# ENDF/B-III Cross Section Measurement Standards* M. K. Drake ${ }^{\text {rt* }}$ <br> National Neutron Cross Scction Center Brookhaven National Laboratory Upton, New York 11973 

July 1972

## Contributors

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## I. INTRODUCTION

Members of the Cross Section Evaluation Working Group (CSEWG) subcommittee on Normalization and Standards have made a study of the various cross sections that are considered to be standards for cross section measurements. Various members of the Subcommittee evaluated the best available experimental data for the standards. The Subcommittee carefully reviewed the evaluations and recommended that certain sets of evaluated data be incorporated into the ENDF/B-III Library and that these data be considered as standards for future work in the area of experimental measurements and evaluations.

This report contains brief sumaries describing the various cross section standards. In certain cases only limited docunentation was available. Detailed documentation, when it was available, is given in the Appendices.

The Normalization and Standards Subcommittee have made recommendations on seven standards. Not all of the standards are known to the desired accuracy. Table 1 shows the various standard cross sections which have been incorporated into the regular ENDF/B-III materials. Each material is complete in that it contains all significant partial cross sections covering entire energy ranges and all secondary angular and energy distributions.

Section II of this report contains a brief description of each standard. Some graphical displays are included which show the recommended cross sections and recent sets of experimental data. In some instances (e.g., ${ }^{6}$ Li, ${ }^{197}$ Au, and ${ }^{235} \mathrm{U}$ ) experimental data not shown on these figures was given the greatest weight in the evaluations. Hence the evaluations, considering all axporimental data, are far better than these figures might suggest. Figures 1 and 2 show the recomuended standard cross sections.

*Desired accuracy as requested by the CSEWG Normalization and Standards Subcommittee and the U.S. Nuclear Data Committee (the USNDC was formerly the Nuclear Cross Section Advisory Committee (NCSAC) which was formerly the Neutron Cross Section Advisory Group (NCSAG).)

## II. DESCRIPTION OF RECOMMENDED STANDARDS

A. Hydrogen Total and Differential Elastic Scattering Cross Sections
 hydrogen. This evaluation was done by Stewart, La Bauve, and Young. The recommended total cross section covers the energy range from 0.1 to 20 MeV . The recommended differential elastic scattering cross section is given in File 4 and covers the energy range from 3 to 20 MeV .

The total cross section was taken to be the sum of the elastic scattering and radiative capture cross sections. The elastic scattering cross section was obtained from a theoretical analysis of measurements made by llopkins and Breit. (2) A set of phase shifts obtained by Seamon et al. (3) was used in this analysis. The differential angular distributions were obtained from the same analysis.

Figure 3 shows a comparison between the recommended total cross section and the results from a measurement made by Davis and Barschali. (4) The uncertainty in the recommended total cross section is believed to be about one percent.
B. ${ }^{3} \mathrm{He}(n, p)$ Cross Section

MAT $=1146$ contains the evaluation of ${ }^{3}$ He cross sections by L. Stewart (LASL). Although this material contains all of the cross section data for $3^{H e}$, only the ( $n, p$ ) cross section from $1.0 \times 10^{-5} \mathrm{eV}$ to 50.0 keV represents the recommended standard cross section.

Below one keV, the ( $n, p$ ) cross section was taken to be $1 / v$ and normalized to 5327 barns at 0.0253 eV . This value was taken from a measurement by Als-Nielsen and Dietrich. (5) The thermal cross section is believed to be known to within one percent.

Between 50 keV and 1 keV , the recommended standard cross section was ' $\mathrm{H} \cdot \mathrm{A}$ on the corrected results from a measurement made by Gibbons and

Macklin. (6) Below 10 keV , the recommended ( $n, p$ ) is believed to be known to within $3 \%$. Above 10 keV , the uncertainty in the recommended cross section rapidly increases beyond the desired three percent.

Figure 4 shows the recommended cross section.
C. ${ }^{6}{ }^{\text {Li }(n, \alpha)}{ }^{3} \mathrm{H}$ Cross Section

The recommended standard cross section for ${ }^{6} \operatorname{Li}(n, \alpha)$ covers the energy range from $1.0 \times 10^{-5} \mathrm{eV}$ to 200 keV , and is contained in MAT $=1115$. Above 200 keV the cross section is not known to the accuracy required to be considered a cross section standard. The recommended cross section was taken directly from an analysis by Uttley et al. (7) of . the total, elistic scattering and ( $n, \alpha$ ) cross section data. Up to 100 eV the cross section was calculated from the formula

$$
\sigma(n, \alpha)=(149.56 / \sqrt{E})-0.024 b .
$$

The $2200 \mathrm{~m} / \mathrm{sec}$. cross section value of 940.25 b was based on a total cross soction measurement of Utteley and Diment. ${ }^{(8)}$

Between 100 eV and 500 keV , the recommended ( $n, r$ ) cross section was obtained by a resonance parameter fit ${ }^{(8)}$ to the best measured results for the total, elastic scattering, and ( $n, \alpha$ ) cross sections. As can be seen in Figure 5, the fitted curve is not in good agreement with the available experimental results for the ${ }^{6} \mathrm{Li}(\mathrm{n}, \alpha)$ reaction.

The $2200 \mathrm{~m} / \mathrm{sec}$. value is believed to be known to within $0.5 \%$. Between thermal and 10 keV the uncertainty is about one percent and increases to two percent at 100 keV .

Although the Uttley and Diment analysis produced a total cross section that was in excellent agreement with the available experimental data,
the disagreement in the $(n, \gamma)$ cross section across the 247 keV resonance precludes the ( $n, \alpha$ ) cross section from being considered a standard above 100 keV at this time.
D. ${ }^{10} B(n, \alpha)$ Cross Section

The ${ }^{10^{-}}(n, \alpha)$ standard cross section is part of MAT $=1155$; the standard is given for the energy range from $1.0 \times 10^{-5} \mathrm{ev}$ to 100 keV . The recommended cross section was based on an analyses by Sowerby et al. ${ }^{(9,10)}$ The $2200^{\circ} \mathrm{m} / \mathrm{sec}$. value of the $(n, \alpha)$ cross section was taken to be $3836.45 \mathrm{~b}^{(11)}$ and this value is believed to be known to within one percent.

Above the thermal neutron energy range, the recommended cross section was obtained from the formula $(9,10)$
$\sigma_{n, \alpha}=\frac{13.837}{\sqrt{E}}-0.312-1.014 \times 10^{-2} \sqrt{E}+\frac{2.809 \times 10^{5}}{\sqrt{E}\left[(170.3-E)^{2}+2.243 \times 10^{4}\right]}$ where $\sigma$ is given in barns and $E$ in keV.

Sowerby et al. obtained the constants for the above equation by making a least squares fit to selected data sets. The absorption cross section measured by Mooring, Monahan and Huddleston ${ }^{(12)}$ was used. The absorption cross section resulting from subtracting the scattering cross section measured by Asami and Moxon ${ }^{(13)}$ from the total cross section measured by Diment ${ }^{(14)}$ also was used. The most important data set used was the ${ }^{10} B_{(n, \gamma)}$ cross section derived from the Sowerby et al. ${ }^{(9,10)}$ ${ }^{6} \mathrm{Li}(\mathrm{n}, \alpha) /{ }^{10} \mathrm{~B}(\mathrm{n}, \alpha)$ ratio measurement. In this case the ${ }^{6} \mathrm{Li}(\mathrm{n}, \alpha)$ cross section was taken to be the same cross section as was used in ENDF/B-III for the standard. The resulting recommended cross sections for ${ }^{6} \mathrm{Li}(\mathrm{n}, \alpha)$ and ${ }^{10} B(n, \gamma)$ cross section contain a high degree of consistency.

The recommended standard cross section is believed to be known to within one percent for neutron energies less than one keV. This uncertainty rises from two percent at 10 keV to three percent at 100 keV .
E. ${ }^{12} \mathrm{C}$ Total Cross Section

The ${ }^{12} \mathrm{C}$ total cross section from $1.0^{-5} \mathrm{eV}$ to 2.0 MeV has been designated as a cross section standard. The differential angular distribution data for elastic scattering is also a preferred standard, however, the ENDF/B data is not considered adequate as a standard at this time. The recommended cross section is part of MAT $=1165$ and was taken from an analysis made by Francis et al. (15)

The recommended free atom scattering cross section is 4.729 barns. This value is believed to be known to within 0.5 percent.

Francis et al. (1.5) obtained the recommended cross section by analyzing several recent experimental results. $16,17,18,19) \quad$ R-matrix and coupled channel model calculations were used in the analysis. The recommended cross section is believed to be accurate to within one percent and likely to within 0.5 percent.
F. ${ }^{197} \mathrm{Au}(\mathrm{n}, \gamma)$ Cross Section

MAT $=1166$ contains the recommended standard cross section for 197
$\mathrm{Au}(\mathrm{n}, \mathrm{y})$ which extends from 10.0 keV to 5.4 MeV . The resonance region radiative capture cross section for $M A T=1166$ has not been evaluated in a manner necessary to be considered a standard cross section. The recommended cross section for the above energy range was obtained by F. J. Vaughn and H. A. Grench ${ }^{(20)}$ (Lockheed Palo Alto Research Laboratory).

Vaughn and Grench analyzed the results from 15 differential cross section measurements. ${ }^{(21,35)}$ In addition they used the results from

19 different monoenergetic measurements. The results from each experiment was examined and the data sets were renormalized, when necessary to take into account information which became available after publication. Renomali\%alions were based on new information about the standards used, in particular the ${ }^{235}$ U fission cross section.

The complete energy range was divided into three overlapping energy intervals, i.e., $10-150 \mathrm{keV}, 123-560 \mathrm{keV}$, and $400-1800 \mathrm{keV}$. A least squares polynomial fit was made for each energy interval. An iterative procedure was used to obtain a "best curve" through the experimental data. In the iterative procedure, an adjustment factor for each data set was obtained. Each point in a particular data set was multiplied by its adjustment factor. The adjusted data sets were re-fit and new adjustment factors obtained. This process was cnntinued until the results converged. For example, the final adjustment factors for the data sets measured by Fricke et al., (21) Poenitz, (22) and Kompe ${ }^{(35)}$ were $1.01567,1.07871$, and 1.11054 respectively.

Figure 6 shows the recommended standard cross section for the energy range from 10 keV to 5 MeV . Also shown are selected sets of experimental results (unadjusted).
G. 235 Fission Cross Section

The CSEWG Normalization and Standards Subcommittee has recomended that the ${ }^{235} U$ fission cross section from MAT $=1157$ be considered as an interim standard for the energy range from $1.0 \times 10^{-5} \mathrm{eV}$ to 15 MeV . It is difficult to estimate the uncertainty in the recomended cross section, but it is considerably greater than the preferred uncertainty of one percent.

The evaluation of the fission cross section for the thermal energy region was made by Leonard. (36) The $2200 \mathrm{~m} / \mathrm{sec}$ values for the fission cross
section and the "Westcott-g factor" were taken from the 1969 IAEA review, ..... (37)
i.e., $\sigma_{f}=580.2$ barns and $g_{f}=0.9768$.
The recommended single level resonance parameters were obtained by
Smith and Young ${ }^{(38)}$ for the neutron energy range from 1.0 eV to 82 eV .
Several sets of experimental data were used in the fits, includingde Saussure et al., (39) Michaudon, (40) Blons et al., (41) and Cao et al. (42)In the energy range from 80 eV to several keV the fission cross
section was based on the data of de Saussure et al. (39) Between several
keV and 100 keV the results measured by White, (43) Perkin, (44) and
Silver et al. ${ }^{(45)}$ were used for absolute values. A fine structure was
introduced into the recommended data. This structure was taken from a
measurement by Lemley et al. ..... (46)Above 100 keV , the recommended data were based on measurements byWhite, (43) Szabo, (47) and Smith et al. (48) (as corrected by Hansen in
1968). Figures 7,8 , and 9 show the recommended fission cross section
above 4 keV and selected sets of experimental results.

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Figure 1: Cross Section Standards


Figure 2: Cross Section Standards

Figure 3: ${ }^{1}$ H TOTAL CROSS SECTION


Figure 4: ${ }^{3} \mathrm{He}(n, p)$ CROSS SECTION




Figure 6: ${ }^{197}$ Au ( $n, \hat{y}$ ) Cross Section


Figure 7: ${ }^{235} \mathrm{U}$ Fission Cross Section



Appendix A

## Hydrogen

# LOS ALAMOS SCIENTIFIC LABORATORY of the University of California <br> los alamos • New mexico 

# Evaluated Nuclear Data for Hydrogen in the ENDF/B-II Format 

by
L. Stewart
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P. G. Young

## EVALUATED NUCLEAR DATA FOR HYDROGEN IN THE ENDF/B-II FORMAT

## by

L. Stewart, R. J. LaBauve, and P. G. Young

## ABSTRACT

The following nuelear data are given for hydrogen in the energy range from $1.0 \times 10^{-5} \mathrm{eV}$ to 20.0 MeV .
File 1 . The general information file includes a brief description of the data to follow.
File 2. Values for nuclear spin and effective scattering radius are given in the resonance file.....
File 3. Smooth cross-section data are given for the total cross section, the free-atom elastic scattering cross section, and the radiative capture cross section; data for $\bar{\mu}, \xi$, and $\gamma$ are also included.
File 4. The angular distributions for elastic scattering are given as probability vs cosine of the scattering angle.
File 7. The r reenatom-scattering cross section is the only information provided at thermal.
File 12. Secondary gama-ray production multiplicities for capture, which are equal to one, are given in this file.
File 14. Gamma-ray angular distributions are provided for the single radiative capture gamma ray.

## INTRODUCTION

This evaluation for hydrogen (MAT $=1148$ ) difFers fram the previous ENDF/B evaluation (MAT = 1001) in that the elastic scattering data were taken from recent work by Hopkins and Breit and the data for radiative capture were taken from recent work by Horsley. ${ }^{2}$ Also, gamma-ray production data, not given in the MAT $=1001$ evaluation, are included. A complete listing for MAT $=1148$ is given in the Appendix.

FILE 1: GENERAL INFORMATION
A brief summary of the data to follow is given in file l. The atomic mass for hydrogen was taken to be 1.007825 from the May 1969 "Chart of the Nuclides. ${ }^{13}$

FILE 2: RLSONANCL Information
Nulloat rain and eflective ocultering radius are given in thio lile. An effective scattering
radius of $1.2756 \times 10^{-12} \mathrm{~cm}$ is consistent with a potential scattering cross section of 20.449 b , as determined from $4 \pi a^{2}$. Singlet and triplet scattering radii are not included.

## FILE 3: SMOOTH CROSS SECTIONS

Total cross sections ( $M T=1$ ) were obtained by adding the elastic scattering and radiative capture cross sections at all energies $\left(1.0 \times 10^{-5} \mathrm{eV}\right.$ to $\mathbf{2 0 . 0} \mathrm{MeV}$ ). The hydrogen total cross sections are shown in Fig. 1.

The elastic scattering cross sections (MT = 2) were taken from an extensive theoretical treatment of fast neutron measurements by Hopkins and Breit.' In this work, a consistent set of cross sections and angular distributions were obtained by using a set of phase shifts previously determined at Yale Universily.' tabular valuer of the "ldelit watll. ing cross section are given in Ref. I tor unly a
few energies, the two lowest points being 100 and 200 keV . The phase shift program and the Yale phase shifts were provided by Hopkins' so that many intermediate points could be calculated. At 0.1 $k \cdot V$, the lowest energy recommended for running this program, the scattering cross section is 20.4488 b . This value is in excellent agreement with the thermal cross section ( $20.442+0.023$ b) derived by Davis and Barschall ${ }^{5}$ from a revised value of the effective range obtained by determining the best values of the neutron energies from many experiments below 5 MeV performed since 1950. Therefore, for this evaluation, the free-atom-scattering eross section is assumed to be constant below 100 eV and equal to the value calculated from the Yale phase shifts at 100 eV , giving a thermal cross section of 20.449 b . At higher energies, these theoretical predictions are in excellent agreement with the recent measurements of Davis ${ }^{6}$ giving an average value of 0.84 for the square of the deviation for energies below 20.0 MeV . The elastic cross section for hydrogen from $1.0 \times 10^{-5} \mathrm{eV}$ to 20.0 MeV is shown in Fig. 2.

The cross sections for radiative capture (HT $=$ 102) were taken from the 1966 publication of Horsley, ${ }^{2}$ where a value of 332 mb was adopted for the thermal value. Deuteron photodisintegration cross sections were also employed in deriving radiative capture in Horsley's report. Although the Nuclear Data article by Horsley ${ }^{2}$ was referenced for $M A T=1001$, the values were taken from an early version described in AWRE 0-23/65, and these were later revised for the Nuclear Data article. The latter report (Ref. 2) has been used for this evaluation, as suggested by Horsley. The radiative capture cross section for MAT $=1148$ From $1.0 \times$ $10^{-5} \mathrm{eV}$ to 20.0 MeV is shown in Fig. 3.

The average value of the cosine in the laboratory system ( $\bar{\mu}_{L}$ ) for elastic scattering (MT = 251) was derived from the secondary angular distributions in file $4(M T=4)$. Values for $\bar{\mu}_{L}$ from $1.0 \times$ $10^{-5} \mathrm{eV}$ to 20.0 MeV are shown in Fig. 4 .

Values for $\xi$, the average logarithmic energy change per collision ( $M T=252$ ), and for $\gamma$, the Goertzel-Greuling constant ( $M T=253$ ), are taken equal to 1 over the range $1.0 \times 10^{-5} \mathrm{eV}$ to 20.0 MeV , following the $M T=1001$ evaluation.

## FILE 4: SECONOARY ANGULAR DISTRIBUTIONS

Anguiar distributions of secondary mutrons resulting from elastic scattering are tabulated fron $1.0 \times 10^{-5} \mathrm{eV}$ to 20.0 MeV . Distributions at $0.1,5$, 10,20 , and 30 MeV are provided by Ref. 1; additional and intermediate data were calculated by using the Hopkins-Breit phase shift program and the Yale phase shifts. As shown in Figs. 5 through 16 , the angular distributions above 100 keV are neither isotropic below 10 MeV , nor are they symmetric aboul $90^{\circ}$ at higher energies as assumed in the earlier version (MAT $=1001$ ). At 100 keV , the angular distributions are assumed to be isotropic because the $180 / 0^{\circ}$ ratio is very nearly unity ( 1.0011 ). At 500 keV, this ratio approaches 1.005 ; therefore, the pointwise normalized probabilities as a function of the cosine of the scattering angle are provided at $1.0 \times 10^{-5} \mathrm{eV}$ (isotropic), 100 keV (isotropic), 500 keV , and at $1-\mathrm{MeV}$ intervals from 1 to 20 MeV .

File S: thermal data
Free-atom cross sections specified from 1.0 x $10^{-5} \mathrm{eV}$ to 5 eV are included in this file.

FILE 12: PHOTON PRODUCTION CROSS SECTIONS
A multiplicity representation is used to describe the single hydrogen radiative capture gamma ray from $1.0 \times 10^{-5} \mathrm{eV}$ to 20.0 MeV . The multiplicity is referred to MT $=102$ in File 3 and is unity at all neutron energies. To adequately represent the gamma-ray energy for MeV-incident neutrons, the neutron energy region from 0.2 to 20 MeV is divided into 16 different energy bands, and the gamma-ray energy is tabulated for each neutron energy band as

$$
\bar{E}_{Y}=2.225 \times 10^{6}+\bar{E}_{n} / 2 \quad \text { (eV) , }
$$

where $\vec{E}_{n}$ is the neutron energy at the midpoint of the band in eV. The value $2.225 \times 10^{6} \mathrm{eV}$ corresponds to the deuteron binding energy; that is, the small energy change due to the nuclear recoil that accompanies gamma emission has been ignored.

## FILE 14: GAMMA-RAY ANGULAR DISTRIBUTIONS

The gamma-ray angular distributions are assumed to be isotropic at all neutron energies from $1.0 \times$ $10^{-5} \mathrm{eV}$ to 20.0 MeV .


Fis. 1. Fotal cross suction (MT $=1$ ).


Pig. 2. Elastic scattering cross section (MT = 2).
average value of mu $\left(10^{-2}\right)$





Fis. $i$. Ansular distribution for L.ll M.l.






Fif. 12. Angular alistribution for 12.0 NeV.



Fig. I6. Anguiar distribution for 20.0 MeV .

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APPENDIX
LISTING OF HYDROGEN EVALUATION





| 1.9000F+06 | $3.0030 \mathrm{E}+00$ | 2.0000E*06 |  |  |  |  | 120 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ?.4000E+06 | 2.6220F-10 |  |  | 06 | $2.3920 E+001148$ | 31 | 121 |
| $3.0000 \mathrm{E}+06$ | 2.2930F. 00 | 3.2000E +06 | 2.2036. 64 | - $000 \mathrm{E}+06$ | 2.1200E +001148 | 31 | 122 |
| 3.6000E+06 | $2.0430 \mathrm{E}+00$ | 3. $\mathrm{AOOOE}+06$ | 1.97:30. | -. 0 OOOF +06 | $1.9070 \mathrm{E} * 001148$ | 31 | 123 |
| 4.2000F+06 | 1.8450E + 00 | 4.4000E.06 | 1.7EBúE | 4. $5000 \mathrm{E}+06$ | $1.7340 \mathrm{E}+00114$ | 31 | 124 |
| $4.8000 E+06$ |  | 5.0000F-S |  | $5.2000 \mathrm{E}+06$ | $1.5890 E+0011$ | 31 | 25 |
| $5.4000 \mathrm{E}+06$ | $1.5470 \mathrm{E}+00$ | $5.6000 \mathrm{E}+06$ | 1.53600-00 | 5. $4000 \mathrm{E}+06$ | 1.4670F+001148 | 31 | 126 |
| $6.0000 \mathrm{E}+06$ | 1.4300F-00 | 6.2000E-06 | $1.3950 F+00$ | 6.4000E + 06 | 1.3620E 001148 | 31 | 127 |
| $6.6000 F+06$ | 1.3290E | O. $\mathrm{HOOOE}+06$ | 1.2990E +0 | 7.0000 E | 1.2690 F | 31 | 128 |
| 7.5000F+06 | 1.2010 | $8.00005+06$ | $1.1390 f+00$ | 8.b $000 \mathrm{E}+06$ | 1 | 31 | 129 |
| Q.0nOOE - On | $1.0720 \mathrm{~F}+00$ | 9.5000F-06 | 9. $\mathrm{H} 590 \mathrm{E}=01$ | 1.0000F-07 | 9.4320E-011148 | 31 | 130 |
| 1.0500F-07 | 9.0.350F-01 | $1.1000 E+07$ |  | 1.15UUF+07 | B.3230F-011148 | 31 | 131 |
| 1.2000F+07 | R.cns0f-n1 | $1.2500 F+07$ | $7.7100 F-01$ | $1.3000 \mathrm{E}+07$ | $7.4130 \mathrm{~F}-01114 \mathrm{~A}$ | 31 | 132 |
| 1.3500F407 | 7.1730F-01 | 1.4000F-07 | 6.4290E-01 | 1.4SU0F. 07 | 6.6480E-011148 | 31 | 133 |
| 1.5000F-07 | 6. 4 R $00 F-01$ | $1.5500 \mathrm{~F}+07$ | 6.2740F-01 | $1.6000 F \cdot 07$ | 6.0780E-011148 | 31 | 134 |
| OOF +07 | S.8970E-01 | 1.7000E+07 | 5.7170F-01 | 1.7500F. 07 | 5.549nE-01114A | 31 | 135 |
| 1. $1.9 \cap O E * O 7$ | 5.3900E-01 | $1 . H 500 E+07$ | 5.2340E-01 | $1.9000 F+0$ | 5.0430E-01114H | 31 | 136 |
| 1.9500E-07 | $4.9550 \mathrm{E}-01$ | $2.0000 \mathrm{~F}+07$ |  |  | 0.114 d | 31 | 137 |
|  |  |  |  |  | 11 | 30 | 131 |
| 1.0 |  |  |  |  | 0114 A | 32 | 134 |
| 0.1 | 0. | 0 | 0 |  | 14114H | 37 | 1411 |
|  | 5 |  |  |  | 114 H | 32 | 41 |
| 1.0000F-05 | 2.0444 E 01 | ל | 2.0449E. 0 |  | 2.0449F-01114A | 32 | 142 |
| $1.0000 \mathrm{~F}-04$ | 2.0444E + 01 | 2.0000F-0 | $2.0449 \mathrm{E}+0$ | 0000F-04 | 449F-011148 | 7 ? | 143 |
| $1.0000 \mathrm{~F}-03$ | ?.044YE+01 | 2.0000E-03 | 2.0 | 5.0000F-03 | 2.0444F-01114 | 32 | 144 |
| $1.00005-0$ ? | 7.0449F-01 | 2. 4 300E-02. | 2.0449E-01 | $1.0000 F+02$ | 2.0444F.01114 | 32 | 145 |
| $1.0000 \mathrm{E} \cdot 03$ | 2.0329F-01 | C.0000E+1)3 | 2.0198F+01 | 3.0noof 03 | 2.0116.HF-01114H | 3 ? | 146 |
| $4.0000 \mathrm{~F}+03$ | 1.9941F.01 | $5.0000 \mathrm{~F}+03$ | 1.4 Hl 15F*01 | h.GOUOF. 0.3 | 1.9691F. | 32 | 147 |
| $8.0000 \mathrm{~F} \cdot 03$ | $1.944 \mathrm{BE}+01$ | $1.0000 E+04$ | $1.9213 F+0$ | 1.5000F-04 | +1.8051F+011149 | 32 | $14 R$ |
| $2.0000 \mathrm{E}+04$ | 1.8120E+01 | 2.5000E+04 | 1.7634F. 01 | $3.0000 \mathrm{~F}+04$ | 1.7172t 011148 | 32 | 9 |
| 3.5000F+04 | 1.6737E+01 | $4.0000 E+04$ | $1.6327 E+01$ | 4.5000F.04 | 1.5441E+011148 | 32 | 150 |
| $5.0000 \mathrm{E}+04$ | $1.5575 \mathrm{E}+01$ | $5.5000 \mathrm{E}+04$ | $1.5228 E+01$ | 6.0000F+04 | 1.4900F-011148 | 32 | 151 |
| 6.5000F+04 | $1.4587 E+01$ | 7.0000E-04 | 1.4291F-01 | 7.5000E +04 | 1.4(108F. 01114 H | 32 | 152 |
| 8.0000E-04 | $1.3738 \mathrm{E}+01$ | C. $5000 E+04$ | $1.34 \mathrm{HIF}+01$ | 9.0000E-04 | 1.3235E+01114 | 32 | 153 |
| $9.5000 \mathrm{~F}+04$ | 1.2999E*01 | $1.0000 \mathrm{E}+05$ | 1.2774E+01 | $1.1000 F+05$ | 1.2351F.011144 | 32 | 154 |
| $1.2000 \mathrm{E}+05$ | $1.1964 \mathrm{E}+01$ | $1.3000 \mathrm{E}+05$ | $1.1607 E+01$ | $1.4000 E+05$ | 1.1275F-01114H | 32 | 155 |
| $1.5000 \mathrm{~F}+05$ | 1.0965E+01 | 1.A000E+05 | $1.0673 \mathrm{~F}+01$ | $1.7000 F+05$ | 1.0.198E*01114 | 32 | 156 |
| 1. $8000 \mathrm{~F}+05$ | $1.0140 \mathrm{~F}+01$ | 1.4000F.05 |  | 2.0nOOF + 05 | 9.6710F+00114 | 32 | 157 |
| 7.2 | 9. $3580 F+00$ | 2.4000E 05 | H. $\mathrm{H92OF}+00$ | 2. $2000 \mathrm{~F}+0 \mathrm{~b}$ | H.5b20f.0 | 32 | 15 H |
| P. $\mathrm{AOOOF}+05$ | H. ThPOF + no | . $3.0000 \mathrm{~F}+05$ | 7.9870F-00 | 3-2nuor 05 | 7.7340F+001144 | 3 ? | 154 |
| $3.4000 \mathrm{~F} \cdot 05$ | 7.5n10F•00 | 3.ACOOE 05 | 7. $2 \mathrm{H} 40 \mathrm{E} \cdot 00$ | 3. HCOOE-OS | 7.0030F-001148 |  | 160 |
| $4.0000 \mathrm{~F} \cdot 05$ | 6. $\mathrm{B} 970 \mathrm{t} \rightarrow 0 \mathrm{O}$ | 4.20005 | $6.7350 F \cdot 00$ | $4.4 n U O F+05$ | B.5050F-00114 |  | 161 |
| $4.4000 \mathrm{~F}+05$ | 6.4150F.nn | $4 . H 000 E+05$ | -. $2750 \mathrm{C}+00$ | 3.0000F.05 | 6.1430F400114 | 3 ? | 16\% |
| 5.50010 OF | 5.8450E•00 | $6.0000 F+05$ | $5.5440 F \cdot 00$ | 6.b000F 0 O | $5.3540 F+00114 \mathrm{H}$ |  | 163 |
| 7.0nODF*05 | $5.14 \mathrm{AOE} \cdot 100$ | 7.5000E +05 | $4.9640 E+00$ | H.0000F.05 | 4.7470E +001148 |  | 164 |
| H.5000E-05 | 4.6450F. 00 | 9.0000E-05 | $4.5060 F+00$ | 9. $5000 \mathrm{E}+05$ | $54.3180 E+001148$ | 32 | 105 |
| 1. 0 ODOE +06 | 4.2610E.On | 1. 1000E 06 | $4.0510 \mathrm{E}+00$ | 1.2000E +06 | 3. H (GOEE*00114 | 32 | 156 |
| $1.3000 \mathrm{~F}+06$ | 3.70R0E-00 | $1.4000 E+06$ | $3.5610 E \cdot 00$ | $1.5000 E+00$ | 3.4290E*0 | 32 | 167 |
| 1.6000E+06 | 3. $3090 \mathrm{E}+00$ | $1.7000 E+06$ | 3.1980E 00 | 1.8000E*06 | 3.0470E +001148 | 32 | 168 |
| 1.9000F+06 | $3.0030 \mathrm{E}-00$ | 2. $0000 \mathrm{E}+06$ | 2.9150E+00 | 2.2000E+0G | 2.7590F+001148 | 37 | 169 |
| 2.4000F*06 | 2.5220E*00 | 2.6000E*06 | $2.5010 E+00$ | $2.8000 E+06$ | 6-3420E*001148 | 32 | 170 |
| 3.0 OnOF 06 | 2.2930E-00 | $3.2000 \mathrm{E}+0$ | 2.2030E*00 | $3.4000 \mathrm{E}+06$ | 2.120nE*001148 | 3 | 171 |
| $3.60005+06$ | 2.0430E-00 | 3.4000E*06 | 1.9730E-00 | 4.0000F*06 | 1.9070F-001148 |  | 172 |
| 4. 2 OOOE +06 | 1. A450F. 00 | $4.4000 F \cdot 06$ | 1.7RROE 00 | 4.6000F. 06 | $61.7340 E+001148$ | $\cdots$ | 17.7 |
| 4. ROOOF OG | $1.6830 E+00$ | S.UOROE 06 | 1.6.350E $\cdot 00$ | 5. $2000 \mathrm{~F}+06$ | 1.5890F+001148 | 3 | 174 |
| $5.4000 \mathrm{~F} \cdot 0 \mathrm{O}$ | $1.5470 \mathrm{E}+00$ | $5.0000 F \cdot 06$ | $1.5060 F \cdot 00$ | $5.8000 \mathrm{f}+06$ | 61.4670F+00114R | 32 | 115 |
| $6.0000 \mathrm{~F}+06$ | $1.4300 E \cdot 00$ | $6.2000 E+06$ | 1.3950F-00 | $6.4000 F+06$ | 61.3020F+00114R | 2 | 17 H |
| G. $\mathrm{BOQOOF}+$ Of | 1.3290E+ 00 | 6. $\mathrm{HOCOF}+06$ | $1.2990 \mathrm{E} \cdot 00$ | $7.0000 \mathrm{~F}+06$ | $61.2040 F+00114 \mathrm{~B}$ | 32 | 177 |
| 7.5000F*06 | 1-2010F.00 | S.0000F*06 | $1.1390 \mathrm{~F}+00$ | 8. $5000 \mathrm{~F}+06$ | 1.OB30E-00114 | 32 | 178 |
| $9.0000 \mathrm{~F} \cdot 06$ | 1.0320F-00 | 9.5000F+06 | $9.85901-01$ | $1.0000 F+07$ | 9.43?0F-01114 | 32 | 179 |
| $1.0500 \mathrm{~F}+07$ | 9.0.750E-01 | $1.1000 \mathrm{~F} \cdot 07$ | H.6650F-01 | $1.1500 \mathrm{E}+07$ | 8. $3630 \mathrm{~F}-01114$ |  | 180 |
| 1.2000F407 | 8.0050F-01 | $1.2500 t+07$ | 7.7100E-01 | 1.3000F-07 | $77.4330 \mathrm{~F}-01114 \mathrm{H}$ | . 2 | 181 |
| 1. $3500 \mathrm{~F}+07$ | 7.1770F-01 | 1.4000F-07 | 6.9290E-01 | 1.4500E 07 | 7 h. $6980 \mathrm{O}-01114 \mathrm{~A}$ | 3 2 | 188 |
| 1.5000F. 07 | 6.4800E-01 | $1.5500 \mathrm{E} \cdot 07$ | 6.2740E-01 | $1.6000 \mathrm{E} \cdot 07$ | $7 \mathrm{6.OTHOE}-011144$ | J | 183 |
| $1.6500 \mathrm{E}+07$ | $5.8930 \mathrm{E}-\mathrm{DL}$ | 1.7000E*07 | 5.7170E-01 | 1.7500F. 07 | $75.5490 \mathrm{E}-011148$ | , | 184 |
| $1.8000 E+07$ | 5.3900F-01 | 1. $\mathrm{CSOOE}+07$ | 5.2380E-01 | $1.4000 F+07$ | 75.0930F-01114 | 32 | 145 |
| $1.95005+07$ | 4.9550F-01 | 2.0000E +07 | 4. | , | -0. 1149 | 2 | 186 |
|  |  |  |  |  | 1148 | 30 | $1 \mathrm{H}_{7}$ |
| 001 E+03 | $9.9917 \mathrm{E}-01$ |  | 0 |  | 01148 | 3102 | 1 Hf |
| 0.0 | 2.2247 E406 | 0 |  |  | 134114 A | 3102 | 189 |
| 134 |  |  |  |  | 1148 | 3102 | 190 |


| 1.0000 | F－05 | 1．FKGQ4F－01 | 2．0000E－05 | 1．1ROAF 01 | 5.0000 F | 001148 | 3102 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1.0000 | E－04 |  | 2．0000E－04 | $3.7341 \mathrm{~F}+00$ | 5．0000E－04 | 2．3016E＋001148 | 3102 | 142 |
| 1.0000 | F－03 | 1．AROQF +00 | $3.0000 \mathrm{E}-0.3$ | 1．1808F－00 | $5.0000 \mathrm{E}-03$ | 7．46H2E－011148 | 3102 | 193 |
| 1．000n | F－0？ | 5． $280 \mathrm{HF}-01$ | ？． $5300 \mathrm{OF}-02$ | 3． $3200 \mathrm{~F}-01$ | $1.0000 \mathrm{E}+02$ | S． $27708-031148$ | 3102 | 194 |
| 1.0000 | F +03 | 1．6590F－03 | 2．0000F－03 | 1．19？6F－03 | 3．0000F＋03 | 9．7008F－041148 | 310？ | 195 |
| 4.0000 | OFP 03 | H．3P90F－04 | 5．0000F－03 | $7.3747 \mathrm{E}-04$ | 6．0000F． 03 | G．hti $9 \mathrm{E}-041148$ | 3102 | 196 |
| H．0000 | FF－0．3 | ＇，大51 BF－04 | 1．0000F 04 | 4．95ROF－04 | 1．5000E +04 | 3．87R2E $=041148$ | 3102 | 147 |
| 2．0000 | OF +04 | 3．2386F－04 | 2． $2000 \mathrm{E} \cdot 04$ | 2．ROS4t－04 | $3.0000 E+04$ | 2．4909E－041148 | 3102 | 198 |
| 3.5000 | F +04 | ？．7フロ日F－04 | 4．0000E－04 | 2．0553F－04 | $4.5000 \mathrm{E} * 04$ | 1．8970E－041148 | 3102 | 149 |
| 5.0000 | F +04 | 1．7H46E－04 | 5．3000E•04 | 1．6518E－04 | t． $0000 \mathrm{E}+04$ | 1．5544E－041148 | 3102 | 200 |
| 5.5000 | F +04 | 1．4593E－04 | 7．0000E－ 04 | 1．3942E－04 | $7.5000 E+04$ | 1．3273E－041148 | 3102 | 201 |
| 8.0000 | OF +04 | 1．26738－04 | R． $50005+04$ | 1．2132F－04 | $9.0000 \mathrm{E}+04$ | 1．1640E－041148 | 3102 | 202 |
| 9.5000 | F +04 | 1．1191E－04 | 1．0000E＊05 | 1．0780E－04 | 1．1000E＊05 | 1．0019E－041148 | 3102 | 203 |
| 1.2000 | $F \cdot 05$ | 9．3717E－05 | 1．3000E－05 | 8．8131E－05 | $1.4000 \mathrm{E} \cdot 05$ | 8．325thF－051148 | 3102 | 204 |
| 1．5000 | F－05 | 7．R9GOF－0S | 1． $6000 \mathrm{~F}+05$ | $7.5142 \mathrm{E}-05$ | $1.7000 E+05$ | 7．17PらE－051148 | 3102 | 205 |
| 1.9000 | F +05 | 6．RALSE－05 | $1.9000 \mathrm{E}+05$ | $6.5853 E-05$ | 2．0000E 05 | $6.3110 \mathrm{E}-05114 \mathrm{~A}$ | 3102 | 206 |
| 2.2000 | OF＋ 05 | 5．9051E－05 | $2.4000 F \cdot 05$ | 5．5591E－05 | 2．6700E－05 | $5.2131 E-051148$ | 3102 | 207 |
| 2.8000 | F +05 | $5.0330 E-0.5$ | $3.0000 \mathrm{E}+05$ | $4.8290 E-05$ | 3．2000F．05 | 4．6538E－05114R | 3102 | 20 H |
| 3.4000 | OF－05 | $4.5018 F-05$ | 3． $0000 \mathrm{E}+05$ | 4．3691F－05 | 3．8000E－05 | 4．2523E－05114A | 3102 | 209 |
| 4.0000 | OF +05 | $4.1490 \mathrm{E}-05$ | 4．P000E＋05 | 4．0607F－05 | 4．4000E．05 | 3．9828E－05114E | 3102 | 210 |
| 4.6000 | OF +05 | $3.9138 \mathrm{E}-05$ | $4.8000 \mathrm{E}+05$ | 3．8525E－05 | 5．0nOOE＋OS | 3．7480E－05114R | 3102 | 211 |
| 5.5000 | OF +05 | 3．7396F－05 | 6．0000E +05 | 3．6870E－05 | $6.5000 F+05$ | 3－6163F－05114R | 3102 | 212 |
| 7.0000 | OE +05 | 3．5520E－05 | 7．4000E＋05 | 3．5167E－05 | H．0000E 05 | 3．4840E－05114日 | 3102 | 21.3 |
| H． 5000 | DF－05 | 3．47425－05 | 9．0000E＊05 | 3．4650E－05 | 9．5000E 05 | 3．455？E－051148 | 3102 | 214 |
| 1.0000 | OF－06 | 3．4460F－05 | $1.1000 \mathrm{E}+06$ | 3．4440E－05 | $1.2000 F+06$ | 3．4410F－051148 | 3102 | 215 |
| 1.7000 | Of： 0 n | 3．4400f－05 | 1．4000E＋OK | 1．4360E－05 | $1.5000 \mathrm{~F}+06$ | ．3．4140E－051148 | 3102 | 216 |
| 1．hnon | OF + On | 3．4．310E－04 | 1．1000E＋06 | 3．4290E－05 | 1．8000F－06 | 3．4270F－05114日 | 3102 | c17 |
| 1．9nno |  | 3．4250F－05 | 2．0000E－06 | 3．4230F－05 | 2． $2000 \mathrm{~F}+00$ | 3．4570F－051144 | 3102 | P14 |
| 2.4000 | OF + Of | 3．4910F－05 | $2.0000 \mathrm{E}+05$ | 3．5100E－05 | $2.8000 E+00$ | 3．5300F－05114H | 3102 | 219 |
| 3.0000 | OF－06 | 3．5AROE－05 | $3.2000 E+06$ | $3.5800 \mathrm{~F}-05$ | 3．40U0E－06 | $3.5410 E-05114 \mathrm{H}$ | 3102 | 220 |
| 3.6000 | OF－Ot | $3.6030 \mathrm{~F}-05$ | 3．R000t－06 | 1．h140F－05 | 4．0000E 00 | 3．6260E－051144 | 3102 | 221 |
| 4.2000 | OF +06 | 3．6．690E－05 | $4.4000 E+06$ | 3．6320E－05 | $4.6000 \mathrm{~F}+06$ | 3．6360E－051148 | 3102 | 222 |
| 4.8000 | DE＋ 06 | 3．6390E－05 | 5．0000E＋06 | 3．64？0E－05 | S． $2000 \mathrm{E}+06$ | 3．6290E－051148 | 3102 | 223 |
| 5.4000 | DF + Df | 3．6160E－05 | 5．6000E＋06 | $3.6040 \mathrm{~F}-05$ | 5． $2000 \mathrm{E}+06$ | 3．5410F－051148 | 3102 | 224 |
| 6.0000 | jF +06 | 3．5780E－05 | G．POOOE 06 | 3．5070E－05 | 0．4000E＋ 06 | 3．5SG0E－05114R | 3102 | 225 |
| 6.6000 | DF +06 | 3．5450E－05 | 6． $4000 \mathrm{E}+\mathrm{Oh}$ | $3.5340 \mathrm{~F}-05$ | $7.0000 E+06$ | 3．5230F－051148 | 3102 | 226 |
| 7．50no | OF－06 | 3．4590F－05 | H．0000F－O6 | $3.3940 \mathrm{E}-05$ | H．5000F＋06 | 3．3640E－05114B | 3102 | 227 |
| 9.0000 | OF＊OF | $3.3330 \mathrm{~F}-05$ | $9.5000 E+06$ | $3.2960 \mathrm{~F}-05$ | $1.0000 E+07$ | 3．2590E－05114R | 3102 | 22R |
| 1.0500 | OE＋ 07 | 3． $3710 \mathrm{E}-05$ | 1．1000E＋07 | 3．1820f－05 | 1．1500E．07 | 3．1450F－051148 | 3102 | 229 |
| 1.2000 | OF + n7 | 3．1明EE－05 | $1.2500 E+07$ | 3．0630F－05 | $1.3000 \mathrm{E}+07$ | 3．01ROE－051148 | 3102 | 230 |
| 1.3500 | OF +07 | 3．0010E－05 | $1.4000 \mathrm{E}+07$ | 2．9830F－05 | $1.4500 \mathrm{E}+07$ | 2．9400E－051148 | 3102 | 231 |
| 1.5000 | SF－07 | 2．R9K0E－05 | $1.5500 \mathrm{E}+07$ | $2.8630 \mathrm{~F}-05$ | $1.6000 \mathrm{E}+07$ | $2.8300 \mathrm{E}-051148$ | 3102 | 232 |
| 1.6500 | DE＋07 | 2．7880E－05 | $1.7000 \mathrm{E}+07$ | $2.7450 \mathrm{E}-05$ | 1．7500E＋07 | 2．7360F－051148 | 3102 | 233 |
| 1.9000 | OE＋ 07 | 2．7P60E－05 | $1.8500 \mathrm{E}+07$ | $2.6730 \mathrm{E}-05$ | $1.9000 \mathrm{E}+07$ | 2．6200E－051148 | 3102 | 234 |
| 1.9500 | OE $* 07$ | 2．61POE－05 | 2．0000E＋0？ | $2.6040 \mathrm{~F}-05$ |  | －0． 1148 | 3102 | 235 |
|  |  |  |  |  |  | 1148 |  | 236 |
| 1.001 | F＋03 | $9.9917 \mathrm{E}-01$ | 0 | 0 | 0 | 0114 月 | 3251 | 237 |
| 0.0 |  | 0.0 | 0 | 0 | 1 | 14114 A | 3251 | 238 |
|  | 14 | 3 |  |  |  | 114 A | 3251 | 239 |
| 1.0 | F－05a | H．f5？ $13 F-01$ | 1.0 F．05t | H．Kら213F－01 | 5.0 E．05 | 6．64H99F－011148 | 3251 | 340 |
| 1.0 | $\mathrm{F} \cdot 0 \mathrm{hr}$ | H．h4h20r－01 | $2.0 \quad F \cdot 0$ ¢ | h． 64149 F －01 | $4.0 \quad F \cdot 06$ | 6．63．155F－01114H | 3251 | 241 |
| h． 0 | F．06R |  | H．OE E．Oft | H．R201RF－01 | $1.0 \quad F .07$ | 6．61 138F－011148 | 3251 | 2.42 |
| 1.2 | F．07A | A．A104SE－01 | $1.4 \mathrm{~F}+\mathrm{OH}_{6}$ | A．RO449F－01 | 1．6 F＊07 | 6． $59979 \mathrm{~F}-011148$ | 3241 | 2.43 |
| 1.9 | $\mathrm{E}+07 \mathrm{~b}$ | ．59540E－01 | 2.0 E．07r | ． 5919 GE－01 |  | 1148 | 3751 | 244 |
|  |  |  |  |  |  | 114 H |  | 245 |
| 1.001 | $F \rightarrow 03$ | $9.9917 E-01$ | 0 | 0 | 0 | 01148 | 3252 | 246 |
| 0.0 |  | 0.0 | 0 | 0 | 1 | 21148 | 3252 | 247 |
|  | 2 | 2 |  |  |  | 1148 | 3252 | 248 |
| 1.0 | E－05 | 1.0 | 2.0 F． 07 | 1.0 |  | 1148 | 3252 | 249 |
|  |  |  |  |  |  | 1148 |  | 250 |
| 1.001 | $F+03$ | $9.9917 E-01$ | 0 | 0 | 0 | 0114 H | 3253 | 251 |
| 0.0 |  | 0.0 | 0 | 0 | 1 | 21148 | 3253 | 252 |
|  | 2 |  |  |  |  | 114 A | 3253 | 25.3 |
| 1.0 | F－05 | 1.0 | 2.0 E＊07 | 1.0 |  | 114 A | 3253 | 254 |
|  |  |  |  |  |  | 1148 | 30 | 255 |
|  |  |  |  |  |  | 1148 | 0 | 256 |
| 1.001 | $F \cdot 03$ | $9.9917 F-01$ | 11 | 2 | 0 | 01148 | 42 | 257 |
| 0.0 |  | 0.0 | 0 | 2 | 0 | 01148 | 42 | 254 |
|  | 0.0 | 0.0 | 0 | 0 | 1 | 141148 | 42 | 259 |
|  | 14 | 2 |  |  |  | 1148 |  | 260 |


| 0.0 | $1.0-050$ | 0 | 11114 H | 4 | 2 | 261 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 11 | $?$ |  | 1148 | 4 | 2 | 262 |
| －1．0 | ． 5 －．9 | ． 5 －． 6 | ．5114A | 4 | 2 | 263 |
| ． 4 | .50 .7 | － 50 | ．51148 | 4 | 2 | 204 |
| ．？ | ． 50.4 | ． 5 －6 | －51148 | 4 | 2 | 64 |
| ．${ }^{\text {r }}$ | ． 51.0 | ． 5 | 114 H | 4 | 2 | $26 \%$ |
| 0.0 | 1．0．04 0 | $0 \quad 1$ | 11114 A | 4 | 2 | 267 |
| 11 | $\checkmark$ |  | 114 H | 4 | $?$ | ChH |
| －1．11 | $\therefore \quad-H$ | ． 5 －．6 | ． 51144 | 4 | 2 | chu |
| －． 4 | －－－－ | ． 50. | ． 51148 | 4 | 2 | 270 |
| ．$?$ | .5 ． 5 | .5 ． 6 | ． 51148 | 4 | 2 | 271 |
| ． 8 | .51 .0 | － 5 | 1148 | 4 | 2 | 212 |
| 0.0 | 5.0 .050 | 01 | 11114 H | 4 | 2 | 273 |
| 11 | 2 |  | 114 H | 4 | 2 | 274 |
| －．10n00F－01 | －50117E－00－． $\mathrm{H} 0000 \mathrm{OE}+00$ | ． $50097 \mathrm{E}+00 \mathrm{O}$－GOOUNOF． 00 | ．50066F－00114 | 4 | 2 | 275 |
| $-.40000 \mathrm{~F} \cdot 00$ | $.50045 \mathrm{~F}+00-.70000 \mathrm{E}+00$ | ． $50025 \mathrm{~F}+000$. | ． $50005 \mathrm{~F}+00114 \mathrm{~g}$ | 4 | 2 | 275 |
| －20000F－00 | ． $49974 \mathrm{~F}+00.40000 \mathrm{E}+00$ | ． $49953 \mathrm{E}+00 \cdot+0000 \mathrm{~F}+00$ | ． $49433 \mathrm{E}+001148$ | 4 | 2 | 277 |
| －RODOOF－00 | －49902F－00＝10000E＋01 | ． $49882 \mathrm{SE}+00$ | 114 H | 4 | 2 | 278 |
| 0.0 | 1.0 .060 | 01 | 111148 | 4 | 2 | 279 |
| 11 | $2$ |  | 1148 | 4 | 2 | 240 |
| －． $10000 \mathrm{~F}+01$ | ． $50218 \mathrm{E}+00-.80000 \mathrm{E}+00$ | ． $50176 \mathrm{E}+00-.60000 \mathrm{~F}+00$ | ．50131E－001148 | 4 | 2 | 281 |
| －． $40000 \mathrm{~F}+00$ | －50089E＋00－． $200000+00$ | $.50044 \mathrm{~F}+000$ | ．50000F－001148 | 4 | 2 | 282 |
| $.20000 F+00$ | ．49957F＋00．40000F＋00 | $.49913 \mathrm{E}+00.60000 \mathrm{~F}+00$ | ． $49869 \mathrm{~F}+001148$ | 4 | 2 | 283 |
| $.80000 E+00$ | －49823E＊00－10000E＋01 | ． $49779 \mathrm{E}+00$ | 1148 | 4 | 2 | 284 |
| 0.0 | $2.0+060$ | 0 | 111148 | 4 | 2 | 285 |
| 11 | 2 |  | 1148 | 4 | 2 | 285 |
| $-.10000 \mathrm{~F} \cdot 01$ | $.50387 \mathrm{E}+00-\mathrm{H0000E+00}$ | ． $50311 \mathrm{E}+00-.60000 \mathrm{E}+00$ | ． $50234 \mathrm{~F}+001148$ | 4 | 2 | 287 |
| －． $40000 \mathrm{~F}+00$ | －50158F＋00－． $20000 \mathrm{E}+00$ | $.50081 F+000$. | ．50003F．001148 | 4 | 2 | 248 |
| ． $20000 \mathrm{~F} \cdot 00$ | ． $49925 \mathrm{~F}+00.40000 \mathrm{E}+00$ | ． $4984 \mathrm{FF}+00.60000 \mathrm{~F}+00$ | ． $49764 \mathrm{E}+00114 \mathrm{H}$ | 4 | 2 | 289 |
| －RODOOF－00 | ． $49682 \mathrm{C}+00 \cdot 10000 \mathrm{~F}+01$ | ． $44598 \mathrm{~F}+00$ | 1148 | 4 | 2 | 290 |
| 0.0 | 4.0 －06 0 | $0 \quad 1$ | 111149 | 4 | ？ | 291 |
| 11 | 2 |  | 1148 | 4 | 2 | 292 |
| －． $10000 \mathrm{~F} \cdot 01$ | ． $50677 \mathrm{~F}+00 \times$－ $20000 \mathrm{E}+00$ | ． $50542 \mathrm{E}+00-.60000 \mathrm{~F}+00$ | ．50407F－00114R | 4 | 2 | 293 |
| $-.40000 \mathrm{~F}+00$ | －5027SF＋00－． $20000 \mathrm{FF}+00$ | ． $50143 F+000$. | －bOU1JE＋0ワ114R | 4 | $?$ | 244 |
| －20000F＋90 | ． $49873 \mathrm{E}+00.40000 \mathrm{E}+00$ | ． $49734 \mathrm{~F}+00 \cdot 60000 \mathrm{~F}+00$ | ．49ち89E＋0n114A | 4 | ？ | 245 |
| ． $90000 \mathrm{~F}+00$ | ． $49441 F+00 \cdot 10000 E+01$ | ． $442 \mathrm{CHE}+00$ | 1149 | 4 | ？ | 296 |
| 0.0 | 6．0－06 0 | 01 | 111148 | 4 | 2 | 297 |
| 11 | 2 |  | 1148 | 4 | 2 | 244 |
| $-.10000 \mathrm{~F}+01$ | ． $50999 \mathrm{E}+00-.40000 \mathrm{E}+00$ | ． $50770 \mathrm{E}+00-.60000 \mathrm{E}+00$ | －50560E＋001148 | 4 | 2 | 299 |
| $-.40000 E+00$ | $.50 .362 \mathrm{E}+00 \sim+20000 \mathrm{E}+00$ | ． $50177 \mathrm{~F}+000$. | ． $49493 \mathrm{E}+001148$ | 4 | 2 | 300 |
| － 2000 DF＋00 | ． $4980 B F+00.40000 E+00$ | ． 496 PRE＋ 00 － $60000 \mathrm{E}+00$ | － 49439 F － 0 O1148 | 4 | 2 | 301 |
| － $80700 \mathrm{C}+00$. | $.49246 E+00.10000 E+01$ | ． $49040 \mathrm{~F}+00$ | 1148 | 4 | 2 | 302 |
| 0.0 | $8.0+060$ | $0 \quad 1$ | 111148 | 4 | 2 | 303 |
| 11 | 2 |  | 1148 | 4 | 2 | 304 |
| －． $10000 \mathrm{~F}+01$ | －5128日E＋ $00-80000 \mathrm{E}+00$ | ． $50952 \mathrm{E}+50-.70000 \mathrm{E}+00$ | － $00676 \mathrm{~F}+001148$ | 4 | 2 | 305 |
| －． $40000 \mathrm{~F}+10$ | $.50434 \mathrm{~F}+00-20000 \mathrm{E}+00$ | ． 50207 F ＋000． | ． $49487 \mathrm{~F}+001148$ | 4 | 2 | 306 |
| ． $20000 \mathrm{~F}+90$ | ． $49766 F+00.40000 \mathrm{C}+00$ | $.49546 \mathrm{~F}+00 .+0000 \mathrm{E}+00$ | ．49514E401148 | 4 | 2 | 307 |
| ． $80000 \mathrm{~F}+00$ | $.49077 \mathrm{E}+00-10000 \mathrm{E} \cdot 01$ | ． $48818 \mathrm{FF}+00$ | 1148 | 4 | 2 | 308 |
| 0.0 | 10．0．06 0 | 0 | 111148 | 4 | 2 | 309 |
| 11 | 2 |  | 1148 | 4 | 2 | 310 |
| －．10000E＋ 01 | ． $51727 \mathrm{E}+00-\mathrm{HOOOOE}+00$ | ．51201F－00－．600U0F－00 | ． $50794 \mathrm{~F}+00114 \mathrm{H}$ | 4 | 2 | 311 |
| －． $40000 \mathrm{~F}+00$ | ． $50461 \mathrm{~F}+00-20000 \mathrm{E}+00$ | ． $50168 F+000$ ． | ． $49908 \mathrm{~F}+001148$ | 4 | ？ | 312 |
| ． $20000 \mathrm{~F}+00$ | ． $49669 \mathrm{E}+00.40020 \mathrm{E}+00$ | ． $44435 \mathrm{~F}+00.60000 \mathrm{~F} \cdot 00$ | －49202E＋001148 | 4 | 2 | 313 |
| ． $800000+00$ | ． $49964 \mathrm{E}+00-10000 \mathrm{E}+01$ | ． 48716 F －00 | 114 A | 4 | 2 | 314 |
| 0.0 | 12．0＋06 0 | $0 \quad 1$ | 121149 | 4 | 2 | 315 |
| 12 | 2 |  | 1148 | 4 | $\stackrel{2}{2}$ | 316 |
| $-.10000 E+01$ | ． $52272 \mathrm{E}+00-.90000 \mathrm{E}+00$ | ． $51825 E+00-.80000 E+00$ | ． $51464 F+001148$ | 4 | 2 | 317 |
| －． $60000 \mathrm{E}+00$ | ． 50 R9BF＊ $00-.40000 \mathrm{E} * 00$ | $.50475 \mathrm{~F} \cdot 00-.20000 \mathrm{E}+00$ | ． $50129 \mathrm{C}+001148$ | 4 | 2 | 314 |
| 0. | ．49R2JE $+00.20000 \mathrm{E}+00$ | ． $49556 \mathrm{FE} 400.40000 \mathrm{~F} \cdot 00$ | ． $49313 F+001148$ | 4 | 2 | 319 |
| $.60000 \mathrm{~F}+00$ | $.49085 \mathrm{~F}+00 \cdot+00000 \mathrm{E}+00$ | $.4 R R G 6 E+00.100 U O E+01$ | $.48654 \mathrm{~F}+001148$ | 4 | 2 | 320 |
| 0.0 | $\begin{equation*} 14.0 * 06 \tag{0} \end{equation*}$ | $0$ $1$ | $121148$ | 4 | 2 | 121 |
| 12 | 2 |  | 1148 | 4 | $?$ | 322 |
| －．10000e＋ 01 | ．52R23F＊00－． $90000 \mathrm{E}+00$ | ． 521 BBF －00－．大0000F＊00 | ． $51048 t+001148$ | 4 | 2 | 323 |
| －． $60000 \mathrm{~F}+00$ | ． $509 R 2 \mathrm{CF}+00-.40000 \mathrm{~F}+00$ | ．50456F＊00－． $200 \cup 0 F+00$ | ． $5004 \mathrm{HE}+00114 \mathrm{H}$ | 4 | 2 | 324 |
| 0 ． | ．49703E $+00.20000 \mathrm{E} \cdot 00$ | ．49431F－00－40000F－00 | ． $44195 \mathrm{~F}+00114 \mathrm{H}$ | 4 | $\%$ | 324 |
| ．60000F－00 | ． 4900 SF． 00 ． $\mathrm{HOOOOE} * 00$ | ． 4 HR33E＋00 ． 10000 E 01 | ． $4 \mathrm{HOHHE}+0 \mathrm{O} 1148$ | 4 | $?$ | 326 |
| 0.0 | 16．0．0h 0 | $0 \quad 1$ | 1P1148 | 4 | $\checkmark$ | 3 C 7 |
| 12 | 2 |  | 114 H | 4 | 2 | 3er |
| －． 10000 Col | ．53433F＋00－． $40000 \mathrm{E} \cdot 00$ | ．52575F－00－． $40000 \mathrm{E}+00$ | ． $51424 \mathrm{~F} \cdot 00114 \mathrm{H}$ | 4 | 2 | 324 |
| －．A0000F－00 | ．51024F－00－．40000F＊00 | ． $90404 \mathrm{~F} \cdot 00-.20000 \mathrm{C} \cdot$ On | ．49439F－00114 | 4 | 2 | 3311 |


| 0. | ． 49587 F ． 00 | ．20nOOF－OO | ． 492 SHFP 00 | ． $40000 \mathrm{E} \cdot 00$ | －441）N1t＋00114H | $4 \quad 2$ | 331 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ．60000F－00 | ．4R930F－00 | ． $40000 \mathrm{~F}+00$ | ． 4 HRら．3F－00 | $.100005 \cdot 01$ | ． 4 AH $33 \mathrm{~F}+$ OO114 | 42 | $33 \%$ |
| 0.0 | 1R．0．06 | 0 | 0 | 1 | 12114 H | $4 ?$ | 133 |
| 17 | 2 |  |  |  | 114 H | 42 | 334 |
| －． $10000 \mathrm{~F} \cdot 01$ | ． $54092 \mathrm{~F} \cdot \mathrm{no-}$ | ． 40 OnOF 00 | ．529R2F－00－ | ．H0000F•00 | ． 52134 F ＋ 001148 | 42 | 3．34 |
| －．tnonotato | .51034 F － $00-$ | － 50000 － 00 | ．50327F－00－ | ． $20000 \mathrm{~F} \cdot 00$ | －49H02E＊OD114R | 4 ？ | 3.36 |
| 0. | ． 494 n6F－ 10 | ． $200008 \cdot 00$ | ．49126F＋00 | ． $40000 \mathrm{O}+00$ | － 4 HYA3F＊00114R | 42 | 3.37 |
| －minoorabor | ． 4 R90SF． 00 | ．HODOOE－00 | ． 4 K940F＋00 | ．10000F＋01 | ．44091E＊00114R | 2 | 3.34 |
| 0.0 | poroonm | $n$ | 0 | ， | 12114 ¢ | 42 | 339 |
| 17 | 7 |  |  |  | 114 H | 42 | $34 \pi$ |
| －． 10000 F 01 | ． 54907 E － $00-$ | ．40000F－00 | ． 5334 8F－00－ | ．R0000F•00 | ．52332F＊001148 | 42 | 341 |
| －． $60000 \mathrm{~F} \cdot 00$ | ．510P9F－00－ | ． $40000 \mathrm{O} \cdot 00$ | ． $50221 F+00-$ | －20000F 000 | ． $49635 \mathrm{E}+001148$ | 42 | 342 |
| 0. | ． 49 21 $\mathrm{HF}+00$ | ． $20000 \mathrm{E}+00$ | ．48958F－00 | ． $40000 \mathrm{~F} \cdot 00$ | ． $48840 \mathrm{E}+001148$ | 42 | 343 |
| －GOOODECON | ． 4 RAthfono | ． $\mathrm{HOOOOE}+00$ | ．49075F－00 | ． $10000 \mathrm{~F} \cdot 01$ | ．4446大E－00114B | 42 | 344 |
|  |  |  |  |  | 114 A | 40 | 345 |
|  |  |  |  |  | 114 K | 0 | 346 |
| 1.001 E．03 | $9.99175-01$. | 0 | 0 | 0 | 01148 | 74 | 347 |
| 0.0 | 0.0 | 0 | 0 | 12 | 11148 | 74 | 344 |
| 0.0 | 2．0 E．0． 0 | 9．9917E－01 | 5.0 | 0.0 | $0.0 \quad 1148$ | 74 | 344 |
| 1.0 | 2．0444F＋01 | Ч．4917E－01 | 0.0 | 0.0 | 0.01148 | 74 | 350 |
|  |  |  |  |  | 114 H | 70 | 35： |
|  |  |  |  |  | 1148 | 00 | 35？ |
| 1．0010F．OT | 9．0017F－n1 | 1 | 0 | 17 | 011481 | 12102 | 353 |
| 0. | n． | 0 | 0 | 1 | 211491 | 17102 | 354 |
| $p$ | － |  |  |  | 1144 | 12102 | 345 |
| 1．0000F－04 | 1．0n00F－ 00 | 2．0000t＋07 | 1．0000t 00 |  | 114 H | 12102 | 3 bh |
| 1．1735F－07 | 0 ． | 0 | $?$ | 1 | 3114 H | 12102 | 347 |
| 3 | $?$ |  |  |  | 11481 | 12102 | 354 |
| 1．9000F＋07 | 0 | 1．4001F＊O7 | 1．00not 00 | $2.0000 F+07$ | 1．0000F60011481 | 12102 | 349 |
| 1．0725F＋07 | 0 。 | 0 | ？ | 1 | 41148 | 12102 | 3 HO |
| 4 | $?$ |  |  |  | 1148 | 12102 | 361 |
|  | 0 ． | 1．6001F－07 | 1.0000 P 00 | 1．ROOOF 007 | 1．0000F－0011481 | 12102 | 3n？ |
| 1．AON1F．O7 | 0 。 |  |  |  | 1148 | 12102 | 3 n 3 |
| O．TOCOF．OA | $\bigcirc$ 。 | 0 | $?$ | 1 | 4114 H | 12102 | 3 H 4 |
| 4 | $\dot{\gamma}$ |  |  |  | 114 H | 1－10？ | 3 S 5 |
| 1．4000F＋07 | 11. | 1．40U1F．07 | 1.0000 P 00 | 1．nounfot 07 | 1．0000F－00114． | 12102 | int． |
| 1．4001F＋07 | 0 。 |  |  |  | 1144 | 12102 | ．in！ |
| H．TPSOFPOK | 0. | 11 | $?$ | 1 | 4114 A | 12102 | Itar |
| 4 | $?$ |  |  |  | 114 H | 121．12 | ．7nc |
| 1．3000f 07 | 0. | 1． $2001 F+07$ | $1.0000 \mathrm{t}+00$ | $1.4000 F+07$ | 1．0000\％+00114 A | 1210 ？ | 370 |
| 1．4001F．07 | $n$ |  |  |  | 114 H | 12108 | $3 / 1$ |
| 7．7ア50F．fif | 0. | 0 | 2 | 1 | 4114 H | 12102 | 312 |
| $4$ | $?$ |  |  |  | 1148 | 1210\％ | 173 |
| 1．0nロnf．07 | ＇n． | 1．0001F．07 | $1.0000 \mathrm{~F} \cdot 00$ | 1．2000F．07 | 1．0000f 001148 | 12112 | 314 |
| 1．2001F－07 | 0. |  |  |  | 1148 | 1210？ | 375 |
| 6.97505 .06 | 0. | 0 | 2 | 1 | 41148 | 12102 | 376 |
| 4 | $\gamma$ |  |  |  | 1144 | 1310？ | 377 |
| 4.00006 .06 | 0. | 4．0001F＋06 | 1.0000 F 00 | $1.0000 \mathrm{~F}+01$ | 1．0000f 0011144 | $1210{ }^{2}$ | ． 178 |
| $1.00015 \cdot 07$ | 0. |  |  |  | 1148 | 12102 | 174 |
| H．4750F－06 | 0. | 0 | $?$ | 1 | 4114 A | 1210？ | 3140 |
| 4 | 2 |  |  |  | 1148 | 1く102 | 3H1 |
| B．DOOOEPGK | $\dagger$ 。 | H．0001t－On | $1.0000 \mathrm{t}+00$ | Y．0000F－06 | 1．0000F－001144 | 12102 | 3 H 2 |
| $4.0001 \mathrm{~F} \cdot 0 \mathrm{nc}$ | 0 ． |  |  |  | 1144 | 12102 | 3H． 3 |
| 5．8750F．06 | 0. | 0 | 2 | 1 | 4114 A | 12102 | $3 \mathrm{H}_{4}$ |
| 4 | 2 |  |  |  | 1148 | 1210？ | 345 |
| $7.00005+06$ | 0 ． | 7．0002F－0ts | $1.0000 \mathrm{E} \cdot 00$ | H．0000F＋06 | 1．0000t．00114 | 1217？ | 3n6 |
| 9．0n0lFton | 0. |  |  |  | 1148 | 12102 | 3n7 |
| $5.4750 \mathrm{E} \cdot 06$ | $\Omega$. | 0 | 2 | 1 | 41148 | 1210？ | 3 HH |
| 4 | $?$ |  |  |  | 1149 | 12102 | $3 \mathrm{H4}$ |
| F．0090F－06 | 0. | 6．0001EPOt | $1.0000 \mathrm{~F}-00$ | 7．0000F +06 | 1．0000F +00114 H | 12102 | 340 |
| 7．0001F．06 | 0 |  |  |  | 1148 | ご102 | 341 |
| $4.77505+06$ | 0. | 0 | $?$ | 1 | 41148 | 12102 | $3 \cup 2$ |
| 4 | 2 |  |  |  | 114 A | 12107 | 343 |
| $5.0 n 00 F+n 6$ | 0 。 | $5.00015 * 06$ | $1.0000 \mathrm{~F}+00$ | $6.0000 F+06$ | 1．000nF＋ 00114 R | 12102 | 394 |
| 6．0001F－06 | 0 ． |  |  |  | － 1148 | 12102 | 395 |
| $4.47505+176$ | 0 ． | 0 | 2 | 1 | 41148 | 12102 | 346 |
| 4 | $\cdots$ |  |  |  | 1148 | 12102 | 307 |
| $4.0000 E+06$ | 0. | 4．0001E＋06 | $1.0000 t+00$ | $5.0000 \mathrm{~F}+06$ | 1．0000F．00114R | 12102 | $34 \%$ |
| S．0001E＋06 | n． |  |  |  | 1148 | 12102 | 399 |
| $3.9750 F+06$ | 0. | 0 | 2 | 1 | 41148 | 12102 | 400 |


| 4 | 2 |  |  |  | 114412102 | 417 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 3．0000F－06 | 0. | 3．0001F．06 | $1.0000 \mathrm{t} \cdot 00$ | $4 . V O C O F \cdot O 5$ | 1．0000F－00114A12102 | 492 |
| $4.0 n 01 F+06$ | 0 。 |  |  |  | 114812102 | 403 |
| $3.4750 F+06$ | 0 ． | 0 | 2 | 1 | 4114812102 | 4114 |
| 4 | 2 |  |  |  | 114812102 | 404 |
| P．0000F－06 | 0 。 | ？．n00lfat | $1.00001-00$ | 3．0000r．0h | 1．0000t－00114H1710？ | 406 |
| 3．0001F＋06 | 0. |  |  |  | 114412102 | 4118 |
| P． 475 0F－06 | 0. | 0 | 2 | 1 | 4114817107 | 4010 |
| 4 | $\gamma$ |  |  |  | $114 \mathrm{H1P102}$ | 404 |
| 1．000nF＊06 | ）． | 1．1001t．un | 1．0000t•00 | 2．UOUOF•00 | 1．0000t－00114912102 | 410 |
| 2．000）5．06 | 0. |  |  |  | 114817102 | 411 |
| 2．APCOF－06 | 0 ． | 0 | ？ | 1 | 4114813102 | $41 \%$ |
| 4 | $?$ |  |  |  | 114412102 | 413 |
| 0．0000F－05 | 0. | 6．0001F．05 | 1．0000F－00 | 1．0000f•06 | 1．0000F－00114417102 | 414 |
| 1．0001F．06 | 0. |  |  |  | 114812102 | 415 |
| P．4 $3.50 \mathrm{~F}+16$ | 0 ． | 0 | 2 | 1 | 4114812102 | 416 |
| 4 | 3 |  |  |  | 114412102 | 417 |
| 2．0000F．05 | 0. | ？．0001E＋05 | 1．0000E 00 | 0．0U00F．0S | 1．0000E＋00114813102 | 414 |
| 6．0001F．05 | 0 。 |  |  |  | 114 Bl 2102 | 414 |
| 2．2250F．06 | 0. | 0 | 2 | 1 | 1114412102 | 420 |
| 3 | 2 |  |  |  | $11481<102$ | 421 |
| 1．0000F－05 | 1．0n00F•00 | $2.0000 F+05$ | 1．0000F－00 | $2.0001+05$ |  | $42 ?$ |
|  |  |  |  |  |  | $4{ }_{4}^{4} 3$ |
|  |  |  |  |  |  | $4{ }_{4}^{4}$ |
| 1.0010 F 03 | $4.9917 \mathrm{~F}-\mathrm{nl}$ | 1 | 0 | 0 | $0114 \times 14102$ | $4{ }^{4} 5$ |
|  |  |  |  |  | 1144140 | 4 Co |
|  |  |  |  |  | 114403 | 427 |
|  |  |  |  |  | －1 | 4.24 |

Appendix B

## Summary for ${ }^{3}$ He

* 

SUMMARY DOCUMENTATION FOR He-3
L. Stewart and R. J. LaBauve

Los Alamos Scientific Laboratory


#### Abstract

ABSTRACI

The following nuclear data are given for ${ }^{3} \mathrm{He}$ in the energy range from $1.0 \times 10^{-5} \mathrm{eV}$ to 20 MeV :

File 1. The general description of the data which follow File 2. Values of the nuclear spin and effective scattering radius

File 3. Smooth point-wise data for the total, the free-atom elastic, $(n, p),(n, d)$, and the rediative capture cross sections; data for $\bar{\mu}, \xi$, and $\gamma$ are also included

File 4. The angular distributions for elastic scattering as probability vs cosine of the scattering angle in the center-of-mass system

File 7: The free-atom-scattering cross section at thermal

\section*{I. INTRODUCTION}


These data were translated from an unpublished evaluation completed by L. Stewart in 1968. In 1971, the Standards Subcommittee of CSEWG reviewed the file and concluded that the ( $n, p$ ) cross section was still adequately represented to be recommended as a standard cross section.

## II. TOTAL CROSS SECTION

The total cross section was obtained by suming the partials up to 100 keV . From 100 keV to 20 MeV , the LASL measurements ${ }^{1}$ were used exclusively in this eveluation.

## III. ELASTIC SCATTERING

Available measurements of Seagrave, Cranberg, and Simmons ${ }^{2}$, of Sayres, Jones, and Wu ${ }^{3}$, and of Antolkovic et al. 4 were used. Also the $p+T$ scattering was used to fill the gaps in energy where no $n+3$ He elastic scattering measurements exist, The $p+T$ experiments employed were those of Brolley et al.5, Rosen and Leland', and of Vanetsian and Fedchenko ${ }^{7}$. Wick's Limit was employed at all energies to insure the nonviolation of unitarity. The angular distributions are given as probabilities versus cosine of the center-of-mass scattering angle.

[^0]
## IV. RADIATIVE CAPIURE

Gallmann, Kane, and Pixley ${ }^{8}$ have placed upper limits on the thermal capture cross section of $100 \mu \mathrm{~b}$ and $10 \mu \mathrm{~b}$ for gama and pair emission, respectively. Since these are upper limits and absolute measurements do not exist, no estimate is made here for radiative capture. The gamma-ray production cross sections are also assumed to be negligible and are therefore ignored.

## V. ( $n, p$ ) CROSS SECTION

The ( $n, p$ ) cross section below 10 eV was derived solely from the measurements of Als-Nielsen and Dietrich ${ }^{9}$ giving $5327+10$ b at thermal. A $1 / v$ extrapolation was assumed to 1.7 keV , where the slope was changed to merge with the slope of the curve given by the data of Gibbons and Macklin ${ }^{10,11}$. Many experiments ${ }^{3}, 12-17$ have been performed at higher energies although, all too often, the cross sections were not obtained on an absolyte basis. The most extensive absolute measurements were those of Perry et al. 13 which have been heavily weighted in this evaluation.

## VI. ( $n, d$ ) CROSS SECTION

Only the Columbia data ${ }^{3}$ near 7.5 MeV were available on this reaction. Bradbury and Stewart18, however, employed detailed balance and the LASL measurements on the inverse reaction, that is, the $D(d, n)^{3}$ He reaction, to predict the energy dependent cross section to 15 MeV . Above 15 MeV , these data were extrapolated.

## VII. THREE- AND FOUR-BODY BREAKUP

The $3_{\mathrm{He}}(\mathrm{n}, \mathrm{np}) \mathrm{D}$ and $3_{\mathrm{He}}(\mathrm{n}, 2 \mathrm{n} 2 \mathrm{p})$ reactions have $Q$ values of -5.494 and -7.718 MeV , respectively. Only a few measurements exist on these reactions, and these are usually limited to a search for final-state phenomena. The spectrum is usually observed at one angle very close to zero degrees. Observation of the proton spectrum at 24.4 MeV reveals no clear indication of $n-d$, twonucleon, or three-nucleon final-state interaction. A strong n-p final state is evident from measurements of the deuterron spectrum at $\theta_{d}=5^{\circ}$. An upper limit of 12 mb has also been set on the $3 \mathrm{He}(\mathrm{n}, 2 \mathrm{n} 2 \mathrm{p})$ reaction. In the absence of measurements of the absolute cross sections, these break-up cross sections have been assumed small and are therefore ignored in the present evaluation.

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Appendix $C$

## Summary for <br> ${ }^{6} \mathrm{Li}$

# SUMMARY OF ${ }^{6}$ Li DATA FOR ENDF/B-III ${ }^{*}$ <br> by 

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## INTRODUCTION

The data, in the energy range of $10^{-5} \mathrm{eV}$ to 20 MeV , for $\mathrm{Cli}_{\mathrm{Ii}}$ (MAT 1115) were submitted to the NNCSC in September 1971 for inclusion in ENDF/B-III. These data represent an extensive revision of the earlier file (MAT l005) which was, for the most part, based on a UKAFA evaluation. ${ }^{1}$ A major change was made in that below about 1.7 MeV the cross sections for MATlllf reflect very strongly the detailed review of the available data by Uttley, Sowerby, Patrick and Rae. In choosing the latter data, consideration was also given to the recommendation of the CSEWG Normalization and Standards Subcommittee that Uttley's ( $n, \alpha$ ) data be used in the latter energy range. ${ }^{3}$ The data above 1.7 MeV for MAT 1115 will be discussed in later sections of this document. Forthisfirst pass reevaluation, the ( $n, \gamma$ ) cross sections from the earlier UK evaluation were retained. Following uttley et al., the data will be considered in the following energy intervals: (1) thermal, (2) thermal to 10 keV , (3) 10 to 500 keV , and (4) 500 keV to 1.7 MeV . Above 1.7 MeV , the data will be considered by reaction type.

## THERMAL ENERGIES

The cross sections given in the file at 0.0253 eV are as follows:

$$
\begin{aligned}
\text { Total } & =941.015 \mathrm{~b} \\
\text { Elastic } & =0.72 \\
(n, \alpha) & =940.25 \\
(n, y) & =0.045
\end{aligned}
$$

It is estimated that the ( $n, \alpha$ ) cross section is known to $\pm 0.5 \%$. The choice of Uttley et al. for the ( $n, \gamma)$ cross section is $30 \pm 8 \mathrm{mb}$.

THERMAL ENERGIES TO 10 keV
The ( $n, \alpha$ ) cross sections up to 100 eV were calculated from the formula

$$
\sigma(n, \alpha)=(149.56 / \sqrt{E})-0.024 b
$$

Above 100 eV , the p-wave absorption contribution from the $247-\mathrm{keV}$ resonance becomes increasingly important, and at 10 keV the negative s-wave absorption ( -0.024 b ) is largely cancelled. Uttley et al. assign an uncertainty of $\pm 1 \%$ to the ( $n, \alpha$ ) cross section from thermal to 10 keV . Over this energy range, the data given in the file deviate from a strict $1 / v$ dependence (149.56/ $\sqrt{E}$ ) by a meximum of $-0.4 \%$.

[^1]The elastic cross section is held constant at 0.72 b up to 2 keV with a monotonic increase thereafter to 0.7221 b at 10 keV . The elastic angular distribution $u p$ to 10 keV is specifised as isotropic in the $C M$ system.

10 to ${ }^{\circ} 00 \mathrm{kcV}$
The uncertainties estimated for the $(n, \alpha)$ cross section are:
$\pm 2 \%$ at $100 \mathrm{keV}, \pm 5 \%$ from 100 to 300 keV , and $\pm 10 \%$ at 500 keV .

As can be seen from Fig. 2 of Ref. 2, there is great disagreement among the ( $n, \alpha$ ) measurements in this energy range. The ( $n, \alpha$ ) curve used in this evaluation is the one recommended by uttley et al. based on their careful review of the experimental data. In arriving at the recomended curve, the latter authors have considered the following points:

1. Errors in the energy scales of the measurements.
2. The inadequacy of those experiments in which the cross sections were measured relative to the 235 U fission cross section.
3. The inadequacy of those measurements in which the neutron flux was measured relative to a long counter.
4. The inconsistency of measurements made using thick lithium detectors with total and scattering cross section measurements.

Uttley et al. do, however, state that much experimental work ${ }^{2}$ remains to be done to confirm their recommendations.

In general, the total cross sections are those reported by Uttley and Diment. 4 The scattering cross sections below and above 100 keV reflect the data of Asami and Moxon ${ }^{5}$ and Lane, Langsdorf, Monahan, and Elwyn ${ }^{6}$, respectively. For the elastic angular distributions, Legendre coefficients (from which the normalized probability distributions are reconstructed) were inferred by fitting the data of Lane et al. 6

500 keV to 1.7 MeV
The ( $n, \alpha$ ) data in this energy range were obtained by subtracting the scattering cross section from the total cross section. The total cross section values were essentially those reported by Dinent and Uttley. ${ }^{4}$ For the scattering cross sections the measurements of Lane et al. ${ }^{6}$ and Knitter and Coppola7 were considered. The ( $n, \alpha$ ) cross section uncertainties are estimated to be:
$\pm 10 \%$ at 500 kev , increasing to $\pm 15 \%$ between 700 and 1000 keV , and decreasing to $\pm 10 \%$ by 1.7 MeV .

## TOTAL CROSS SECTION

From 2 to 15 MeV , the primary reference for the data in the file is the work of Foster and Glasgow. ${ }^{\circ}$ The extrapolation to 20 MeV was based on the measurements of Peterson, Bratenahl, and Stoering. 9

## ELASTIC CROSS SECTION

Data given between 4 and 10 MeV are heavily veighted towards the Hopkins, Drake and Condé evaluation. 10 A value of 0.88 b at 14 MeV was used and data smoothly extrapolated to 20 MeV .

Iegendre coeffictents for the angular distributionsup to 2.5 MeV werc determined from the data of Lane et al. 6 Retween 4.83 and 7.5 MeV coefficients were inferred from Hopkins et al. data. 10 Based on 14 MeV elastic scattering data given in BNL-400, optical model calculations (ABACUS code) were performed to infer Legendre coefficients between 10 and 20 MeV . Data for Mt $251\left(\bar{\mu}_{\sim}\right)$, $252(\xi)$, and $253(\gamma)$ were caluclated using the elastic angular distributions given in File 4.

## THE ( $n, 2 n$ ) on CROSS SECTION

The cross sections and angular distributions for this reaction (MT = 24) are the same as in Ref. 1 up to 15 MeV with a smooth extrapolation to 20 MeV . The secondary energy distributions given in Ref. l have been approximated by ENDF/B Law 9 with $\theta=0.21 \sqrt{E}$ (MeV).

## THE $\left(n, n^{\prime}\right) \gamma$ and $\left(n, n^{\prime}\right)$ odd CROSS SECTIONS

The ( $n, n^{\prime}$ ) $\gamma$ cross sections tabulated under $N T=52$ are those measured by Presser, Bass, and Krüger ${ }^{11}$ up to 7 MeV , with a constant 5 mb assumed thereafter. Isotropy in the $C M$ system is specified for the angular distribution.

The ( $n, n^{\prime}$ ) ad data given under $M T=91$ are the same as in Ref. 1 up to 3 MeV , with the data of Hopkins et al. 10 taken into account between 4 and 10 MeV . The value of 433 mb assigned at 14 MeV is higher than the 403 mb given in Ref. 1 , which is higher than the nominal 330 mb reported experimentally. A more detailed evaluation of the $14-\mathrm{MeV}$ cross sections will be required to resolve the latter discrepancy. The extrapolation to 20 MeV of the ( $n, n^{\prime}$ ) od cross section was obtained by subtracting the sum of all other partials from the total cross section. For the angular distributions, the tabulated values of Ref. l, extrapolated to 20 MeV , were used. The secondary energy distributions of Ref. 1 were approximated using ENDF/B Law 9 with $\theta$ values obtained by linear interpolation between the following points:

$$
\begin{array}{ll}
E=1.718 \mathrm{MeV}, & \theta=0.05 \mathrm{MeV} \\
E=4.1 & \theta=0.75 \\
E=20.0 & \theta=8.40
\end{array}
$$

THE ( $n, p$ ) AND ( $n, \alpha$ ) CROSS SECTIONS
The ( $n, p$ ) cross sectionsup to 7 MeV reflect the data of Presser et al. 11 Above 7 MeV the data of Ref. I were used and extrapolated to 20 MeV . For the ( $n, \alpha$ ) cross sections the data between 2 and 15 MeV are those of Ref. 1. Extrapolation to 20 MeV was based on the measurements of Kern and Kreger ${ }^{12}$ between 15 and 18 MeV.

## ACKNONLEDGMENTS

We are indebted to Leona Stewart for providing us with the data below 2 MeV and plots of the angular distributions of Lane et al., and to P. G. Young for performings the optical model calculations.

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## Appendix D

Summary for ${ }^{10}{ }_{B}$

SUMMARY DOCUMERTIATION FOR B-10*

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## ABSTRACT

The following neutron data are given for B-l0 in the energy range $1.0 \times 10^{-5} \mathrm{eV}$ to 15 MeV ;

File 1. General description of data which follow.
File 2. Values of nuclear spin and effective scattering radius only.
File 3. Smooth cross sections for total, free-atom elastic, total inelastic, the inelastic level at 717 keV ( $M P=51$ ), the inelastic continuum, $(n, d),(n, t)$, and $(n, \alpha)$. Also data for $\bar{\mu}, \overline{5}$, and $\gamma$ are included.

File 4. Angular distributions for elastic scattering (expressed as Legendre polynomial coefficients in the center-of-mass system), the first inelastic level, and the inelastic continuum. (The inelastic distributions are given as tabulated functions with the assumption of isotropy in center-of-mass system.)

File \%. Secondary energy distribution for the inelastic continuum (law 9 - evaporation spectrum).

## INTRODUCTION

This evaluation for ${ }^{10}$ B (MAT 1155) is essentially the same as that given for $\mathrm{FNDF} / \mathrm{B}$, Version I in MAT 1009 which was performed in October 1967 by D. C. Irving of ORNL. At the CSEWG meeting of May 19-20, 1971, the Standards Subcommittee requested LASL to modify the ( $n, \alpha$ ) cross sections (MP = 107) of ${ }^{10} \mathrm{O}_{\mathrm{B}}$ (MAT 1009) below 100 keV to conform to a recent evaluation of Sowerby et al. ${ }^{2}$ In carrying out this request, it was also decided to 1 ) extend the range of modification to 150 keV , and 2) change the elastic scattering cross sections ( $M \mathrm{M}=2$ ) in the modified energy range, and, of course, 3) change the total cross section to be equal to (MI2 + MT107).

The MAT 1009 evaluation contained 15 datum points from $2.0 \times 10^{-5} \mathrm{eV}$ to 150 keV . As the Sowerby evaluation is based on experimental measurements of the ratio $\sigma_{\mathrm{n}, \alpha}\left({ }^{( }{ }_{\mathrm{LI}}\right) / \sigma_{\mathrm{n}, \alpha}\left({ }^{\left.10_{\mathrm{B}}\right) \text {, the modified cross sections were put on the same mesh }}\right.$ ( 30 points) currentiy in use in the LASL ${ }^{G_{\mathrm{LI}}}$ (MAT 1115) evaluation. The energy range was also extended to 150 keV as the MAT 1009 data were easier to merge at this point than at 100 keV . Moreover, it was felt that additional points were needed between 100 keV and 150 keV to compare with the ${ }^{6} \mathrm{Li}$ data.

The ( $n, \alpha$ ) modified cross sections, as recommended by Sowerby at al., are given by the formula:

[^2]$\sigma_{n, \alpha}\left(10_{B}\right)=\frac{13.736}{\sqrt{E}}-0.312-1.014 \times 10^{-2} \sqrt{E}+\frac{2.809 \times 10^{5}}{\sqrt{E}\left[(170.3-E)^{2}+2.243 \times 10^{4}\right]}$
with $\sigma$ in barns and F in keV .
Tncidentaliy, the $\sigma_{n,} \alpha$ datn in the old momp/ ${ }^{10} 0_{1}$ evaluation ald not differ greatly from this formula. The largest deviation from formula (l) for MAT 1009 was about $4 \%$. The most significant change was the use of a more descriptive mesh.

This investigation did reveal a need to change the elastic scattering cross section, however. The ENDF/B evaluation was based on 1966 data of Mooring; ${ }^{3}$ whereas this modification was based on more recent (1969) data of Asami and Moxon. 4

FILE 1: GENERAL INFORMATION
for $10_{B}$ A brief summary of these data is given in File 10.0130 . The atomic mass

FILE 2: RESONANCE INFORMATION
A nuclear spin of $s$ and effective scattering radius of $0.399 \times 10^{-1: 2} \mathrm{~cm}$ 1.s given in file 2.

FIIE 3: SMOOMT CROSS SFCTITONS
Below 150 keV :
$\sigma_{\text {elastic }}-\begin{aligned} & \text { from experimental data of Asami and Moxon } \\ & \\ & \text { constant value of } \\ & 2.2 \text { barns below } 100 \mathrm{eV} \text {. with a }\end{aligned}$
$\sigma_{n, \alpha}-$ given by formula (1) (Sowerby et al. ${ }^{2}$ ).
$\sigma_{\text {total }}-\left(\sigma_{\text {elastic }}+\sigma_{n, \alpha}\right)$.
From 150 to 500 keV :

$$
\left.\begin{array}{rl}
\sigma_{\text {elastic }}-\begin{array}{l}
\text { from the smooth curve of Mooring. } \\
\\
\text { With the data of Lane et al. }
\end{array} \\
\sigma_{n, \alpha}-\text { from a smooth curve through the }(n, \alpha) \text { data of Mooring }{ }^{3} \\
& \text { and Gibbons, } 5 \text { with little weight placed on the data } \\
\text { of Cox. } 6
\end{array}\right\}
$$

Above 500 keV :
$\sigma_{\text {elastic }}-\left(\sigma_{\text {total }}-\sigma_{\text {inelastic }}-\sigma_{n, \alpha}-\sigma_{n, d}-\sigma_{(n, t){ }_{2 \alpha}}\right)$.
$\sigma_{\text {total }}$ - from the data displayed in BNL325. 7,8 Some minor adjustments were made to remove spurious wiggles in the elastic cross section above 4 MeV .
$\sigma_{n, \alpha}$ - from the data of Davis 9 increased by 110 mb as suggested by a comparison to the data of Gibbons and Nellis. 10
$\sigma_{(n, t)} \mathbb{Z}_{\alpha}$ - from a smooth curve was drawn by eye tbrough the sparse data of Frye, 11 Wyman, 12 and Perkin. 13
$\sigma_{n, d}-a$ smooth curve was drawn through the $\operatorname{Be}(\alpha, n)$ data of Bardes 14 and Siemsson ${ }^{15}$ and detailed balance used to obtain ${ }^{10}{ }_{B}(n, d)$ values. These were connected by a straight line to a value at $1+4 \mathrm{MeV}$ obtained by integration of the data of Valkovic ${ }^{16}$ which is slightly higher than that of Ribe. 17
$\sigma_{\text {inelastic }}$ - data from Day ${ }^{18}$ on excitation of the first level which agrees with that of Nellis was used to 4.5 MeV . This was connected by a smooth curve to a value at 14 MeV obtained by subtracting ( $n, d$ ), ( $n, \alpha$ ), and ( $n, t$ ) from the nonelastic value of MacGregor. 19

The ( $n, n^{\prime} d$ ) $2 \alpha$ reaction was neglected since it is contained for all practical purposes in the inelastic data.

## FIIE 4: SECONDARY ANGULAR DISTRIBUTIONS

The experimental data for elastic scattering angular distribution is sparse, consisting of a few points between . 5 and 2 MeV (BNL400)20 and one at 14 MeV (reported in Valkovićl6). By all rights an optical model should not be valid for boron. Hovever, the calculations of Agee 21 compare beautifully with the experimental data and have been used in the evaluation. The angular distrim bution was taken to be isotropic below .5 MeV in agreemen' with BNL 400 and the data of Lane et al. (see Mooring).

The secondary angular distribution for inelastic scattering was assumed isotropic in the center-of-mass system.

## FIIE 5: SECONDARY ENERGY DISTRIBUTIONS

Up to 4.5 MeV , inelastic scattering was assumed to proceed via the first level at . 71 MeV . Above 4.5 MeV an evaporation spectrum was used with the temperature as determined by Weinberg and Wigner. 22

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## Appendix E

## Summary for <br> ${ }^{235} \mathrm{U}$

# ENDF/B-III SUMMARY DOCUMENTATION FOR <br> $J^{235}$ 

## Contributors:

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(1) ANCR 1044
(2) AI-AEC-12916; AI-AEC-13013
(3) BNWL-1586
(4) WARD-4210T4-1

## Thermal Data

The cross section shapes in this evaluation were derived, in general, by perturbing the results of an existing evaluation in order to reproduce the desired y values. The results of existing precision experimental differential measurements were used as a guide to the nature of the perturbations.

An objective of the evaluation was to provide cross-section shapes which were smooth in the thermal region and would not produce irregularities in the behavior of thermally averaged quantities as a function, say, of neutron temperature.

Point representation of the data intended to be in the ENDF/B-I files were received from C. R. Lubitz. These files covered the energy range $10^{-3}$ to 5 eV . It was determined that these files contained apparently unintended irregularities as large as one percent in the entries below 4 meV . These entries were smoothed and the data files extended to $10^{-4} \mathrm{eV}$. Calculation of " g " factors showed that the value of $g_{f}$ was about 0.2 percent greater than the desired value. The original evaluation of the low-energy fission cross section was reported to have been made by fiting $\pi_{f}$ data of LRL and Hanford. Accordingly, a new fit was made including these data and also the fission dala of Safford and Melkonian. A smooth fit was obtained, the main difference from the original fit being the reduction of the rise in cross section at energies below about 0.02 eV . The fission cross section was fitted simultaneously with the capture cross section using an alpha variation deduced by Westcott. In order to achieve a smooth fit to the fission data the file was modified slightly for energies up to 0.18 eV . The capture cross section file was modified for energies up to 0.1 eV .

The cross-section shapes derived for sub-thermal neutron energies are based in large part on precision total cross section measurements of Safford, et al. and the $1+\alpha$ measurement of Safford and Melkonian. The total cross section data of Safford, et al. obtained with liquid samples is shown below compared with the ENDF/B values.

|  | $\sigma_{\mathrm{T}} \sqrt{\mathrm{E}}$ |  |
| :--- | :--- | :--- |
| Neutron Energy | Safford, et al. | ENDF/B |
| 0.000818 eV | $115.71 \pm 0.35$ | 115.56 |
| 0.00128 eV | $115.79 \pm 0.29$ | 115.49 |

The ENDF/B scattering cross sections at these energies is 16 b and $\mathrm{a} \pm 3 \mathrm{~b}$ uncertainty in this cross section contributes about $\pm 0.1 \mathrm{~b}-\mathrm{eV}^{1 / 2}$ to calculated $\sigma_{T} \sqrt{E}$ values.

A summary of the $2200 \mathrm{~m} / \mathrm{s}$ coross section parametors is piven bolow.

| $\sigma_{\mathrm{T}}$ | 694.276 |
| :--- | :--- |
| $\sigma_{\mathrm{S}}$ | 15.776 |
| $\sigma_{\mathrm{f}}$ | 580.2 |
| $\sigma_{\mathrm{C}}$ | 98.3 |
| $\nu$ | 2.423 (Delayed + Prompt) |
| $a$ | 0.1694 |
| $\eta$ | 2.07196 |

## Resolved Resonance Region

The resolved resonance region for ${ }^{235} \mathrm{U}$ in ENDF/B-III extends from 1 to 82 eV . The description uses single level parameters plus a smooth file. Parameters were derived from a simultaneous fit to the following sets of data:

1) Simullameous caplure and fission mea surement by de:Saussure, et at. The strength of this experiment is that it measured the two most important partial cross sections of $\mathrm{U}^{235}$ simultancously, under the same conditions of resolution and background. Morcover, care was taken to correct for such effects as backgrounds, resonance self-shielding, and scattering in the fission chamber. These data were used principally to indicate the ratio of capture to fission for the resonances.
2) Total cross section measurements of Michaudon. These data were obtained at liquid nitrogen temperatures and fairly high resolution. They turned out to give in most cases the best indication of the total widths of the resonances. The data were available only as cross section vs. energy, with results from several samples mixed together. Total cross sections are measured from transmission of samples, and the analysis should really be performed on the transmission data for each sample.
3) Fission cross sections measured by Blons, et al. on the Saclay lincar accelerator. These data were obtained at liquid nitrogen temperature, with resolntion similar to that of Michaudon's total cross section measurement. The Blons data are the best resolution fission data, but below about 35 cV the normalization gets progressivcly more erratic because of difficulty in interpreting the backgrounds in the presence of a B-10 filter used to eliminate low energy overlap neutrons.
4) Fission cross sections measured by Cao, et al. on the linear accelerator at C.B.N.M. (Geel). These data are the highest resolution room temperature measurements of $\sigma_{f}$ for $U^{235}$. They are useful for comparing with the Blons data to confirm the effectiveness of the Doppler corrections in the analysis code. They go to a lower energy than the Blons data, 6 eV vs. 17 eV . However, the Cao data are troubled by crratic background corrections in the vicinity of resonances in filters used to determine backgrounds.

## Cross Section Normalization

Since this analysis covered only the resonance region above 1.0 eV , it was necessary to normalize all data to the existing ENDF/B-II low energy file. Of the principal data sets, only the deSaussure measurements extends to this low energy. His fission data were raised by $1.5 \%$ to bring their integral from 0.45 eV to 1.0 eV into better agreement with that from the ENDF/B low energy file. The difference in the capture integrals was $2.4 \%$. Nevertheless, the capture was raised only $1.5 \%$ in order that deSaussure's a ratios might be preserved.

The Cao data were raised $7 \%$ to bring them into agreement with the renormalized deSaussure data. The Blons data, which already agreed well with the renormalized deSaussure data above 40 eV , were given an energy-dependent renormalization. The ratios to deSaussure values of a series of incremental resonance integrals were fit with a fourth degree polynomial. This polynomial was then used to normalize the Blons data. The resulting correction ranged from about $19 \%$ at 18 eV to zero at 40 eV .

Single level parameters were derived by fitting the experimental data by means of the automatic iterative fitting features of the Automated Cross Section Analysis Program (ACSAP). A value of 11.5 b was used for the potential scattering cross section.

ACSAP will upon request print and plot the differences between experimental points and the cross sections calculated from parameters. Such difference outputs were used in constructing the smooth files. In ordex to maintain proper a values, the deSinssure data were used as much as possible in ronstructing docse different files. However, this ideal had to be abandoned above about 35 eV , as degraded resolution spread the intrinsic difference well away from the resonances to which they apply.

The scattering smooth file represents the difference betwecn the singlelevel prediction and a multilevel calculation adjusted to minimize the effects of interference imbalance at the two ends of the resolved resonance region.

The parameters plus smooth file yield resonance integrals between 1 and 82 eV of 170.6 and 104.5 b for fission and capture, respectively. The average alpha is 0.613 in this region, compared to a value of 0.617 calculated directly from deSaussure's data.

The root mean square fractional diference between the fit and the individual data scts averages about $3.5 \%$. This figure comes from an analysis of pirlial resonance integrals from the fit and from the data. Resonance encrgies are probably good to 0.050 eV and resonance widths to $10 \%$. Overall error in cross sections is about $5 \%$.

For further details see the full report, ANCR-1044.

## Unresolved Resonance Region

The ENDF/B-III evaluation in the unresolved energy range is based primarily on the experimental data of deSaussure. At the time of this evaluation, detailed resolution cross sections were not available from the recent measurcments of Perez, Blons or Lemley. For this evaluation, a continuous curve of the fission cross section was constructed so as to reproduce the decimal interval averages of the experimental data.

Differences between the present and ENDF/B-II evaluations are due to inclusion of new experimental data, renormalization of the experimental data to recent ${ }^{10} B(n, \alpha)$ cross sections and to differences in methods of constructing a smooth curve. The ENDF/B-III evaluation was obtained by averaging the data of deSaussure over lethargy intervals and by passing a continuous curve through the intervals.

The unresolved resonance parameters in the ${ }^{235}$ U ENDF/B-III file were modified to yicld the evaluated fission cross section and to reproduce the ENDF/B-H alpha value as closely as possible. The parameters were obtained by adjusing the parameters fo yield the desired fission and alpha values.

The cross sections which were fitted and the s-wave strength function and fission widths resulting from the fitting procedure are given in WARD-42l0T4-1. In this evaluation, the upper energy range for the unresolved resonance parameters was cut off at 25 keV and pointwise data was used above this energy.

## Data Above 25 keV

a. Fission Cross Section

Qualitatively, the experimental data in the energy range $\sim 25 \mathrm{keV}$ to 100 keV falls into two groups: the low fission values of Szabo and Lemley which use ${ }^{6} \mathrm{Li}$ as a standard and the higher fission values such as White and DeSaussure which use hydrogen and ${ }^{10} B$ as a standard. However the recent data of Gwin using a ${ }^{10} B$ standard supports the low fission values. The data of Blons tend to support the lower fission values while the data of Knoll are in good agreement with White's data. The use of ${ }^{6} \mathrm{Li}$ or ${ }^{10} \mathrm{~B}$ as a standard does not appear to be the source of the discrepancy because the ${ }^{10} B$ cross section used for normalization is partially derived from the ${ }^{6}$ Li data consistent with Lemley's normalization.

No clear choice based on the differential data can be made at the present time between the low and high fission values. Integral testing against critical assemblies indicates that use of the lower fission values would require major cross section adjustments - particularly very low capture cross sections for 238
$U$ in order to obtain eigenvalues as close as $1 \%$ less than unity. The latter is particularly true for soft spectrum assemblies typical of interest in LMiFBR design. For this reason, the choice for the present evaluation is based on the data of Perez, White and Knoll.

The experimental data for the $\mathrm{U}^{235}$ fission cross section above 100 keV , as utilized in the present evaluation is based on the measurements of White, Szabo and Smith as corrected by Hansen. The data of Poenitz indicates notably lower fission values than the other measurements in this energy range. The
principal new measurement since the ENDF/B-II evaluation is that of Szabo. The measurements of Kappeler were reported since the present evaluation. The Szabo measurement is the principal source of the differences between the present and FNDF/B evaluations between 0.15 and 1.0 MeV .

Between 1.0 and 10.0 MeV , the only measurements with an accuracy of better then $5 \%$ are those of White at 2.25 and 5.4 MeV with relatively poor shape determination in this energy range. In the present evaluation, the $U^{235}$ fission cross section between 1.0 and 10.0 MeV was evaluated within existing uncertainties in order to increase the cross section by 1 to $3 \%$ primarily to enhance the $\mathrm{U}^{238}$ fission cross section for which the most accurate measurements are relative to $\mathrm{U}^{235}$ fission. Above 10 MeV , the present evaluation is evaluated to obtain a 14 MeV cross section of 2.13 barns.

The detailed structure in the $\mathrm{U}^{235}$ fission cross section is important as it leads to variations of up to about $3 \%$ between groups of multigroup cross sections averaged over quarter lethargy widths typical of many LMFBR calculations and it significantly influences the $\mathrm{Pu}^{239}$ fission cross sections derived from ratios of ${ }^{239} \mathrm{Pu} /{ }^{235} \mathrm{U}$ fission. For the present evaluation, a compromise structure was selected based on the measurements of Perez and Lemley. The evaluated structure was normalized to obtain the 10 keV evaluated average cross section. This structure could possibly be improved when pointwise data from the measurements of Perez and Lemley become available.

## b. Capture Cross Section

The capture cross section above 25 keV was obtaincd by folding the nowly evaluated fission cross section into the ENDF/B-II alpha values. Results of discussions with deSaussure (ORNL) relative to the two sets of ORNL alpha data (the lower values of Weston et al. and the higher values of deSaussuro, of al.) provided no basis of establishing one data set over the other. A weighted averathe of the two sets was used. The alpha evaluation is summari\%ed as follows: l! 40 keV (Schmidt evaluation); $40-60 \mathrm{kcV}$ (joining of new and Schmidt evalualion); $60-200 \mathrm{keV}$ ( $5-7 \%$ higher than Schmidt evaluation); $200-400 \mathrm{keV}$ (smooth joining of new and Schmidt evaluations); above 400 keV (Schmidt evaluation).

## c. Total Cross Section

Above 2 MeV , the data of Glasgow and Foster was used to represent the total cross section. Below 2 McV the total cross section of MAT 1044 was adopled.

## d. Elastic Scattering Cross Section

The elastic scattering cross section was obtained by subtracting from the total cross section the sum of the remaining partial cross sections.


[^0]:    * 

    Work done under the auspices of the United States Atomic Energy Commission

[^1]:    * Work done under the auspices of the United States Atomic Energy Commission

[^2]:    * Work done under the auspices of the United States Atomic Energy Commission

