.1 & \beta^{-} \\ 35.8 \text{ y} & & 5^{(-)} & 0 \\ \hline EC & 150 \\ \hline 63 Eu \\ \hline 150 \\ 64 Gd \\ \hline \alpha \end{array}$
0⁺ 150 62Sm

Decay data used for $^{155}Eu(n, \alpha)^{150}Pm$.						
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁵³ Eu	52.19 6 *99.20 8	Eu ₂ O ₃	¹⁵⁰ Pm	2.698 h 15	831.9 1165.8 1324.5	11.9 8 15.8 11 17.5 12



$^{153}Eu(n, 2n)^{152m1}Eu$

The metastable state ^{152m1}Eu (45.6 keV 0⁻) is not populated from the higher lying metastable level ^{152m2}Eu (147.8 keV 8⁻) which decays to the ^{152g}Eu ground state (0.0 keV 3⁻). On the other hand, the ^{152m1}Eu does not populate the ^{152g}Eu in decay. Therefore, two independent experiments were carried out for measurement of cross-sections of ¹⁵³Eu(n, 2n)^{152m1}Eu (this page) and ¹⁵³Eu(n, 2n)^{152m2+g}Eu (next page).



Decay data used for		153 Eu(n, 2n) 152m1 Eu.		¹⁵² Sm		α
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁵³ Eu		En O	^{152m1} Eu	9.3116 h <i>13</i>	121.8	7.0 8
	52.19 6				344.3	2.4 4
	*99.20 8	Eu_2O_3			841.6	14.2 16
					963.4	11.6 13

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$^{153}Eu(n, 2n)^{152m2+g}Eu$

The main attention of the experiment was concentrated on the accurate measurement of cross-sections which lead to generating long-lived reaction products. The ^{152m2}Eu metastable level (147.8 keV 8⁻) decays to the long-lived ^{152g}Eu with a 100% probability. The cross-section ¹⁵³Eu(n, 2n)^{152m2+g}Eu was determined in this work.

				2	$EC \xrightarrow{8-147.81}{147.81} 72$ $EC \xrightarrow{3-10}{45.5994} 0$ $EC \xrightarrow{63}{147.81} 72$ $BC \xrightarrow{147.81}{72} 72$	[∞] 96 m [−] 9.274 h 13.542 y ″
Decay d	lata used for	¹⁵³ Eu(n, 2n)	^{152m2+g} Eu.	0⁺ 152Sm		0+ 1.08×10 ¹⁴ y 152Gd 64Gd α
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					344.3	26.59 20
¹⁵³ Eu	52.19 6	Eu O	¹⁵² g _E ,	13 517 v 0	778.9	12.93 8
Eu	*99.20 8	Eu_2O_3	Eu	13.517 y 9	1112.1	13.67 8
					1408.0	20.87 9



$^{153}Eu(n, 2n)^{152}Eu$

The 153 Eu(n, 2n) 152 Eu cross-section was not measured in the experiment directly. This is the sum of the 153 Eu(n, 2n) 152m1 Eu and 153 Eu(n, 2n) ${}^{152m2+g}$ Eu cross-sections presented in the previous pages.





¹⁸¹Ta(n, α)^{178m}Lu

A separate determination of 181 Ta(n, α) 178m Lu and 181 Ta(n, α) 178g Lu cross-sections was based on the following circumstances.

There is no isomeric transition from the metastable state (9⁻) to the ground state (1⁺) of 178 Lu.

The remarkable spin difference cause also the different decay paths of the ^{178m}Lu and ^{178g}Lu inside the ¹⁷⁸Hf that are not crossed till the lowest states. The ¹⁷⁸Hf γ -transitions between the high spin states were chosen for characterization of the ^{178m}Lu decay, and those between the low spin states were used in processing the ^{178g}Lu data.

23.1 m $\frac{(9^{-})}{1(1^{+})}$ $\frac{120}{1}$ β^{-} 28.4 m 71 β^{-}	8-	× 1147.423 ▼	• 4.0 s
	<u>0+</u> 1	0 78 72	

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸¹ Ta			netal ^{178m} Lu 2		213.4	81.4 15
	99.988 2	Ta-metal		23.1 m 3	325.6	94.1 <i>16</i>
					426.4	97.0 18

 181 Ta(n, α) 178m Lu.

Decay data used



$^{181}Ta(n,\alpha)^{178g}Lu$

An essential difference between spins of 178m Lu and 178g Lu allows a separate determination of 181 Ta(n, α) 178m Lu and 181 Ta(n, α) 178g Lu cross-sections (see the previous page).

nucleus	[%]	form	product	T _{1/2}	Eγ [keV]	ι _γ [%]
¹⁸¹ Ta	00.000 2	Ta-metal ^{178g} Lu 28.4 m 2	178gr	28.4 m 2	1309.5	1.1 6
	99.988 2		1340.8	3.4 17		

Decay data used for

Та



$^{181}Ta(n,\alpha)^{178}Lu$

The ¹⁸¹Ta(n, α)¹⁷⁸Lu cross-section was not measured in the experiment directly. This is the sum of cross sections ¹⁸¹Ta(n, α)^{178m}Lu and ¹⁸¹Ta(n, α)^{178g}Lu presented in the previous pages.




$^{181}Ta[(n, np)+(n, d)]^{180m}Hf$



Decay data us	ed for 181 T	a(n, x) ^{180m} Hf	•
Abundance	Chemical	Reaction	

Т

Г

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	E _γ [keV]	Υ _γ [%]
					215.4	81.6 20
¹⁸¹ Ta	99.988 2	Ta-metal	^{180m} Hf	5.53 h 2	332.3	94.0 30
					443.2	81.7 25



¹⁸¹Ta(n, p)¹⁸¹Hf



Decay data used for 181 Ta(n, p) 181 Hf.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁸¹ Ta	99.988 2	Ta-metal	¹⁸¹ Hf	42.39 d 6	345.9 482.2	15.12 <i>12</i> 80.5 <i>4</i>



$^{181}Ta(n, 2n)^{180g}Ta$

	9- 1+ EC 86%	$\begin{array}{c} \hline 75.3 \\ 0 \\ 8.19 \\ \hline \textbf{Ta} \\ \beta^{-} \\ 14\% \end{array}$	×10 ¹⁵ y 52 h <u>0+</u> 180	<u>0</u>
<u>0+ 0</u> 180Hf			74	N

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸¹ Ta 99.98	99.988 2	99.988 2 Ta-metal	^{180g} Ta	8.154 h 6	93.3	4.51 16
	,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,				103.6	0.87 24

Decay data used for 181 Ta(n, 2n) 180g Ta



$^{180}W(n, 2n)^{179m}W$

	0.28 EC	× (1/2)-221.926 (7/2)-▼ 0 (7/2)-▼ 0 EC 74	6.40 m 37.05 m
9/2+ 0 179 72 Hf	7/2 ⁺ 1.82 y €C 73Ta		

Decay	data	used	for
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 180 W(n, 2n) 179m W.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
$^{180}\mathbf{W}$	0.12 1	W-metal	$^{179\mathrm{m}}\mathrm{W}$	6.40 m 7	221.5	8.8 7

$^{182}W(n,n\alpha)^{178m2}Hf$

182 W(n, n α) 178m2 Hf					
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ σ _{ref} [%]		
14.74	0.010	60.0	3.77		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.01$					



 $\frac{16^{+} 2446.05}{4.0 \text{ s}} 31 \text{ y}$ $\frac{8^{-} 1147.423}{1} 4.0 \text{ s}$ $\frac{0^{+} 0}{178} \text{ Hf}$

	Deeuy aata		··· (II, IIW)		•	-
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸² W 26 .		26.50 16 W-metal	^{178m2} Hf	31 y 1	325.6	94.1 16
	26.50 16				426.4	96.5 17
	20.50 16				495.0	70.1 16
					574.2	88.6 21

Decay data used for $^{182}W(n, n\alpha)^{178m^2}Hf$.

$^{182}W(n,\alpha)^{179m2}Hf$

$^{182}W(n, \alpha)^{179m2}Hf$					
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ 0 _{rej} [%]		
14.28	0.0022	59.1	3.71		
14.82	0.0080	62.5	3.72		
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.1$					





			(11) (11)			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
182 W	26.50 16	W-metal	^{179m2} Hf	25.05 d 25	315.9	20.3 7
					362.6	39.6 15
					409.7	21.5 8
					453.6	68 <i>3</i>

Decay data used for $^{182}W(n, \alpha)^{179m^2}Hf$



$^{182}W(n,p)^{182m2}Ta$



p) ^{182m2} Ta

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
182 W	26.50 16	W-metal	^{182m2} Ta	15.84 m <i>10</i>	146.8 171.6	36.5 24 48.0 20

¹⁸²W(n, p)¹⁸²Ta





Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
182 W	26.50 16	W-metal	¹⁸² Ta	114.74 d 12	1121.3 1221.4	35.24 <i>8</i> 27.23 <i>10</i>



$^{183}W(n,\alpha)^{180m}Hf$

5.5 h $\frac{8^{-1141.48}}{\sqrt{9^{-9^{-0}}}} = \frac{9^{-75.3}}{180 - 2} > 1.2 \times 10^{15}$	⁵ y
<u>0+ 0</u> 180µf	

14.0

14.2

14.4

14.6

14.8

15.0

Decay data used for $^{183}W(\mathbf{n}, \alpha)^{180m}$ Hf.			⁰Hf.	¹⁰⁰ 72Ht		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸³ W	14.31 4	W-metal	^{180m} Hf	5.53 h 2	332.3 443.2	94.0 30 81.7 25

¹⁸³W(n, p)¹⁸³Ta



(3/2⁻) 1.06 183µf	7 h			
72'''	β– <u>7/2</u> +	5.1 d		
	183 73	βTa β-	<u>∢</u> 11/2+309	. <u>493</u> 52s
			♥ 1/2 [_]	0
			¹⁸³ 74	V

Decav	data	used	for
Docu,	ance	abea	101

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁸³ W	14.31 4	W-metal	¹⁸³ Ta	5.1 d 1	354.0	11.2 7

 $^{183}W(n, p)^{183}Ta.$

$^{184}W(n,\alpha)^{181}Hf$

184 W(n, α) 181 Hf						
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±Δ σ _{ref} [%]			
13.64	0.495	8.64	0.79			
13.88	0.484	8.65	0.77			
14.28	0.701	4.56	0.75			
14.47	0.639	15.6	0.76			
14.68	0.819	5.48	0.77			
14.82	0.850	5.18	0.79			
Ref. CS	is ⁹³ Nb(n,21	$n)^{92m}Nb;$	$\alpha_d = 1.1$			



$\begin{array}{c} \underline{1/2^{-} 42.39 \text{ d}} \\ \hline 181 \\ 72 \\ \text{Hf} \\ \beta - \end{array}$	
	^{7/2+} 181 73

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
184 W			¹⁸¹ Hf		133.0	43.3 5
	30.64 2	2 W-metal		$^{181}\mathrm{Hf}$	42.39 d 6	345.9
				482.2	80.5 4	

Decay data used for $^{184}W(n, \alpha)^{181}Hf$.



¹⁸⁴W(n, p)¹⁸⁴Ta



Decay data used for	¹⁸⁴ W(n, p) ¹⁸⁴ Ta
---------------------	--

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					252.9	44 3
^{184}W	30.64 2	W-metal	¹⁸⁴ Ta	8.7 h <i>1</i>	414.0	72 3
					920.9	32.0 15

$^{186}W(n,n\alpha)^{182m}Hf$

186 W(n, n α) 182m Hf				
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ σ _{ref} [%]	
14.72	0.0063	20.6	8.74	
14.81	0.0065	23.1	8.74	
Ref. CS is ²⁷ Al(n, α) ²⁴ Na; $\alpha_d = 1.5$				



$\begin{array}{c} 61.5 \text{ m} \frac{8^{-} \left(172.88}{1772.88}\right)^{58\%} \beta^{-} \\ 9 \times 10^{6} \text{ y} \frac{0^{+} \text{ v} 0}{182} \text{ Hf} \beta^{-} \end{array} \frac{3^{-} 182}{73} \text{ Ta} \end{array}$	⁰ 114.43 d β-
	0⁺ 182 74₩

Decay data used for		$^{186}W(n,n\alpha)^{182m}Hf.$				¹⁸² 74		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]		
					224.4	38 3		
186337	28 13 10	W-metal	182m t 16	(1 5 15	344.1	46 6		
vv	28.43 19		w-metal	HI	-metai mi	01.5 m 15	455.8	20.0 21
					942.8	23 3		

$^{186}W(n,\alpha)^{183}Hf$



(3/2⁻) 1.06 183 72Hf	<u>7</u> h β–			
	7/2+ 183 73	5.1 d Τα β-	<u>√</u> 11/2+309 ▼	. <u>493</u> 5.2 s
			1/2- 183 74	<u>o</u> V

Decav	data	used	for
Ducay	uata	uscu	101

 $^{186}W(n, \alpha)^{183}Hf.$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
186 W	28.43 19	W-metal	¹⁸³ Hf	1.067 h <i>17</i>	459.1 783.8	27 4 66 9

¹⁸⁶W(n, p)¹⁸⁶Ta





Decay data used for $^{186}W(n, p)^{186}Ta$.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
186 W	28.43 19	W-metal	¹⁸⁶ Ta	10.5 m 3	197.9 615.3	50 5 28 3

$^{186}W[(n, np)+(n,d)]^{185}Ta$

186 W(n, x) 185 Ta					
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ σ ref [%]		
14.72	0.155	9.68	5.82		
14.81	0.250	7.93	5.82		
14.88	0.270	7.33	5.82		
Ref. CS is ²⁷ Al(n, α) ²⁴ Na; $\alpha_d = 1.4$					



(7/2 ⁺) 49.4 185 73 Ta	m β–	
		$\frac{\frac{11/2^{+}}{197.43}}{\frac{3/2^{-}}{74}} 1.67 \text{ m}$ 75.1 d 75.1 d
		^{5/2+} 185 75Re

Decay data used for	¹⁸⁶ W(n, x)
---------------------	------------------------

¹⁸⁶ W(r	$(x)^{185}$ Ta
	ь л/ Ia.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
^{186}W	28.43 19	W-metal	¹⁸⁵ Ta	49.4 m 15	173.9 177.6	22.6 <i>17</i> 25.7 <i>10</i>



$^{186}W(n, 2n)^{185m}W$

$\frac{11/2^{+}}{3/2^{-}} \underbrace{\stackrel{197.43}{1}}_{75.1 \text{ d}} 1.67 \text{ m}}_{75.1 \text{ d}}$
^{5/2+} 185 75

Decay data used for $^{186}W(n, 2n)^{185m}W$.						¹⁶⁵ 75Re
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸⁶ W					131.6	4.33 21
	28 43 10	W motol	185m xx 7	1 67 m 2	131.6 4.33 164.3 0.59	0.59 3
	20.43 19	w-metai	1.67 m 3	173.7	3.26 15	
					187.9	0.81 5

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$^{186}W(n, 2n)^{185}W$

$^{186}W(n, 2n)^{185}W$					
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±∆ 0 _{ref} [%]		
13.64	1968	17.8	3.76	2	
13.88	1892	17.5	3.75	<u> </u>	
14.28	1903	12.0	3.75	ectio	
14.68	1969	13.1	3.75	C. v	
14.82	2073	11.6	3.76		
Ref. CS	is ⁹³ Nb(n,2)	n) ^{92m} Nb;	$\alpha_d = 1.1$		



	<u>11/</u> <u>3/2</u>					$\frac{11/2^{+}}{3/2^{-}} \underbrace{\stackrel{197.43}{7}}_{74} 1.67 \text{ m} \\ 75.1 \text{ d} \\ \beta -$		
]	Decay data use	d for ¹⁸⁶ W	$(n, 2n)^{185}$ W.			^{5/2+} 185 75Re		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]		
186 W	28.43 19	W-metal	185 W	75.1 d 3	125.4	0.0192 7		

185 Re(n, α) 182 Ta

185 Re(n, α) 182 Ta						
E _n [MeV]	σ [mb]	±Δ O _{total} [%]	±∆ 0 _{ref} [%]			
14.48	1.40	35.7	0.681			
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 1.01$						





Decay data used for	185 Re(n, α) 182 Ta.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁸⁵ Re	37.40 2	K ₂ ReCl ₆	¹⁸² Ta	114.74 d 12	1189.0 1221.4	16.49 5 27.23 10



¹⁸⁵Re(n, 2n)^{184m}Re

Most γ -rays used for ¹⁸⁵Re(n, 2n)^{184m}Re cross-section calculation belong to the ¹⁸⁴W. (An exception is a 104.7 keV γ -ray which belongs to the ¹⁸⁴Re). For γ -rays of the ¹⁸⁴W, a special care was taken to exclude γ -transitions which may accompany the ¹⁸⁴Re ground state decay.



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸⁵ Re			184mRe 169 d 8 104.7 184mRe 161.3 536.7 920.9		104.7	13.6 4
	27.40.2	V D ₂ Cl		6.56 23		
	37.40 2	$\mathbf{K}_2 \mathbf{K} \mathbf{e} \mathbf{C} \mathbf{I}_6$		169 a 8	536.7	3.33 11
					920.9	8.2 3

Decay data used for 185 Re(n, 2n) 184m Re.

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185 Re(n, 2n) 184g Re



Gamma-rays used for ¹⁸⁵Re(n, 2n)^{184g}Re cross-section calculation are presented in the table at the page bottom. They all belong to the ¹⁸⁴W. In contrast to the case considered on the previous page, the γ -rays chosen to characterize the ¹⁸⁵Re(n, 2n)^{184g}Re reaction may accompany not only the decay of ^{184g}Re but also this of ^{184m}Re, with a far less intensity though.

The 185 Re(n, 2n) 184g Re cross-section data was corrected for a possible admixture of the 184m Re decay.



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸⁵ Re			Cl ₆ ^{184g} Re	^{184g} Re 35.4 d 7	792.1	37.7 11
	37.40 2	K ₂ ReCl ₆			894.8	15.7 5
					903.3	38.1 12

Decay data used for 185 Re(n, 2n) 184g Re.



¹⁸⁵Re(n, 2n)¹⁸⁴Re

The ¹⁸⁵Re(n, 2n)¹⁸⁴Re cross-section was not measured in the experiment directly This is the sum of cross-sections ¹⁸⁵Re(n, 2n)^{184m}Re and ¹⁸⁵Re(n, 2n)^{184g}Re presented in the previous pages.



187 Re(n, α) 184 Ta



					(5 ⁻⁾ 8.7 h 184Ta β	
	Decay data used for 187 Re(n, α) 184 Ta.					⁰⁺ 184 74
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸⁷ Re	62.60 2	K ₂ ReCl ₆	¹⁸⁴ Ta	8.7 h 1	414.0	72 3



 187 Re(n, p) 187 W

				3/2- 23.72 h 187 74 β-		
Decay data used for 187 Re(n, p) 187 W.			$\frac{\frac{5/2^{+}}{187}}{75} Re^{\frac{1}{9}} \beta^{-} \frac{\frac{1}{187}}{76} Os$			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁸⁷ Re	62.60.2	KaReCla	¹⁸⁷ W	24.000 h 4	479.5	26.6 4
ite i	02.00 2	is zite CI6	**	27.000 II 7	685.8	33.2 5

¹⁸⁷Re(n, 2n)^{186g}Re



0⁺ 186₩ 74	$\begin{array}{c c} (8^+) & \stackrel{\checkmark}{\swarrow} 149 \\ \hline 1^- & \stackrel{\checkmark}{\blacktriangledown} 0 \\ \hline 186 \\ EC \\ 6.9\% \\ 93.1\% \end{array} 2.0 \times 10^5 \text{ y} \\ 90.64 \text{ h} \\ 90.64 \text{ h} \\ 93.1\% \end{array}$
	^{0⁺} 186Os

Decay data used for	Decav	data	used	for
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 187 Re(n, 2n) 186g Re.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁸⁷ Re	62.60 2	K ₂ ReCl ₆	^{186g} Re	3.7183 d 11	137.2	9.47 3

¹⁸⁶Os(n, p)^{186g}Re



0⁺ 186₩ 74₩	$\begin{array}{c c} (8^+) & \swarrow & 149 \\ \hline 1^- & \bigstar & 0 \\ 90.64 \text{ h} \\ \hline 90.64 \text{ h} \\ \hline 186 \text{ Re} \\ 6.9\% & 93.1\% \end{array}$
	0⁺ <mark>186</mark> Os 76

Decay dat	a used	for
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 $^{186}Os(n, p)^{186g}Re.$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁸⁶ Os	1.59 3	Os-metal	^{186g} Re	3.7183 d 11	137.2	9.47 3

¹⁸⁸Os(n, p)¹⁸⁸Re



The possible contribution of 189 Os(n, np) 188 Re + 189 Os(n, d) 188 Re is expected to be as small as <1%.



Decay data used for 188 Os(n, p) 188 Re.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁸⁸ Os	13.24 8	Os-metal	¹⁸⁸ Re	17.0040 h 22	155.0	15.61 <i>18</i>



¹⁸⁹Os(n, p)¹⁸⁹Re

Possible contribution of ${}^{190}Os(n, np){}^{189}Re + {}^{190}Os(n, d){}^{189}Re$ reactions is estimated as about 5% which is several times less than experimental uncertainties.

				5	²⁺ 24.3 h ¹⁸⁹ ₇₅ Re _β -	
Ι	Decay data used	d for ¹⁸⁹ Os	(n, p) ¹⁸⁹ Re.		<u>9/</u> <u>3/</u>	2-
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					147.1	1.24 21
¹⁸⁹ Oc	16 15 5	Oc motol	¹⁸⁹ D o	24.2 h 4	216.7	5.5 9
Us	10.15 5	Os-metal	ĸe	24.3 ft 4	219.4	4.5 7
					245.1	3.5 7

$^{190}Os(n,\alpha)^{187}W$



$\frac{3/2^{-}}{74} \frac{23.72 \text{ h}}{\beta^{-}}$	
	$\frac{5/2^{+}}{{}^{187}_{75}\text{Re}} \xrightarrow{4.35 \times 10^{10} \text{ y}}{\beta} - \frac{1/2^{-}}{{}^{187}_{76}\text{Os}}$

Decay data u	used for
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 190 Os(n, α) 187 W.

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁹⁰ Oc	26.26.2	Oc motal	187	24.000 h 4	479.5	26.6 4
Us	26.26 2 C	Os-metal	W	24.000 11 4	685.8	33.2 5

¹⁹⁰Os(n, p)^{190m}Re





Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]	
	26.26 2 Os-m	Os motol	Os-metal ^{190m} Re		361.1	12.1 11	
190				224.2	569.3	13.7 11	
Us		Os-metal		Ke 3.2 h 2	3.2 n 2	605.1	14.9 11
					673.1	9.4 6	

Decay data used for $^{190}Os(n, n)^{190m}Re$



$^{191}\mathrm{Ir}(n,\alpha)^{188}\mathrm{Re}$



Decay data used for	191 Ir(n, α) 188 Re.
Deedy data abea 101	II (II, W) IV.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁹¹ Ir	37.3 2	Ir-metal	¹⁸⁸ Re	17.0040 h 22	155.0	15.61 18



¹⁹¹Ir(n, p)¹⁹¹Os

3/2 - 74.38 9/2 - 10	13.10 h 15.4 d		
¹⁹¹ 76Os	β_		
		3/2+	0
		¹⁹¹ 77 lr	

•				L			
	Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
	¹⁹¹ Ir	37.3 2	Ir-metal	¹⁹¹ Os	15.4 d 1	129.4	26.50 4

Decay data used for	¹⁹¹ Ir(n.	$(p)^{191}Os$
Decay data used for	11(11,	(\mathbf{p}) US



191 Ir(n, 2n) 190m2 Ir



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
		3 2 Ir-metal	^{190m2} Ir		361.2	86.72 21
¹⁹¹ Ir	37.3 2			3.087 h 12	502.5	89.35 20
					616.5	90.14 22

Decay data used for	¹⁹¹ Ir(n	$(2n)^{190m^2}$ Ir
Decay data used for	11(11)	, 211) 11

191 Ir(n, 2n) $^{190g+m1}$ Ir





Eγ

[keV]

361.1

371.2

518.6

569.3

Yγ

[%]

13.0 6

22.8 7

34.0 15

28.5 13

-	11 (11, 211)	•	<u> </u>	
	Chemical form	Reaction product	T _{1/2}	

^{190g}Ir

11.78 d 10

Abundance

[%]

37.3 2

Target

nucleus

¹⁹¹Ir

 191 Ir(n, 2n) $^{190g+m1}$ Ir.

Ir-metal

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



191 Ir(n, 2n) 190 Ir

The ¹⁹¹Ir(n, 2n)¹⁹⁰Ir cross-section was not measured in the experiment directly. This is the sum of cross sections ¹⁹¹Ir(n, 2n)^{190m2}Ir and ¹⁹¹Ir(n, 2n)^{190m1+g}Ir presented in the previous pages.



A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI



¹⁹³Ir(n, α)^{190m}Re

Gamma transitions with energies coinciding with those of the 191 Ir(n,2n) 190m2 Ir reaction (T_{1/2}=3.087 h) were excluded.



Decay data used for 193 Ir(n, α) 190m Re.						¹⁹⁰ 76Os
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁹³ Ir	62.7 2	Ir-metal	^{190m} Re	3.2 h 2	119.1 673.1	11.1 <i>11</i> 9.4 6
¹⁹³Ir(n, p)¹⁹³Os



_{3/2−} 30.5 h	
¹⁹³ Os	
70 p-	4
	<u>11/2−№ 80.22</u> 10.53 d
	3/2+ ♥ 0
	193 ₁
	77"

Decay data used for	193 Ir(n, p) 193 Os.
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Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁹³ Ir	62.7 2	Ir-metal	¹⁹³ Os	30.11 h <i>1</i>	321.6 460.5	1.251 <i>18</i> 3.88 5

A. A. Filatenkov. - Neutron Activation Cross Sections Measured at KRI







Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁹³ Ir	62.7 2	Ir-metal	^{192g} Ir	73.829 d 11	316.5 468.1	82.86 <i>3</i> 47.84 <i>3</i>

Decay data used for	193 Ir(n, 2n) $^{192g+m1}$ Ir.
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192 Pt(n, 2n) 191 Pt



	3/2	- 2.9 d
Å	EC	¹⁹¹ 78 Pt
<u>11/2-</u> 171.28	4.94 s	
3/2+ 0		
¹⁹¹ 77 lr		

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
		Pt-metal	¹⁹¹ Pt	2.83 d 2	129.4	3.3 3
192 D 4	0.793 7				359.9	6.4 5
Pt 0.782 7	0.782 /				409.4	8.8 7
					538.9	15.9 12

Decay data used for 192 Pt(n, 2n) 191 Pt.

$^{194}Pt(n,\alpha)^{191}Os$



Poor statistics of the gamma peak located in the region of high background counts is the main reason for the large cross section uncertainty.

$\begin{array}{c c} 3/2^{-} & 74.38\\ \hline 9/2^{-} & 0\\ \hline 191\\ \hline 76 \\ \hline \end{array} \begin{array}{c} 13.10 \text{ h}\\ 15.4 \text{ d}\\ \hline \beta_{-} \end{array}$	<u>11/2- <, 171.28</u> 4.94 s
	3/2+ 0
	¹⁹¹ 77 lr

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁹⁴ Pt	32.97 10	Pt-metal	¹⁹¹ Os	15.4 d <i>1</i>	129.4	26.50 4

Decay data used for 194 Pt(n, α)¹⁹¹Os.

194 Pt(n, p) 194m2 Ir



The reaction product is the nuclide of a high spin state (10^+) . The main reason of the large uncertainty is the poor statistics of gamma peaks. The discharging gammas are in a cascade. Cross sections were corrected for gamma summing.





194 Pt(n, p) $^{194g+m1}$ Ir



Decay data used for 194 Pt(n, p) $^{194g+m1}$ Ir.						¹ 78 ⁷ Pt
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					293.5	2.5 3
¹⁹⁴ Pt	32.97 10	Pt-metal	^{194g} Ir	19.28 h 13	328.4	13.1 17
					645.1	1.18 16

 194 Pt(n, p) 194 Ir



The present data for 194 Pt(n, p) 194 Ir cross section were obtained by summing the experimental data measured separately for 194 Pt(n, p) ${}^{194g+m1}$ Ir and 194 Pt(n, p) 194m2 Ir. This was done for extending the comparative base because a limited number of evaluations exist in which the cross sections are split for ground and metastable states. Most evaluations contain the data for the total reaction cross-sections only.





 195 Pt(n, p) 195m Ir

The data were not corrected for possible contribution of ${}^{196}Pt(n, np){}^{195m}Ir + {}^{196}Pt(n, d){}^{195m}Ir$. According to EAF-2010 evaluation, this contribution can be neglected.

$\begin{array}{c} 11/2 - & 5^{(6)} & 95^{(6)} \\ 3/2^{+} & 0 & \beta - \\ \hline 195 \\ 77 \\ 77 \\ \end{array}$	$^{\frac{6}{-}}$ 3.8 h 2.5 h $\frac{13/2^{+} \stackrel{4}{\leftarrow}_{259.30}}{↓}$ 4.02 d
	<u>1/2− 0</u> 195Pt 78Pt

Decay data used for 195 Pt(n, p) 195m Ir.					78Pt		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]	
¹⁹⁵ Pt		Pt-metal	Pt-metal ^{195m} Ir	267 b 0	319.9	9.8 16	
	22.02.1				364.9	9.7 15	
	33.83 I Pt-1			3.0/ 11 8	432.9	9.8 15	
					684.9	9.8 16	

196 Pt(n, α) 193 Os



The large uncertainty of relevant gamma peak in the spectra propagates in the total cross section uncertainty.

3/2-30.5 h ¹⁹³₇₆Os β-0

Decay data used for 196 Pt(n, α)¹⁹³Os.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁹⁶ Pt	25.24 4	Pt-metal	¹⁹³ Os	30.11 h 1	460.5	3.88 5

196 Pt(n, p) 196m Ir							
E _n [MeV]	σ [mb]	$\pm \Delta \sigma_{total}$ [%]	±Δ σ ref [%]				
13.47	0.013	26.5	3.57				
13.88	0.017	22.5	3.57				
14.28	0.026	29.8	3.58				
14.47	0.031	15.4	3.58				
14.68	0.037	10.7	3.58				
14.82	0.045	10.7	3.58				
Ref. CS is 93 Nb(n,2n) 92m Nb; $\alpha_d = 2.0$							

196 Pt(n, p) 196m Ir



Two Figures are given to show the noticeable difference between our data and the data of three other experimental groups. The upper Figure presents all the data, the lower one gives the Figure fragment zoomed in.

When analyzing possible reasons for the significant difference in the data, the following circumstances were revealed:

1. In other experiments, the γ -peak with energy 356 keV was used for cross section determination. Our γ -countings have indicated that this peak consists of three close γ -lines, two of which are not related to the reaction studied. In our work, γ -peak with energy 355.9 keV was not used for cross section determination

3. Four other γ -lines were used for cross section determination in our work (see Table below). Cross sections obtained for any of these four γ -lines agree within errors.

4. The recommended spin value of the metastable state 196m Ir is high (10, 11⁻). The isomeric ratio for the 196 Pt(n, p) 196 Ir is expected to be low, about 0.01 [ref.24, p.23]. Our data agree well with the expected value.





Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
	25.24 4	Pt-metal	^{196m} Ir	140 h 2	393.5	97 4
19604					447.1	94 4
Pt				1.40 n 2	521.4	96 3
					647.3	91 4

Decay data used for 196 Pt(n, p) 196m Ir.

196 Pt(n, p) 196g Ir



Only one experimental result obtained fifty years ago was presented in EXFOR-file for the 196 Pt(n, p) 196g Ir cross section. The aim of our experiment was to supplement the measurement results for the 196 Pt(n, p) 196m Ir reaction discussed above. The value agrees within errors with the EXFOR data and with the available evaluations. Also, the isomeric ratio determined in our measurements agrees with the branching ratio systematics [ref.24, p.23].



196 Pt(n, p) 196 Ir



The present experimental data for 196 Pt(n, p) 196 Ir cross section is the sum of 196 Pt(n, p) 196m Ir and 196 Pt(n, p) 196m Ir cross-sections measured separately.

The agreement of the sum experimental data with a more number of available evaluations can be considered as a strengthened evidence for the correctness of our results obtained for the ¹⁹⁶Pt(n, p)^{196m}Ir cross section where the prominent discrepancy with the results reported by other authors was revealed (see above).



$^{197}\!Au(n,\alpha)^{194m2}Ir$



The studied reaction leads to population of a high spin state. The cross section is small and the main part of uncertainty is related to the poor peak statistics in the measured gamma spectra.

$$\frac{10,11) 190}{100} \stackrel{\beta}{}_{}^{-1,71 d} \stackrel{\beta}{$$



$^{197}\mathrm{Au}(n,\alpha)^{194\mathrm{g}+\mathrm{m}1}\mathrm{Ir}$

It should be noted that the first metastable state of ¹⁹⁴Ir with excitation energy 112 keV and half-life 31.85 ms was earlier omitted sometimes. This invoked the change of notation. So, the old ^{194m}Ir state can correspond to the present ^{194m2}Ir state and the ¹⁹⁷Au(n, α)^{194g}Ir cross-section in the old notation can be the ¹⁹⁷Au(n, α)^{194g+m1}Ir cross-section in the new notation. One must keep the attention to the possible variations of the used notations.

 $\begin{array}{c|c} \underline{(10,11) \quad 190} & \beta^{-} & 171 \ d \\ \hline (4,5,6) & 112 \\ \hline 1^{-} & 10 \\ \hline 194 \\ 77 \\ \hline 77 \\ \hline \beta^{-} \end{array} \qquad 31.85 \ \text{ms} \\ 19.15 \ \text{h} \\ \hline 194 \\ \hline 78 \\ \hline 194 \\ \hline 194 \\ \hline 78 \\ \hline 194 \\ \hline 194 \\ \hline 78 \\ \hline 194 \\ \hline 194 \\ \hline 78 \\ \hline 194 \\ \hline$

	2		< <i>/ /</i>			10
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁹⁷ Au	100 A		^{194g} Ir		293.5	2.5 3
		Au-metal		19.28 h 13	328.4	13.1 17
					645.1	1.18 16

Decay data used for 197 Au(n, α) $^{194g+m1}$ Ir.

$^{197}\mathrm{Au}(n,\alpha)^{194}\mathrm{Ir}$



The $^{197}Au(n,\alpha)^{194}Ir$ cross section presented on this page is the sum of the $^{197}Au(n,\alpha)^{194g+m1}Ir$ and $^{197}Au(n,\alpha)^{194m2}Ir$ cross sections measured separately (see above). This provides the comparison with a more number of evaluations. The present experimental data agree with most available evaluations within errors.





 $^{197}\mathrm{Au}(\mathbf{n},\mathbf{p})^{197\mathrm{m}}\mathrm{Pt}$

Two gamma lines are given in the Table of the decay data used. The first gamma line marked out by bold font had the determining influence on the total and reference uncertainty. The second one that has the energy 279 keV was used for increasing the result reliability.



Decay data used for

 $^{197}Au(n, p)^{197m}Pt.$

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
¹⁹⁷ Au	100	Au-metal	^{197m} Pt	95.41 m 18	346.5 279.0	11.1 <i>3</i> 2.4 <i>6</i>



 $^{197}Au(n, p)^{197}Pt$

The ¹⁹⁷Au(n, p)¹⁹⁷Pt cross section was measured after the total discharge of the metastable excitation state ^{197m}Pt. In this process, 96.7% of the ^{197m}Pt have populated the ¹⁹⁷Pt ground state via isomeric transitions but other 3.3% of the ^{197m}Pt have decayed to the ¹⁹⁷Au levels directly. The gammas listed in the Table below correspond to decay of the ¹⁹⁷Pt ground state exclusively. When we used them, then the cross sections of [¹⁹⁷Au(n, p)^{197g}Pt + 0.967¹⁹⁷Au(n, p)^{197m}Pt] were determined. To obtain the ¹⁹⁷Au(n, p)¹⁹⁷Pt cross section, the data were added by the 0.033 ¹⁹⁷Au(n, p)^{197m}Pt cross section measured earlier (see the previous page).



Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
197	100	Au motol	197g D +	10 9015 h 70	191.4	3.7 4
Au	100	Au-metai	Pl	19.8915 11 19	268.8	0.23 3

Decay data used for $^{197}Au(n, p)^{197}Pt$.



$^{197}Au(n, 2n)^{196m2}Au$



Decay	data used for	197 Au(n, 2n) ¹	¹⁹⁶ ₇₈ Pt			
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁹⁷ Au	100	Au-metal	^{196m2} Au	0661	147.8	43.5 15
					188.3	30.0 15
				9.0 11 1	285.5	4.4 5
					316.2	3.0 3



¹⁹⁷Au(n, 2n)¹⁹⁶Au

The cross-section data presented in the Table above are the weighted average of [1] and [3] corresponding data. The results were also corrected for the new reference data.



De	ecay data used	for ¹⁹⁷ Au(1	19	²⁶ ₇₈ Pt		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
¹⁹⁷ Au	100	Au-metal	^{196g} Au	6.1669 d <i>6</i>	333.0 355.7	22.9 9 87 3
					426.1	6.6 3



 205 Tl(n, α) 202 Au

The data accuracy is limited by a very weak gamma peak statistics which is the result of the small cross section, the low gamma intensity and the short half-life. An additional problem was impossibility to use irradiated samples repeatedly because of high background of the competing reaction 203 Tl(n, 2n) 202 Tl that produce gamma rays of the same energy. The half life of 202 Tl is 12.23 d.

(1 ⁻) 28.8 s 202 79Au β-	
	₀⁺ <mark>202</mark> Hg

Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]	
					439.5	9.2 5	
²⁰⁵ Tl	70.48 <i>1</i> T	Tl-metal	²⁰² Au	28.4 s 12	1125.3	2.12 7	
					1306.4	2.07 6	

Decay data used for 205 Tl(n, α) 202 Au



 205 Tl(n, p) 205 Hg

A large uncertainty in the reference γ -ray intensity and very weak peak statistics in the counted spectra are causes of high uncertainties in the measured cross section values.

					1/2⁻ 0 205 80Hg	5.2 m β–
	Decay data us	sed for ²⁰⁵ 7	fl(n, p) ²⁰⁵ Hg.			^{1/2+} 205 TI 81
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
²⁰⁵ Tl	70.48 1	Tl-metal	²⁰⁵ Hg	5.14 m 9	203.7	2.2 10



204 Pb(n, 2n) 203 Pb

Although the cross-section is convenient for measurement, the experimental results are scattered significantly. Perhaps, this is caused by the remarkable uncertainty of ²⁰⁴Pb abundance in the natural mixture. Anyway, the abundance uncertainty is a major contributor to the total uncertainty in our experiment.



Decay data used for	204 Pb(n, 2n) 203 Pb.
---------------------	---------------------------------

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²⁰⁴ Pb	1.4 1	Pb-metal	²⁰³ Pb	51.92 h 3	279.2 401.3	80.9 <i>19</i> 3.35 <i>10</i>



206 Pb(n, α) 203 Hg

There are rather comfortable conditions for this cross-section measurement. The problem may arise from the intersection with the reaction 204 Pb(n, 2n) 203 Pb which has a much higher cross-section, a coinciding gamma radiation but a shorter half-live. To provide the result reliability, the cooling time should be not less than 20 d. The main contribution to the data uncertainty is related to the weak gamma peak statistics.



Decay data used for	206 Pb(n, α) 203 Hg.
Boody data abou 101	1 × (11, W) 11-B.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²⁰⁶ Pb	24.1 1	Pb-metal	²⁰³ Hg	46.594 d 12	279.2	81.56 5



 209 Bi(n, α) 206m Tl

A very small cross section was measured. A possible contamination with 206 Pb that could generate a 206 Pb(n, p) 206m Tl reaction and disturb the data was found to be neglected for the bismuth sampes used. The weak gamma peak statistics is the main contributor to the experimental uncertainties. The gammas are cascading and corrections for gamma summing are necessary.

Only two evaluations were found for the 209 Bi(n, α) 206m Tl cross-section. EAF-2010 agrees satisfactory with the experimental data. TENDL-2014 evaluation is tens times higher (is not shown in the Figure). Note that the spin of the 206m Tl is recommended as high as (12⁻). If so then the isomeric ratio would be expected of order of 0.01. The corresponding value is approximately 0.012 in EAF-2010.

				<u>(</u>	$\begin{array}{c} (12^{-}) & 2643.11 \\ \hline & 2643.11 \\ \hline & & 3.74 \text{ m} \\ \hline & & & & \\ \hline & & & & \\ \hline & & & & \\ \hline & & & &$	n
De	ecay data used	for ²⁰⁹ Bi(n	$, \alpha)^{206m}$ Tl.			0+ 206 82 Pb
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					216.4	74 3
²⁰⁹ Bi	100	Pi O	206m 1	274 m 2	453.3	93 5
DI	100	$B1_2O_3$	11	5./4 M 3	457.2	22 3
					686.5	91 5

209 Bi(n, 2n) 208 Bi



The main problems of the present measurement are the extremely long half-life of the 208 Bi and overlapping its radiation with the very power background γ -line 2614.5 keV.

The experimental result is inside the comparatively narrow band formed by many evaluations.



Decay data used for 209 Bi(n, 2n) 208 Bi.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²⁰⁹ Bi	100	Bi ₂ O ₃	²⁰⁸ Bi	3.68•10 ⁵ y 4	2614.5	99.785 2

²⁰⁹Bi(n, 3n)²⁰⁷Bi



It is a rare situation for the neutron energy region covered in the present work when the (n, 3n) reaction can be excited.



Decay data used for 209 Bi(n, 3n) 207 Bi.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²⁰⁹ Bi	100	Bi ₂ O ₃	²⁰⁷ Bi	31.55 y 4	569.7 1063.7	97.75 <i>3</i> 74.5 <i>3</i>



232 Th(n, 2n) 231 Th

The main difficulty of the present experiment is high background and very high self absorption.



Decay data used for 232 Th(n, 2n) 231 Th.

Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²³² Th	100	Th-metal	²³¹ Th	25.52 h 1	84.2 102.3	6.6 <i>4</i> 0.436 <i>24</i>



 238 U(n, 2n) 237 U

The reaction is studied rather well. The evaluations and experimental data are in a good accordance.



Target	Abundance	Chemical	Reaction	T _{1/2}	Ε _γ	Υ _γ
nucleus	[%]	form	product		[keV]	[%]
²³⁸ U	99.2745 15	U-metal	²³⁷ U	6.752 d 2	208.0	21.2 3

Decay data used for $^{238}U(n, 2n)^{237}U$.

$^{241}Am(n,\alpha)^{238}Np$



This is one of the results of the KRI pioneering cross section measurements with the ²⁴¹Am [4]. In more detail, the experiment of high complexity is described in pp.26-30 of this paper.

For the 241 Am(n, α) 238 Np cross-section, we managed to determine only the upper limit. It is equal to 1.5 mb.

					2+ 238 93	2.117 d Np β-	
Dec	ay data used fo	or ²⁴¹ Am(n	$, \alpha)^{238}$ Np.				0+ 87.7 y 238 94 Pu α
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}		Ε _γ [keV]	Υ _γ [%]
						984.5	25.2 3
²⁴¹ Am	99.9 <i>1</i>	$Am(NO_3)_3$	²³⁸ Np	2.099 d	2	1025.9	8.75 7
						1028.5	18.23 12

$^{241}Am(n, 3n)^{239}Am$



Experiment of high complexity is described in detail on pp.26-30 of this work.

	(5/2)- EC 95 99.990%	Am α 0.010%
1/2 ⁺ 24110 y 239 94 Pu α		

Decay data used for $^{241}Am(n, 3n)^{239}Am$.				ά		
Target nucleus	Abundance [%]	Chemical form	Reaction product	T _{1/2}	Ε _γ [keV]	Υ _γ [%]
					226.4	3.3 4
²⁴¹ Am	99.99	$Am(NO_3)_3$	²³⁹ Am	11.9 h <i>1</i>	228.2	11.3 13
					277.6	15.0 17



 $^{241}Am(n, 2n)^{240}Am$

The pioneering neutron activation cross-section measurements with ²⁴¹Am [4] demonstrated the ability to obtain the cross section data for material of small mass ~3 mg but of very high own activity $\sim 10^8$ Bk. The experiment is described in detail in pp.26-30 of this paper.

Results of new experiments reported recently have generally confirmed the KRI data.



50.8 h 3

888.9

987.8

 $24.7 \ 5$

72.2 13

Decay data used	for ²⁴	¹ Am(n.	$(2n)^{240}$ Am.
Doou'y dulu ubou	101	T NTTT(110	#11 / 111

 $Am(NO_3)_3$

[%]

99.99

Target

nucleus

²⁴¹Am

²⁴⁰Am

ATTACHMENT. Reference cross sections

	KRI (old)		IRDFF v1.05			KRI (old)		IRDFF v1.05		
En			(nev	W)		En			(nev	W)
[MeV]	CS	Relative	CS	Relative		[MeV]	CS	Relative	CS	Relative
	[mb]	Error [%]	[mb]	Error			[mb]	Error	[mb]	Error
13.45	125.8	1.13	125.46	0.795		14.29	118.5	1 01	118 47	0.410
13.46	125.8	1.12	125.46	0.795		14.42	116.1	1.03	116.63	0.375
13.47	125.7	1.11	125.46	0.795		14.43	115.9	1.04	116.49	0.375
13.48	125.7	1.11	125.46	0.795		14.44	115.6	1.04	116.35	0.375
13.49	125.6	1.10	125.46	0.795		14.45	115.4	1.04	116.20	0.375
13.50	125.5	1.09	125.46	0.795		14.46	115.1	1.04	116.06	0.375
13.56	125.2	1.04	125.33	0.525		14.47	114.9	1.04	115.91	0.375
13.58	125.0	1.04	125.29	0.525		14.48	114.8	1.05	115.77	0.375
13.64	124.6	1.04	125.07	0.525		14.49	114.6	1.05	115.63	0.375
13.65	124.5	1.04	125.03	0.525		14.50	114.5	1.05	115.48	0.375
13.66	124.4	1.05	124.98	0.525		14.61	112.9	1.06	113.88	0.385
13.67	124.3	1.05	124.94	0.525		14.62	112.8	1.06	113.73	0.385
13.68	124.2	1.05	124.89	0.525		14.63	112.6	1.07	113.59	0.385
13.70	124.0	1.05	124.80	0.525		14.64	112.5	1.07	113.44	0.385
13.72	123.8	1.05	124.67	0.525		14.65	112.4	1.07	113.29	0.385
13.73	123.7	1.05	124.60	0.525		14.66	112.3	1.07	113.15	0.385
13.74	123.6	1.05	124.53	0.525		14.67	112.1	1.07	113.00	0.385
13.75	123.5	1.05	124.47	0.525		14.68	112.0	1.07	112.85	0.385
13.87	122.7	0.98	123.51	0.491		14.69	111.9	1.07	112.71	0.385
13.88	122.6	0.98	123.43	0.491		14.70	111.8	1.07	112.56	0.385
13.89	122.6	0.98	123.34	0.491		14.71	111.7	1.07	112.41	0.385
13.95	122.3	0.98	122.72	0.491		14.72	111.6	1.08	112.27	0.385
13.96	122.3	0.98	122.62	0.491		14.73	111.5	1.08	112.12	0.385
14.04	121.9	0.98	121.71	0.462		14.74	111.4	1.08	111.97	0.385
14.05	121.9	0.98	121.59	0.462		14.77	111.1	1.08	111.53	0.385
14.06	121.8	0.99	121.47	0.462		14.78	111.0	1.08	111.38	0.385
14.07	121.7	0.99	121.35	0.462		14.80	111.0	1.13	111.09	0.385
14.09	121.4	0.99	121.12	0.462		14.81	111.0	1.15	110.94	0.437
14.10	121.3	0.99	121.00	0.462		14.82	111.0	1.17	110.79	0.437
14.19	120.3	1.00	119.83	0.462		14.83	110.9	1.17	110.64	0.437
14.21	119.9	1.00	119.56	0.410		14.84	110.8	1.17	110.49	0.437
14.23	119.6	1.00	119.29	0.410		14.85	110.7	1.17	110.35	0.437
14.25	119.2	1.01	119.02	0.410		14.86	110.6	1.18	110.20	0.437
14.26	119.1	1.01	118.88	0.410		14.87	110.5	1.18	110.05	0.437
14.27	118.9	1.01	118.74	0.410		14.88	110.4	1.18	109.90	0.437
14.28	118.7	1.01	118.61	0.410						

Table I. The 27 Al(n, a) 24 Na reference cross section

Table Ia. Decay data used at the ${}^{27}Al(n, \alpha){}^{24}Na$ cross section determination.

Reaction product	T _{1/2}	E _□ [keV]	Y [%]		
24 N L	14.007 h /2 1368.6 99.9936				
Ina	14.997 n <i>12</i>	2754.0	99.855 <i>5</i>		

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Table II. The ⁹³ Nb(n, 2n) ⁹² Nb reference cross section	^{2m} Nb reference cross section	able II. The ⁹³ Nb(n, 2n
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E.,	KRI ((old)	IRDFF (nev	v1.05 w)	E.,	KRI ((old)	IRDFI (ne	F v1.05 w)
[MeV]	CS [mb]	Relative Error	CS [mb]	Relative Error	[MeV]	CS [mb]	Relative Error	CS [mb]	Relative Error
13.45	451.5	1.40	452.00	0.633	14.29	459.3	1.44	460.10	0.543
13.46	451.6	1.39	452.20	0.633	14.42	460.2	1.22	460.12	0.549
13.47	451.7	1.37	452.40	0.633	14.43	460.2	1.26	460.12	0.549
13.48	451.8	1.36	452.60	0.633	14.44	460.3	1.30	460.12	0.549
13.49	451.9	1.35	452.80	0.633	14.45	460.3	1.35	460.13	0.549
13.50	451.9	1.34	453.00	0.592	14.46	460.4	1.39	460.13	0.549
13.56	452.4	1.26	454.03	0.592	14.47	460.4	1.43	460.13	0.549
13.58	452.4	1.30	454.38	0.592	14.48	460.4	1.42	460.13	0.549
13.64	452.4	1.37	455.31	0.592	14.49	460.3	1.40	460.13	0.549
13.65	452.5	1.36	455.46	0.592	14.50	460.3	1.39	460.13	0.549
13.66	452.6	1.34	455.60	0.592	14.61	459.9	1.22	460.15	0.565
13.67	452.8	1.33	455.75	0.592	14.62	459.8	1.25	460.15	0.565
13.68	452.9	1.32	455.90	0.592	14.63	459.6	1.29	460.16	0.565
13.70	453.1	1.29	456.19	0.592	14.64	459.5	1.32	460.16	0.565
13.72	453.4	1.26	456.43	0.592	14.65	459.3	1.36	460.16	0.565
13.73	453.5	1.25	456.55	0.592	14.66	459.2	1.39	460.16	0.565
13.74	453.6	1.23	456.68	0.592	14.67	459.0	1.43	460.16	0.565
13.75	453.7	1.24	456.80	0.562	14.68	458.9	1.46	460.17	0.565
13.87	454.8	1.35	458.09	0.562	14.69	458.7	1.43	460.17	0.565
13.88	454.9	1.36	458.18	0.562	14.70	458.6	1.41	460.17	0.565
13.89	455.0	1.35	458.28	0.562	14.71	458.4	1.38	460.17	0.565
13.95	455.5	1.25	458.75	0.562	14.72	458.3	1.36	460.18	0.565
13.96	455.6	1.23	458.83	0.562	14.73	458.1	1.33	460.18	0.565
14.04	456.9	1.40	459.34	0.546	14.74	457.9	1.31	460.18	0.565
14.05	457.1	1.42	459.39	0.546	14.77	457.5	1.23	460.19	0.565
14.06	457.2	1.41	459.45	0.546	14.78	457.3	1.20	460.19	0.565
14.07	457.4	1.39	459.50	0.546	14.80	457.1	1.28	460.20	0.565
14.09	457.7	1.36	459.61	0.546	14.81	456.9	1.32	460.21	0.588
14.10	457.8	1.35	459.66	0.546	14.82	456.8	1.36	460.21	0.588
14.19	459.1	1.22	459.95	0.546	14.83	456.4	1.36	460.22	0.588
14.21	459.1	1.27	460.00	0.543	14.84	456.1	1.36	460.23	0.588
14.23	459.1	1.33	460.02	0.543	14.85	455.7	1.36	460.24	0.588
14.25	459.2	1.38	460.05	0.543	14.86	455.3	1.36	460.24	0.588
14.26	459.2	1.41	460.06	0.543	14.87	454.9	1.36	460.25	0.588
14.27	459.2	1.43	460.07	0.543	14.88	454.6	1.36	460.26	0.588
14.28	459.2	1.46	460.09	0.543					

Table IIa. Decay data used at the 93 Nb(n, 2n) 92m Nb cross section determination.

Reaction	T _{1/2}	E _□	Y□		
product		[keV]	[%]		
^{92m} Nb	10.15 d 2	934.4	99.15 <i>4</i>		

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