

# Physics and Results of CASCADE.04 model

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## Physics Models

**Intra-nuclear cascade model**

**Pre-equilibrium (exciton model)**

**Evaporation (Generalized Evaporation Model)**

**Fission model (Fong's Model)**

# Intra-Nuclear Cascade model

## What are the inputs we have?

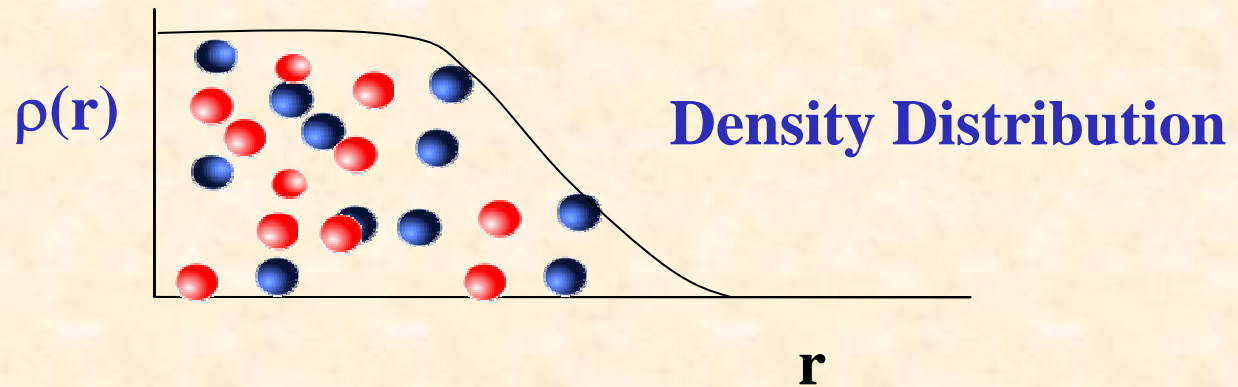
Projectile: 

Charge, mass, energy/momentum

Target: 

Charge, mass, nucleon density distribution

Each nucleon is  
assigned  
position &  
momentum

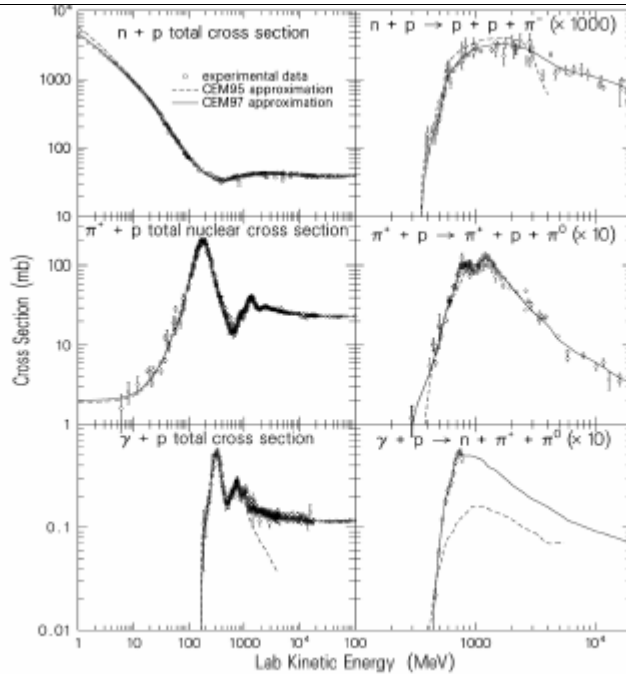
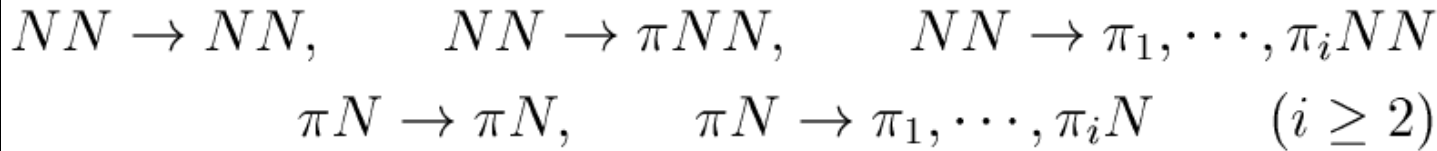


$$\left\{ \begin{array}{l} \rho(r) = \frac{\rho_0}{1 + \exp\left(\frac{r-r_0}{a}\right)} \\ \text{where } r_0 = 1.07A^{1/3} \text{ fm} \\ a = 0.545 \text{ fm} \quad \text{For } A > 10 \\ \rho(r) = \rho_0 \exp\left(-\frac{r^2}{R^2}\right) \quad \text{For } A \leq 10 \end{array} \right\}$$

$$\left\{ \begin{array}{l} P_F(r) = \left(\frac{3\pi^2\rho(r)}{2}\right)^{1/3} \\ E_F(r) = \hbar^2 \frac{(3\pi^2\rho(r))^{2/3}}{2m_N} \end{array} \right\}$$

$$\left\{ \begin{array}{l} V \equiv V_N = E_F + \text{Binding energy} \\ V_\pi = 25 \text{ MeV} \end{array} \right\}$$

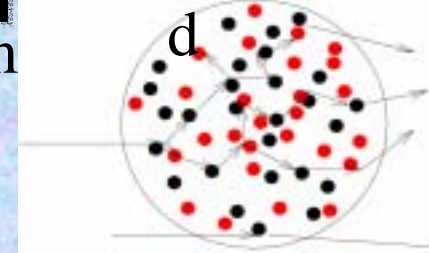
# Intra-Nuclear Cascade model



mcrp Computational Analysis and Simulation (K-3)

Los Alamos

$$\lambda = p/m$$



$$\lambda \ll d$$

$$\lambda \ll \Lambda$$

## Reactions

$p + p = p + p$	Isotropic $E < 0.46$ GeV
$p + p = p + p$	$0.46 < E < 2.8$ GeV
$p + p = p + p$	$2.8 < E < 10.0$ GeV
$p + n = p + n$	$E < 0.97$ GeV
$\pi^+ + p = \pi^+ + p$	$E < 80.0$ MeV
$\pi^+ + p = \pi^+ + p$	$80 < E < 300.0$ MeV
$\pi^+ + p = \pi^+ + p$	$0.3 < E < 1.0$ GeV
$\pi^+ + p = \pi^+ + p$	$1.0 < E < 2.4$ GeV

## Cross-section

$$\left\{ \begin{aligned}
 \cos(\theta) &= 2\xi^{1/2} \left[ \sum_{n=0}^N a_n \xi^n + \left(1 - \sum_{n=0}^N a_n\right) \xi^{N+1} \right] - 1 \\
 a_n &= \sum_{k=0}^N a_{nk} E^k \\
 N=3, M=3
 \end{aligned} \right. \quad \text{Angular distribution}$$

# Pre-equilibrium model (Exciton model)

**Sharp cut off energy (7 MeV) is the criteria to close INC**

n, p, d, t,  $^3\text{He}$ , and  $^4\text{He}$  emission

**Probability of emission is calculated as given below**

$$\Gamma_j(p, h, E) = \int_{V_j^*}^{E-B_j} \lambda_c^j(p, h, E, T) dT,$$
$$\lambda_c^j(p, h, E, T) = \frac{2s_j + 1}{\pi^2 \hbar^3} \mu_j \mathfrak{R}_j(p, h) \frac{\omega(p-1, h, E - B_j - T)}{\omega(p, h, E)} T \sigma_{inv}(T)$$

p=particle, h=hole, n=p+h is exciton number, s=spin,

$\sigma_{inv}$ =cross-section, E=excitation energy, B=binding energy

$$\lambda_+(n_{eq}, E) = \lambda_-(n_{eq}, E)$$
$$n_{eq} \simeq \sqrt{2gE}$$

M. Veselsky, Nucl. Phys. A705 (2002) 193

## Differences with CEM

**Kalbach Systematics below 210 MeV**

the widths for complex-particle emission were changed by fitting the probability of several excitons to "coalesce" into a complex particle that may be emitted during the preequilibrium stage to available experimental data on reactions induced by protons and neutrons;

## Earlier assumption

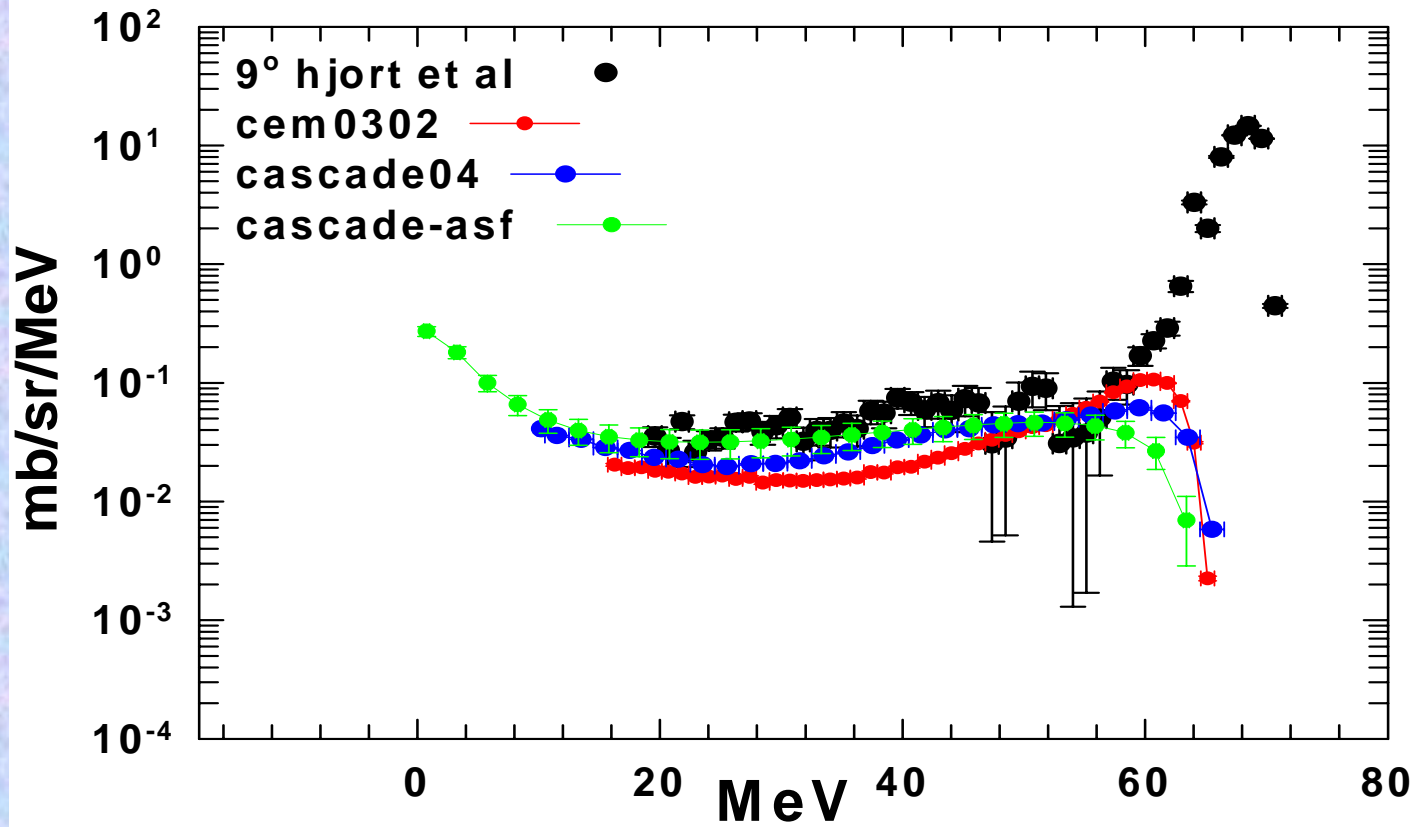
$$P_{pre} = 1 \text{ or } 0$$

Now

$$P_{pre}(n/n_{eq}) = 1 - \exp\left(-\frac{(n/n_{eq} - 1)^2}{2\sigma_{pre}^2}\right)$$

$\sigma_{pre} = 0.22$  (0.4 in CEM)

# P (65MeV)+<sup>56</sup>Fe



# Evaporation model

$$P_j(\epsilon)d\epsilon = g_j \sigma_{inv}(\epsilon) \frac{\rho_d(E - Q - \epsilon)}{\rho_i(E)} \epsilon d\epsilon$$

$$P_f = \frac{\int_0^{U-B_f} \rho(U - B_f - \delta - E) dE}{\rho_j(E)}$$

$$\rho(E) = \frac{c_1 \exp(2\sqrt{a(E - \delta)})}{a^{1/4} (E - \delta)^{5/4}}$$

$$\rho(E) = c_2 \exp((E - E_0)/T)$$

**S is assumed 0 at saddle**

$Z_j$	Ejectiles							
0	n							
1	p	d	t					
2	<sup>3</sup> He	<sup>4</sup> He	<sup>6</sup> He	<sup>8</sup> He				
3	<sup>6</sup> Li	<sup>7</sup> Li	<sup>8</sup> Li	<sup>9</sup> Li				
4	<sup>7</sup> Be	<sup>9</sup> Be	<sup>10</sup> Be	<sup>11</sup> Be	<sup>12</sup> Be			
5	<sup>8</sup> B	<sup>10</sup> B	<sup>11</sup> B	<sup>12</sup> B	<sup>13</sup> B			
6	<sup>10</sup> C	<sup>11</sup> C	<sup>12</sup> C	<sup>13</sup> C	<sup>14</sup> C	<sup>15</sup> C	<sup>16</sup> C	
7	<sup>12</sup> N	<sup>13</sup> N	<sup>14</sup> N	<sup>15</sup> N	<sup>16</sup> N	<sup>17</sup> N		
8	<sup>14</sup> O	<sup>15</sup> O	<sup>16</sup> O	<sup>17</sup> O	<sup>18</sup> O	<sup>19</sup> O	<sup>20</sup> O	
9	<sup>17</sup> F	<sup>18</sup> F	<sup>19</sup> F	<sup>20</sup> F	<sup>21</sup> F			
10	<sup>18</sup> Ne	<sup>19</sup> Ne	<sup>20</sup> Ne	<sup>21</sup> Ne	<sup>22</sup> Ne	<sup>23</sup> Ne	<sup>24</sup> Ne	
11	<sup>21</sup> Na	<sup>22</sup> Na	<sup>23</sup> Na	<sup>24</sup> Na	<sup>25</sup> Na			
12	<sup>22</sup> Mg	<sup>23</sup> Mg	<sup>24</sup> Mg	<sup>25</sup> Mg	<sup>26</sup> Mg	<sup>27</sup> Mg	<sup>28</sup> Mg	

$$a(A_d, Z_d, E) = A_d(0.134 - 1.2110^{-04} A_d)(1 + \frac{S}{E}(1 - \exp(-0.061E)))$$

$$a_f = a_n \left\{ \begin{array}{l} 1.041 + 0.00915X^2 - 0.0005977X^3 \text{ For } Z < 78 \\ 1.0196 + 0.00896X^2 - 0.000585X^3 \text{ For } 78 < Z < 85 \\ 0.9445 + 0.0083X^2 - 0.000542X^3 \text{ For } Z > 85 \end{array} \right\}$$

**Fission barrier  
from Myers and  
Swiatecki  
systematics**

# Fission model

$$\Omega(E_1) = c_1 \exp(2\sqrt{a_1 E_1})$$

$$\Omega(E_2) = c_2 \exp(2\sqrt{a_2 E_2})$$

$$E_1 : E_2 = a_1 T^2 : a_2 T^2 = a_1 : a_2$$

***scission point model***

$$\Omega(E) = c_1 c_2 \int_0^E \exp(2\sqrt{(a_1 + a_2)(E_1 + E_2)})$$

**E = M\*(A, Z)-M(A1, Z1)-M(A2, Z2) - coulomb energy – deformation energy**

$$V_C = \frac{1.44 Z_1 Z_2}{R_{12}}$$

$$E_1 : E_2 = a_1 T^2 : a_2 T^2 = a_1 : a_2$$

$$R_{12} = R_0(1 + \alpha_{2i}(1 - \frac{3}{5}X_i) + \alpha_{3i}(1 - \frac{3}{7}X_i^2))$$

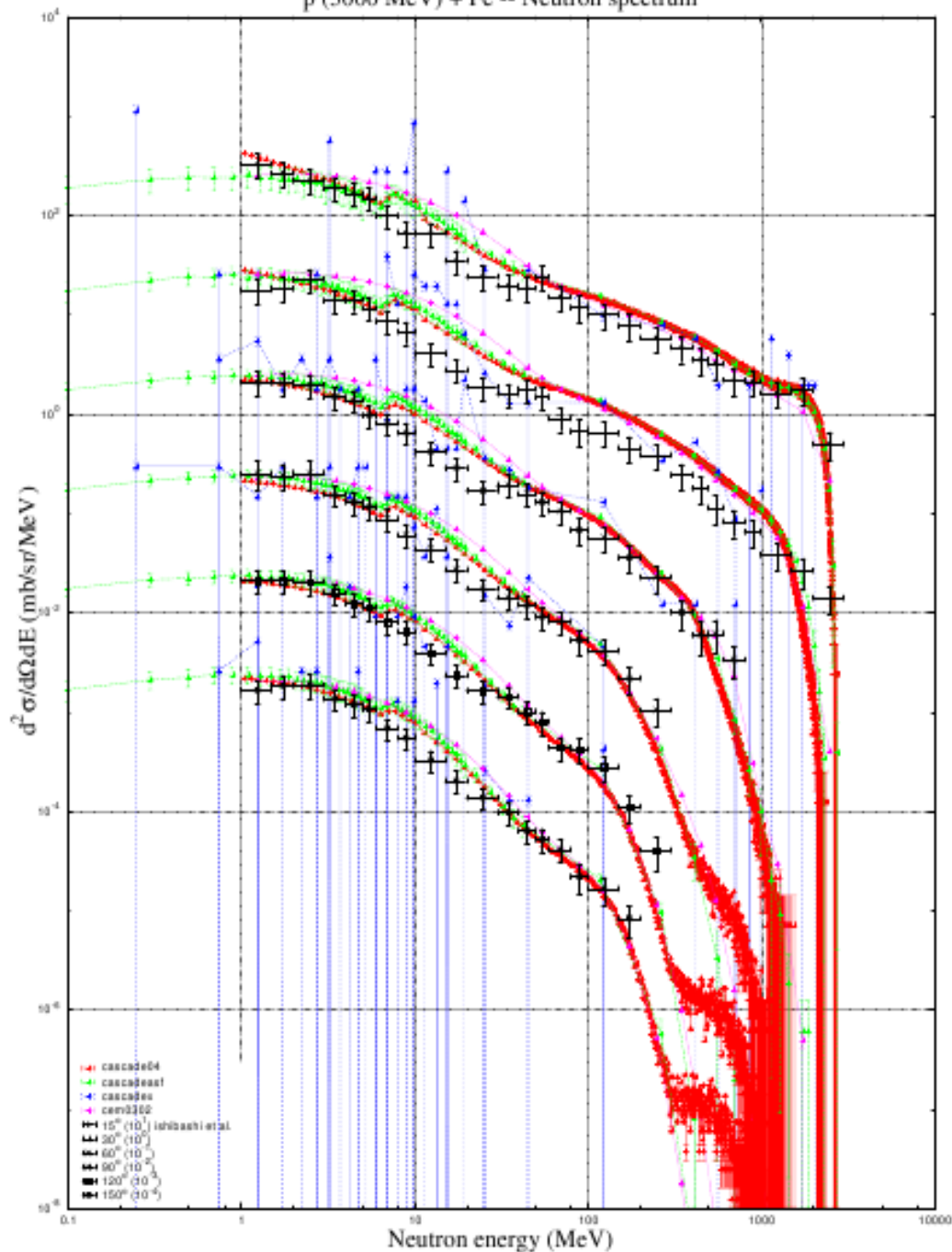
$$X_i = \frac{R_0(i)}{\sum_i R_0(i)(1 + \alpha_{2i} + \alpha_{3i} - \frac{9}{35}\alpha_{2i}\alpha_{3i})}$$

$$R_0(i) = R_0 A_i^{1/3}$$

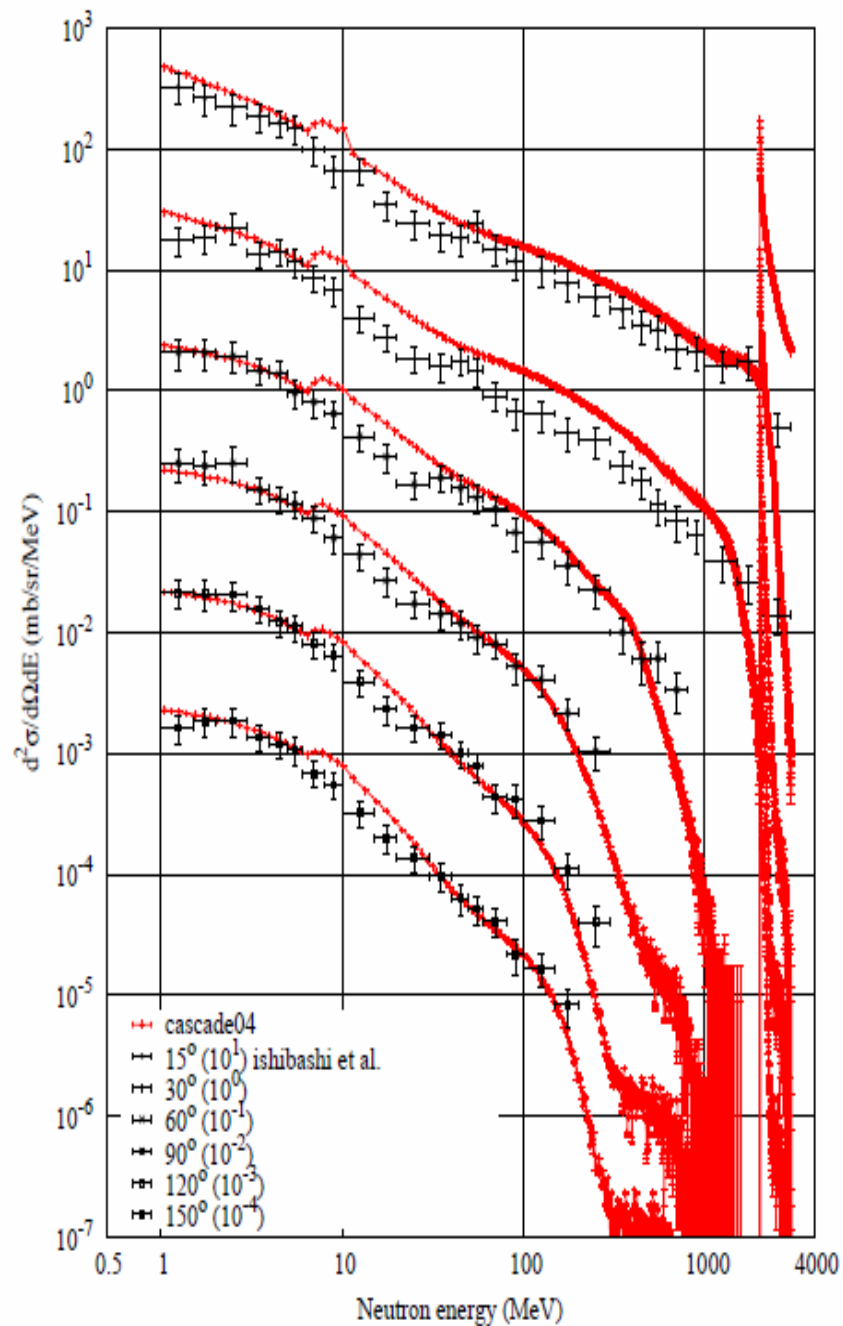
$$\text{and } R_0 = 1.3$$

$$D(i) = (0.4E_S - 0.2E_C)\alpha_{2i}^2 + (0.7144E_S - 0.204E_C)\alpha_{3i}^2$$

p (3000 MeV) + Fe -- Neutron spectrum

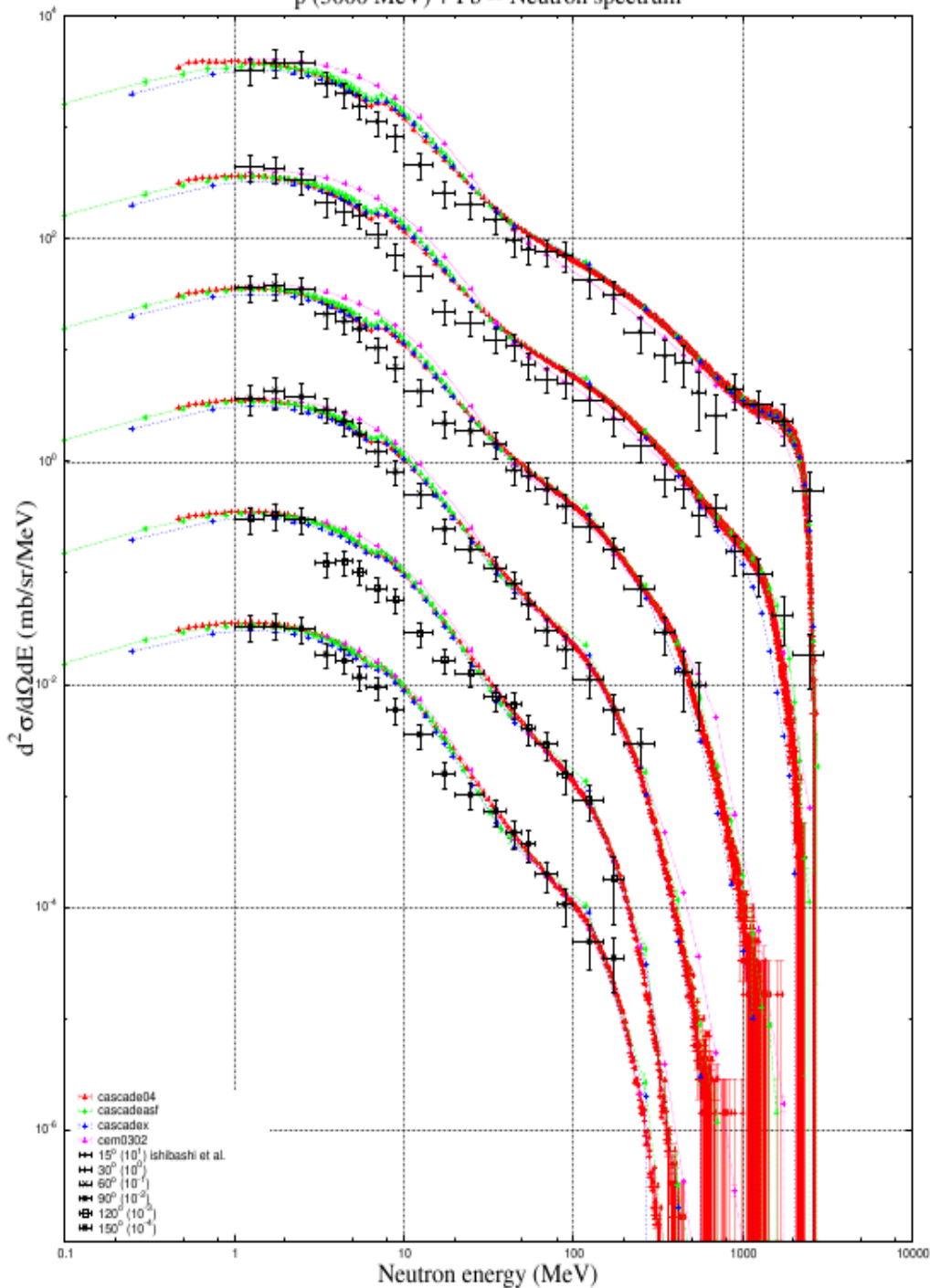


p (3000 MeV) + Fe -- Neutron spectrum

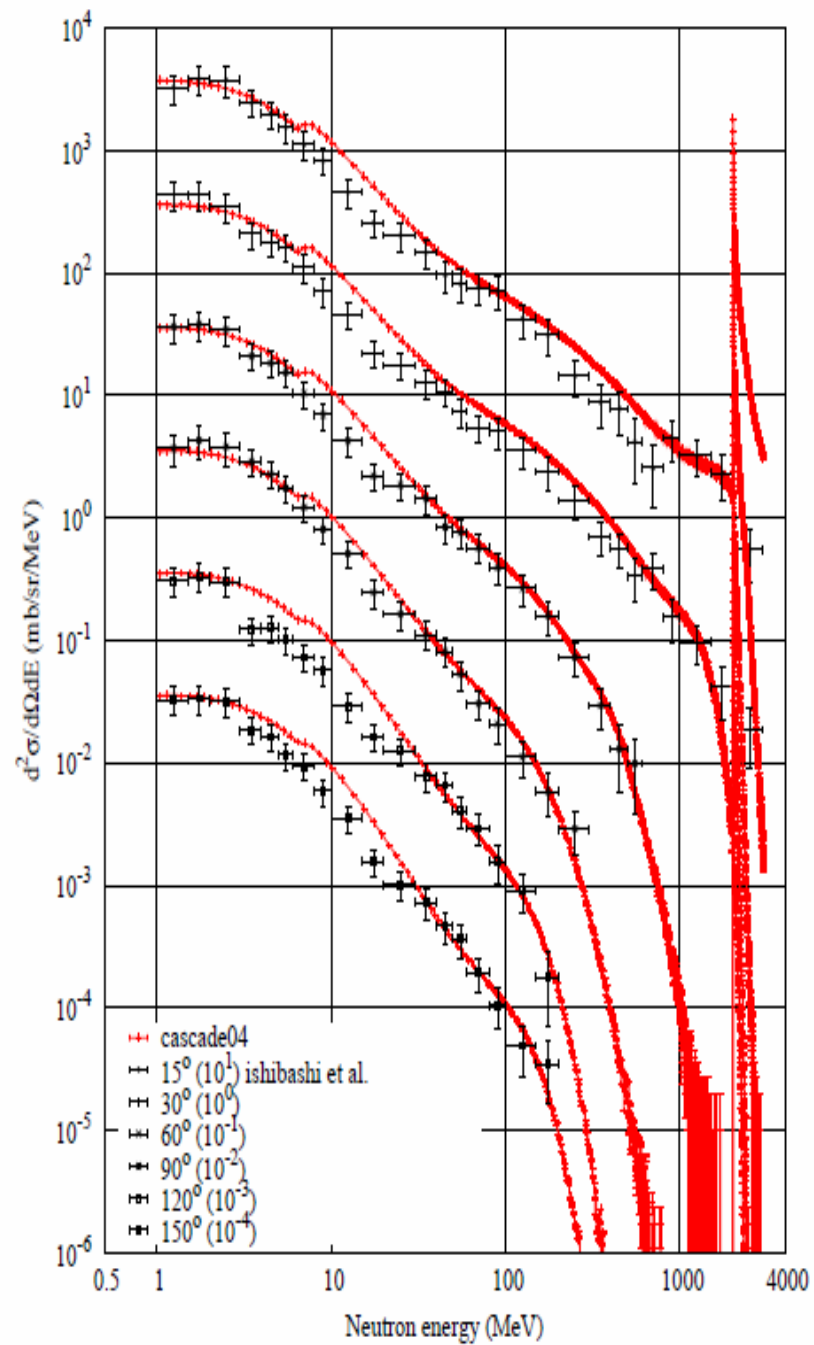




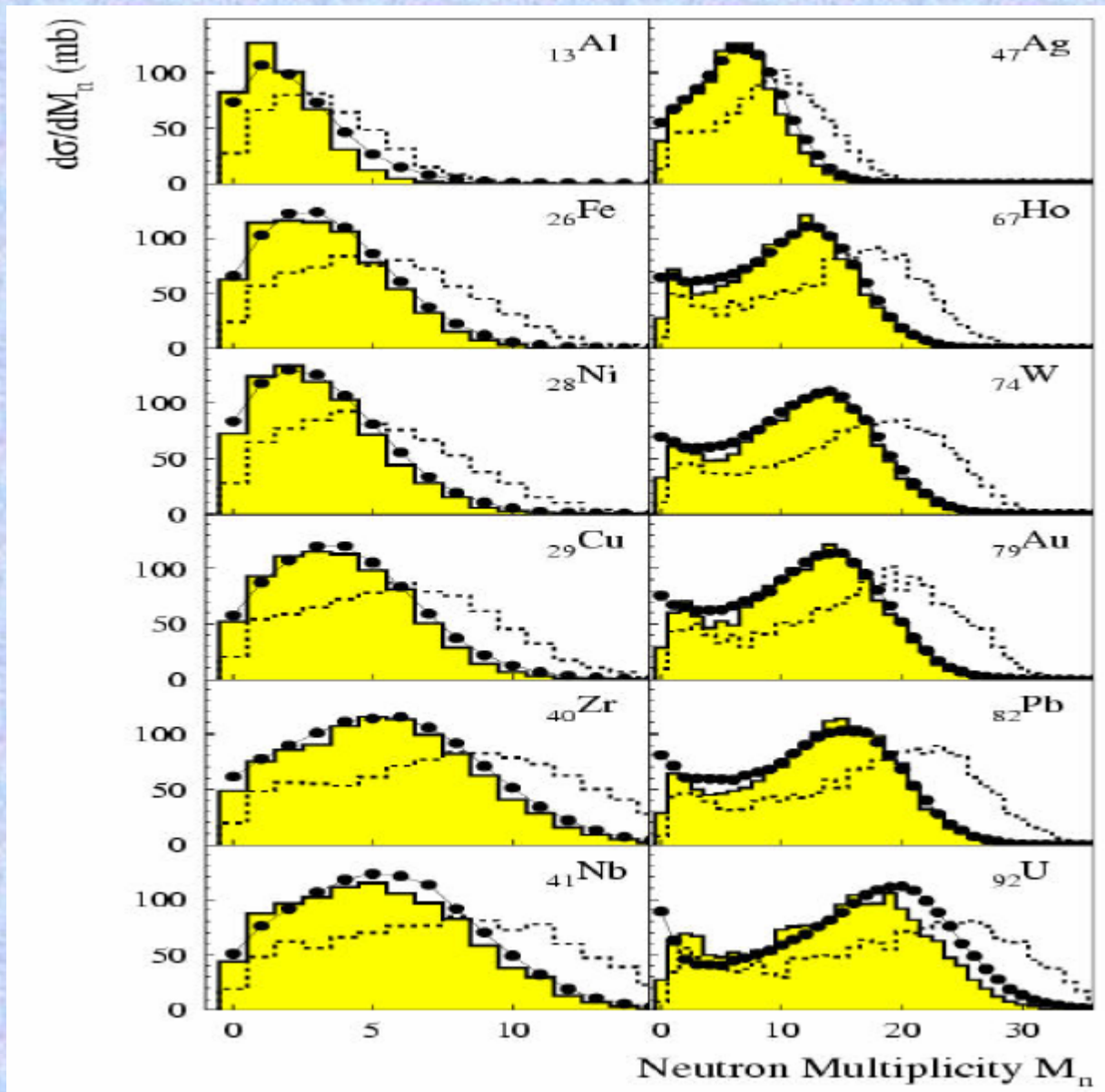
p (3000 MeV) + Pb -- Neutron spectrum



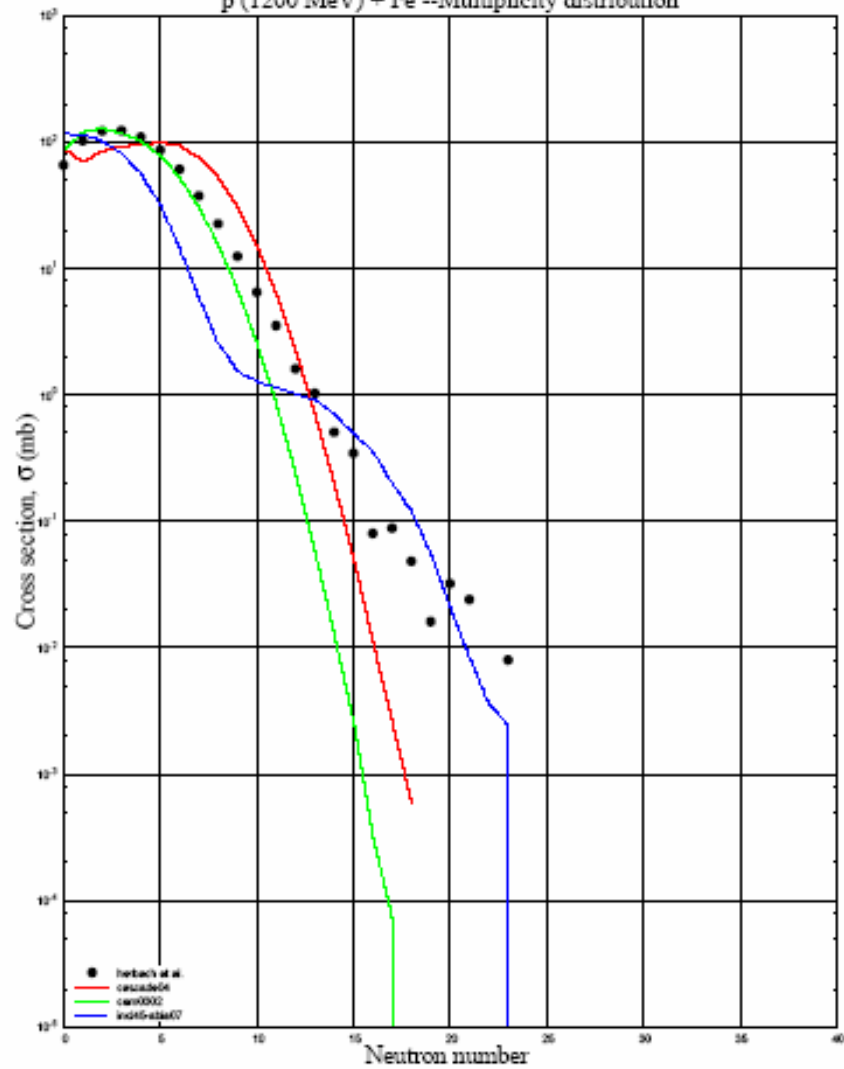
p (3000 MeV) + Pb -- Neutron spectrum



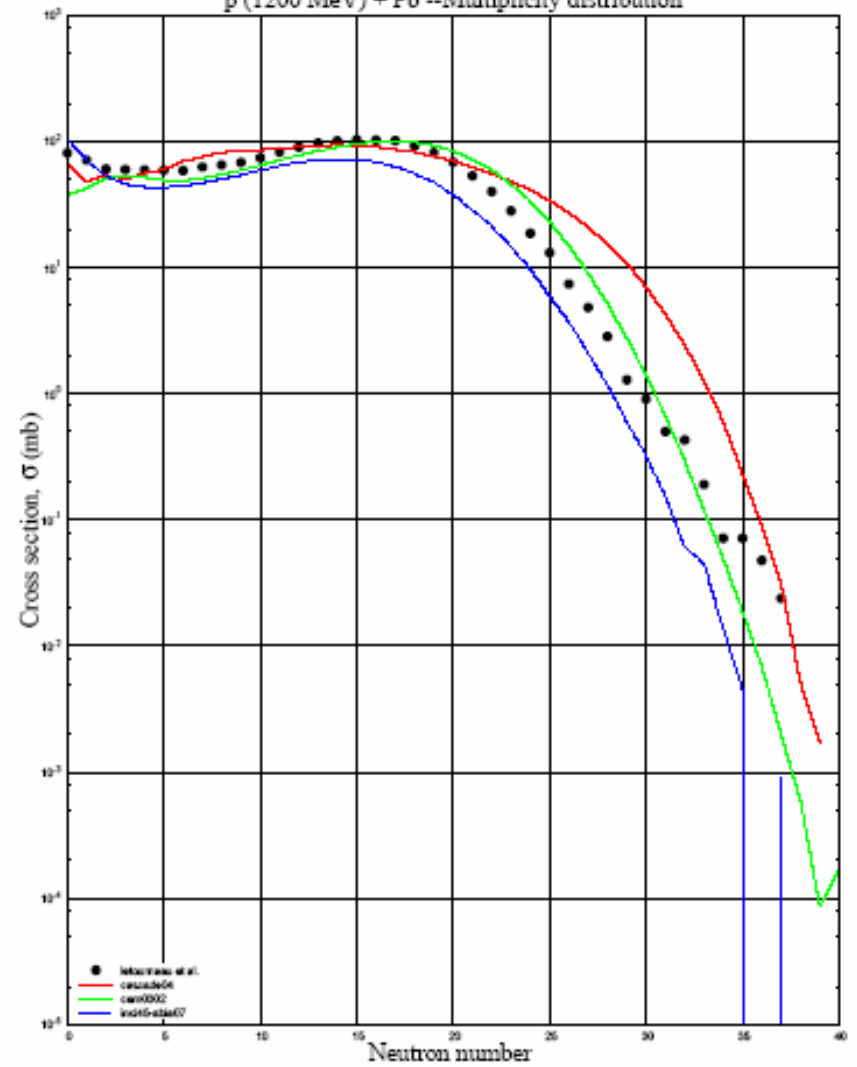
# The multiplicity data impose strong boundary on full code but still not on individual models



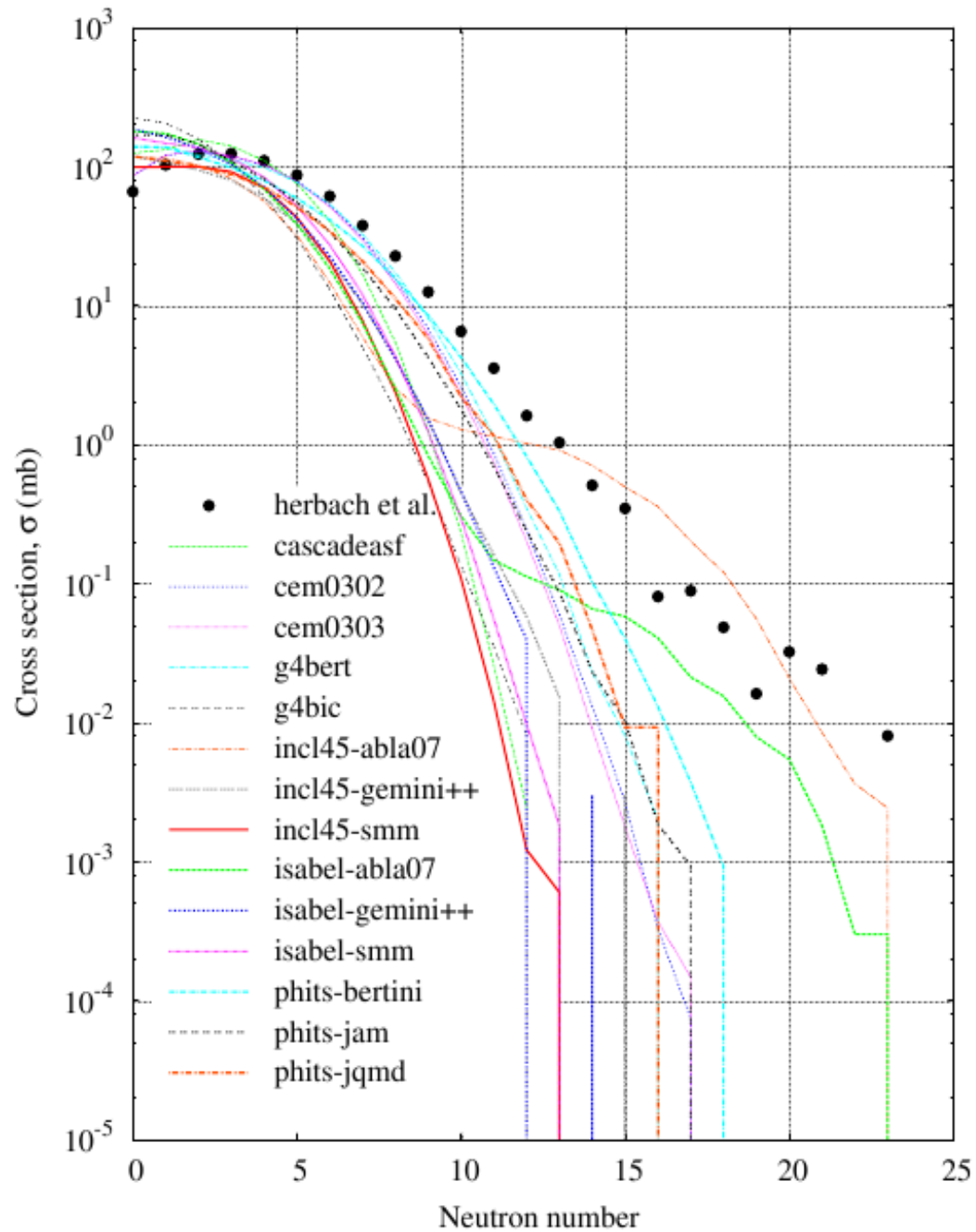
p (1200 MeV) + Fe -- Multiplicity distribution



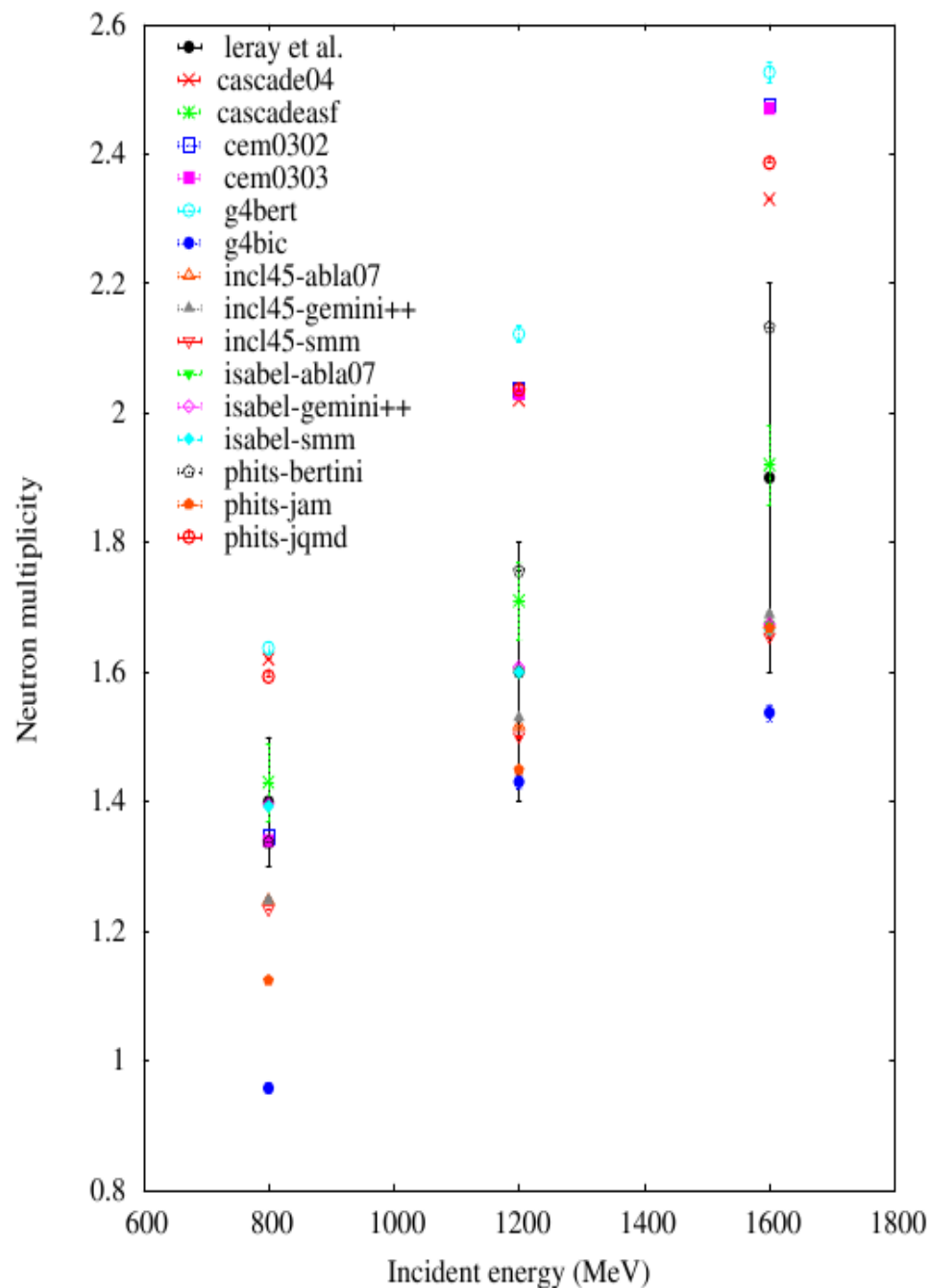
p (1200 MeV) + Pb -- Multiplicity distribution



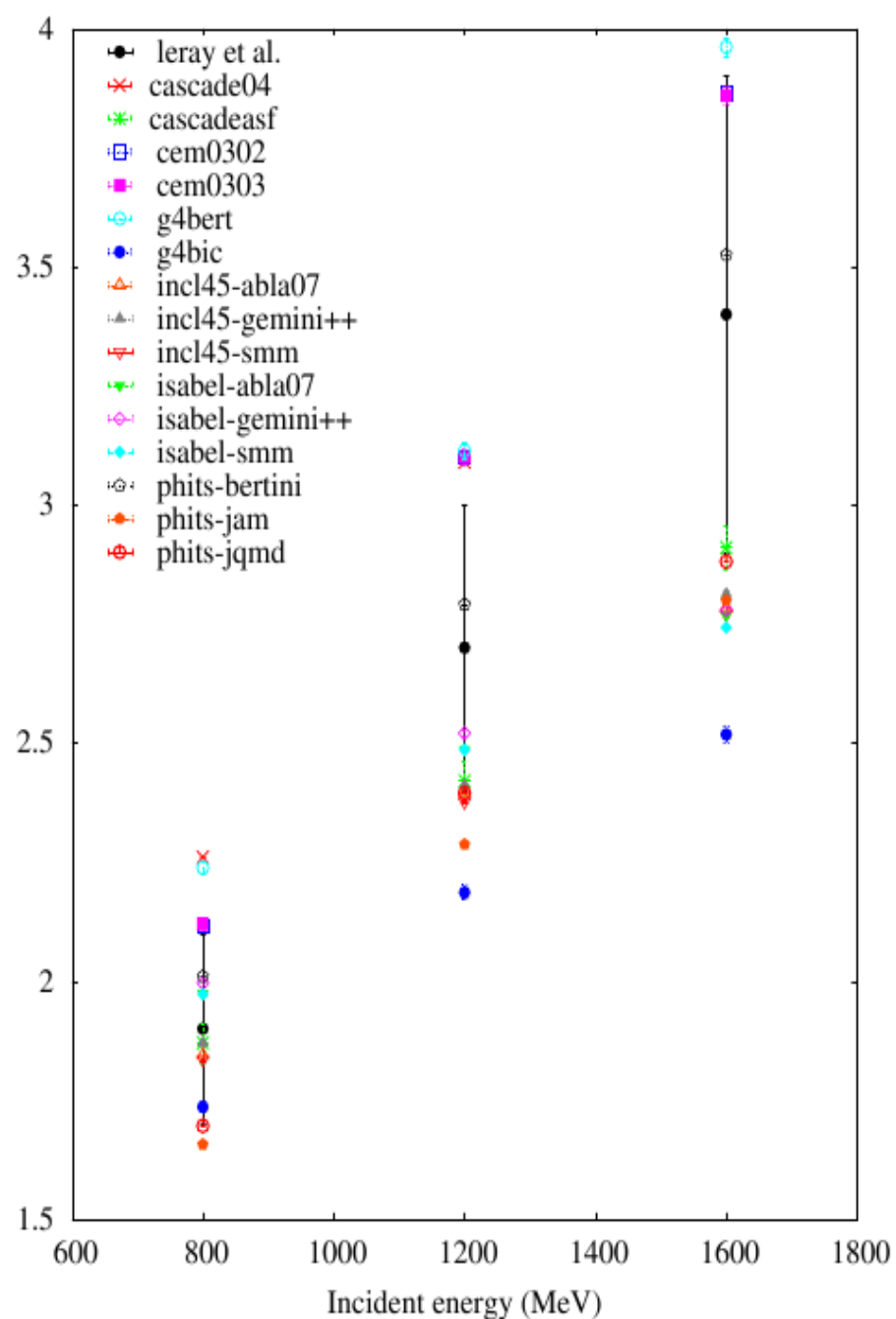
p (1200 MeV) + Fe -- Multiplicity distribution



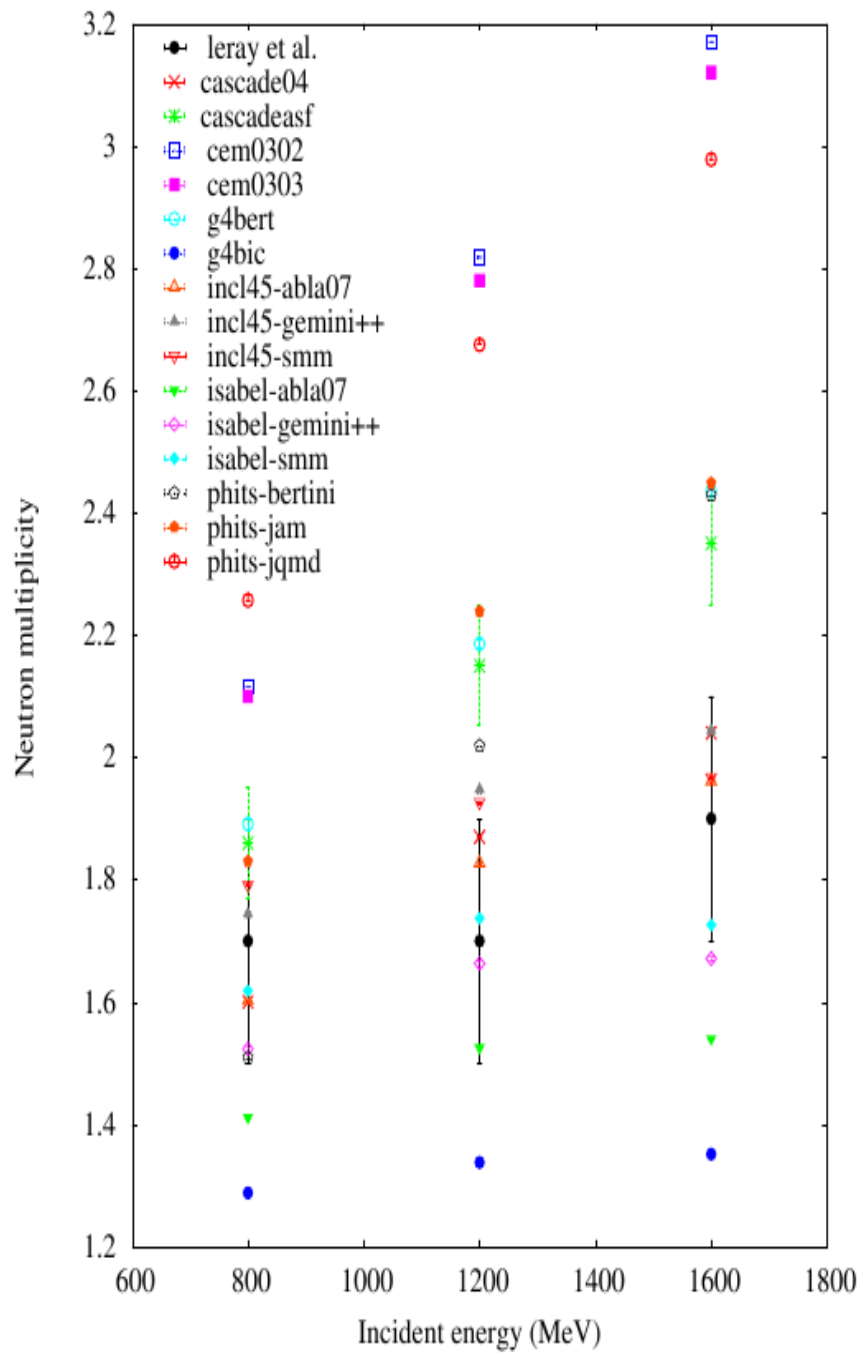
p + Fe -- (20+ MeV) Avg. neutron multiplicities



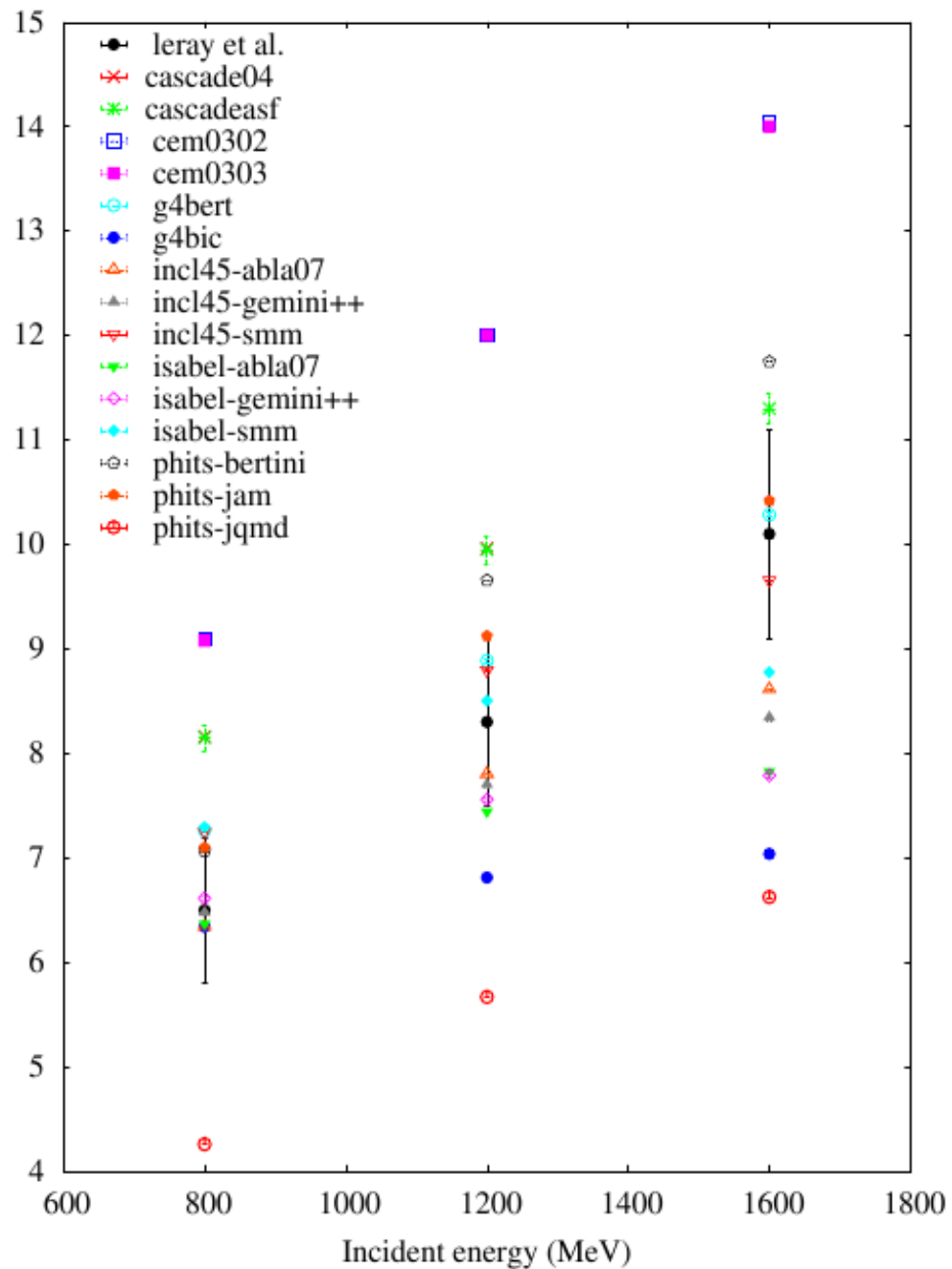
p + Pb -- (20+ MeV) Avg. neutron multiplicities



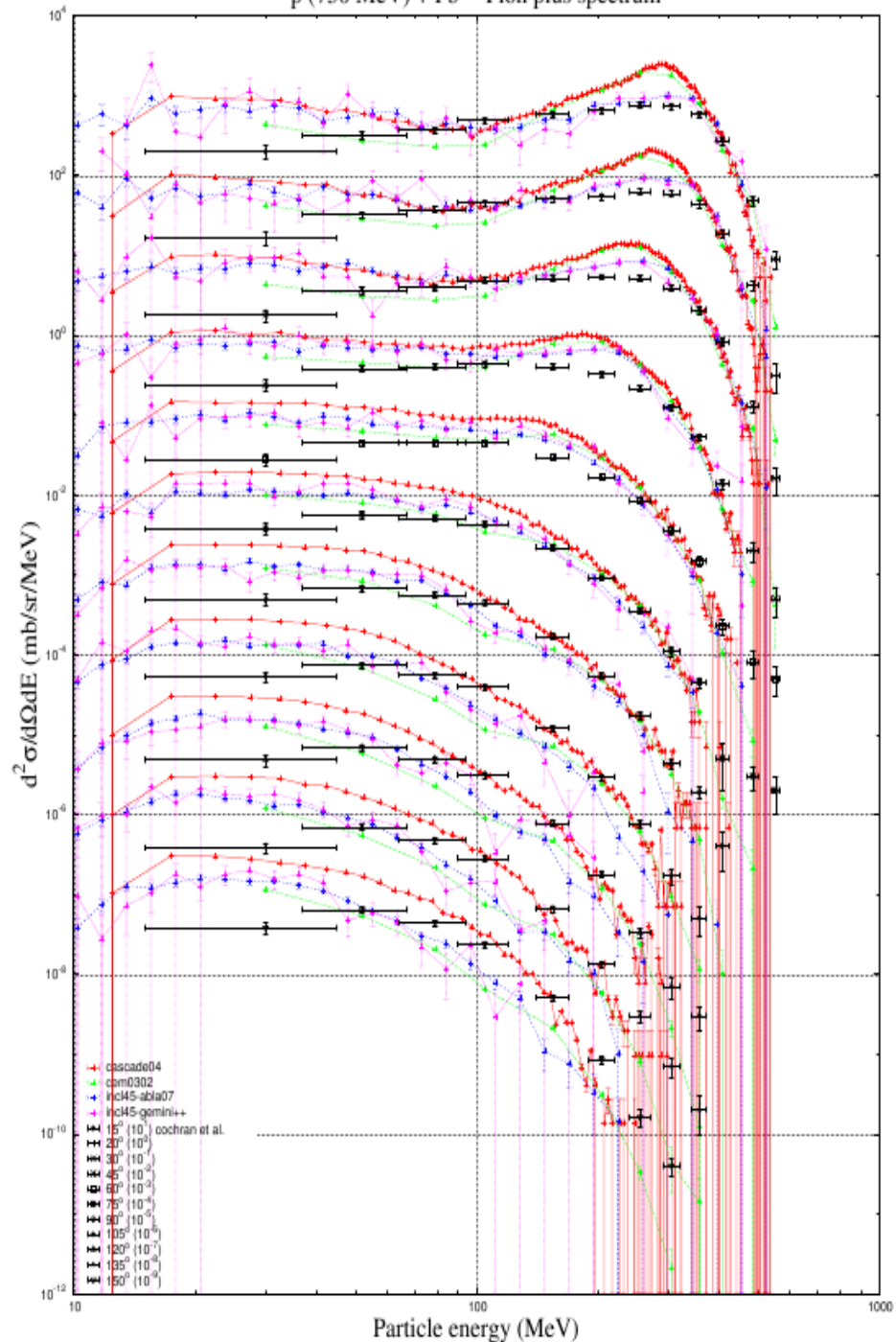
p + Fe -- (2-20 MeV) Avg. neutron multiplicities



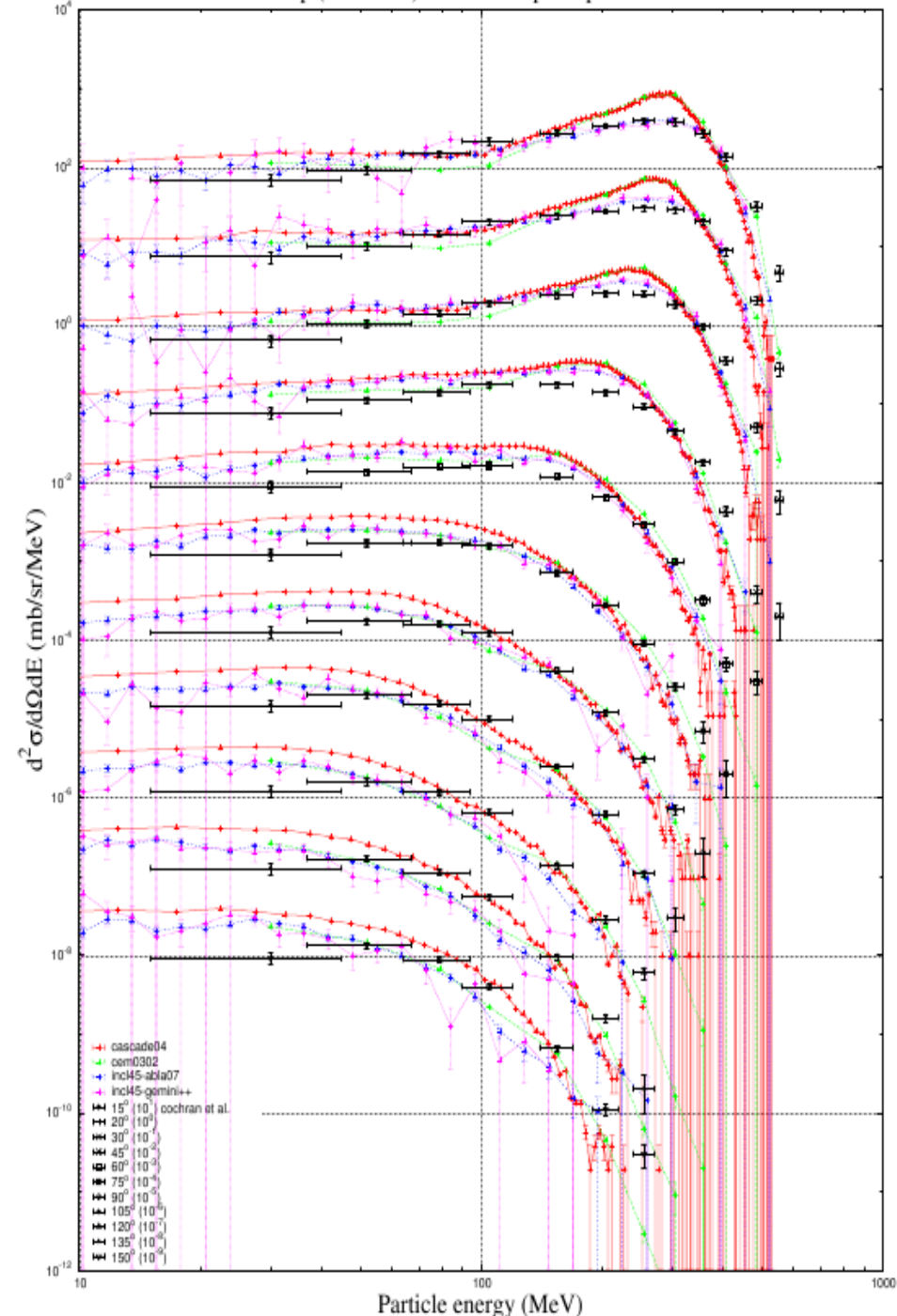
p + Pb -- (2-20 MeV) Avg. neutron multiplicities



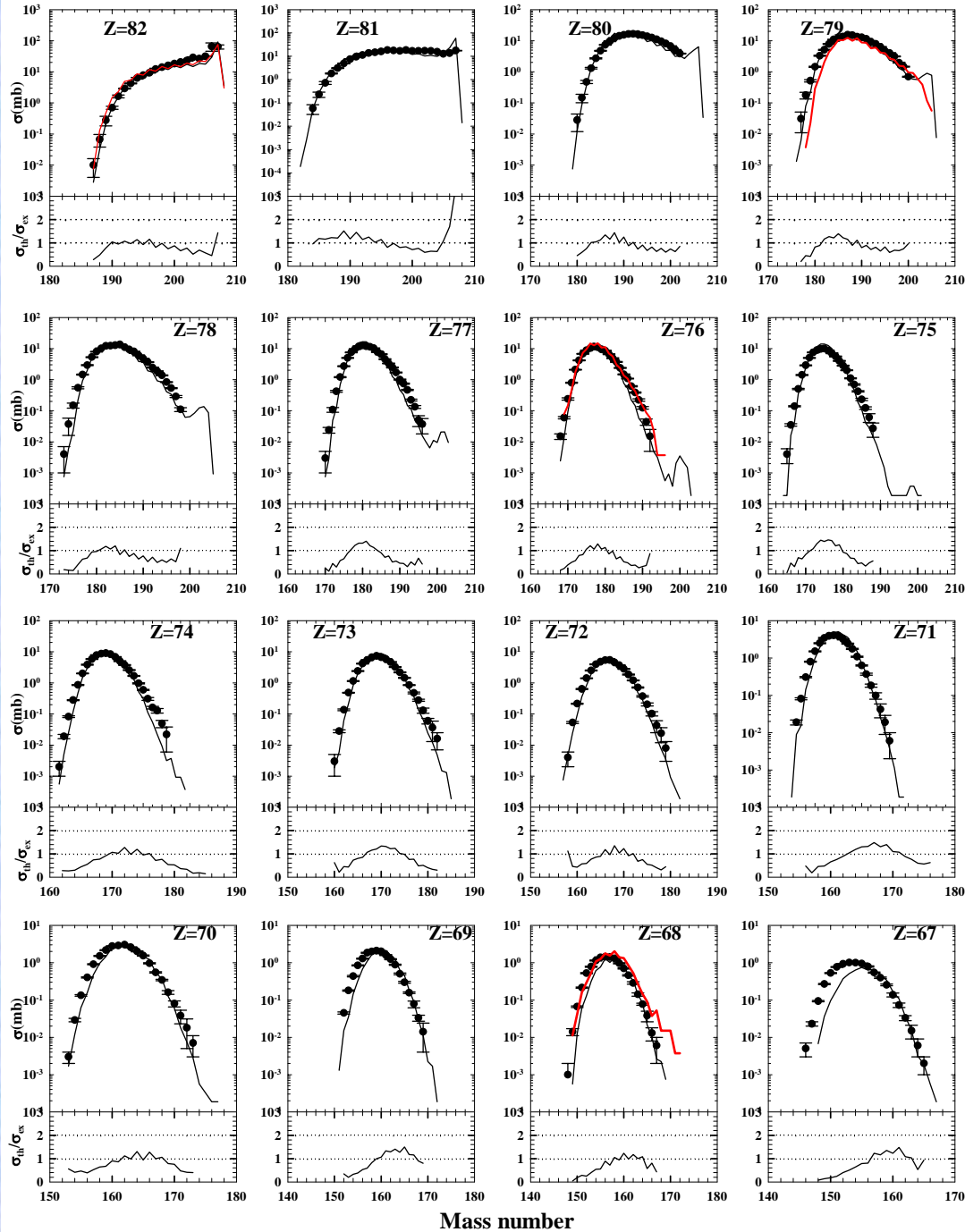
p (730 MeV) + Pb -- Pion plus spectrum



p (730 MeV) + C -- Pion plus spectrum



P (1 GeV) + <sup>208</sup>Pb



Earlier assumption

$$P_{pre} = 1 \text{ or } 0$$

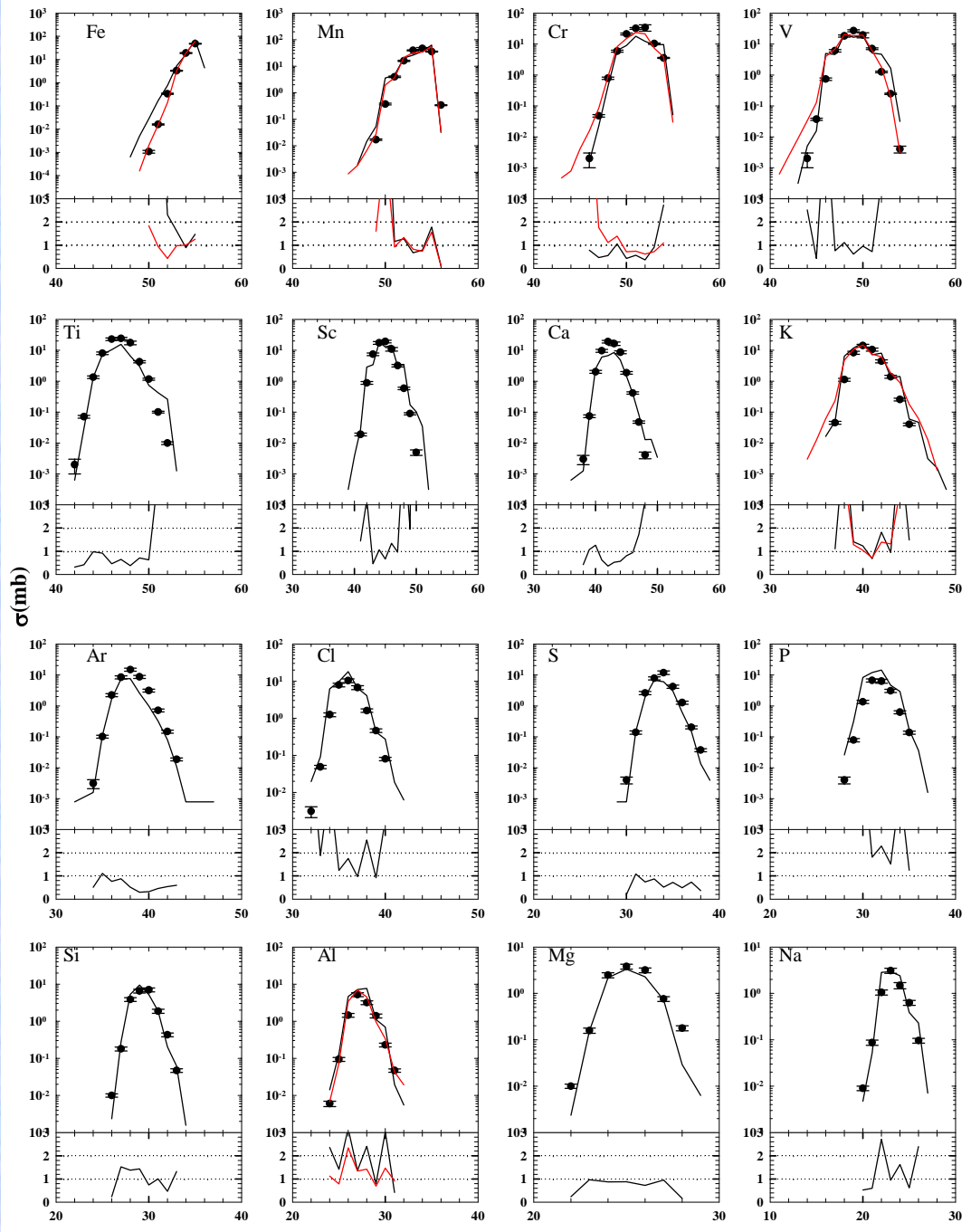
Now

$$P_{pre}(n/n_{eq}) = 1 - \exp\left(-\frac{(n/n_{eq} - 1)^2}{2\sigma_{pre}^2}\right)$$

$P_{pre} = 0.22$  (0.4 in CEM)

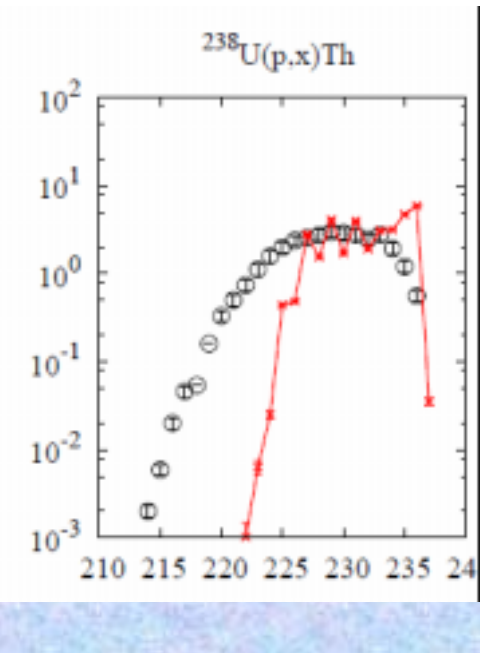
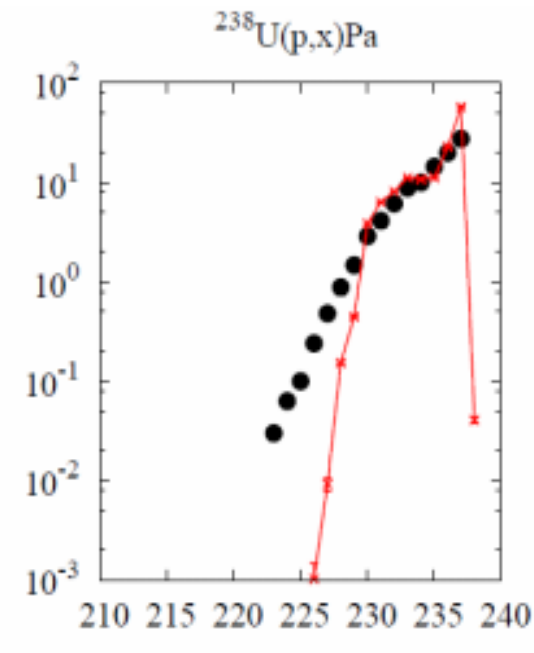
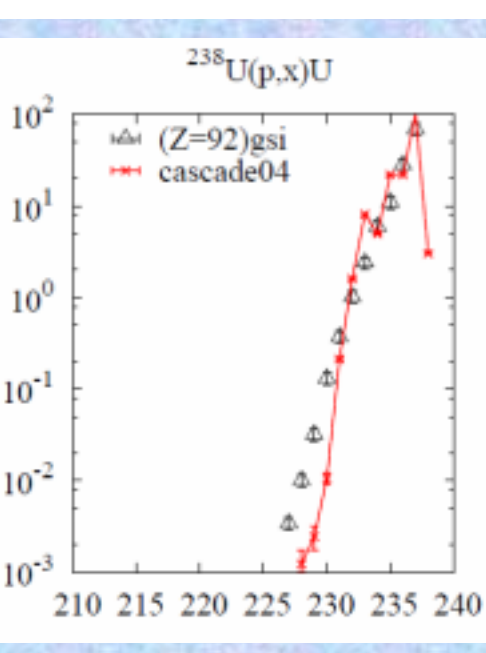
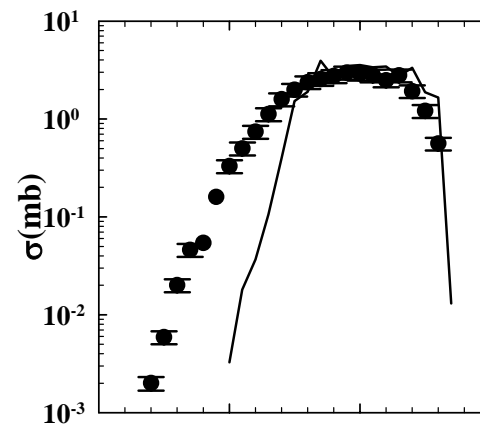
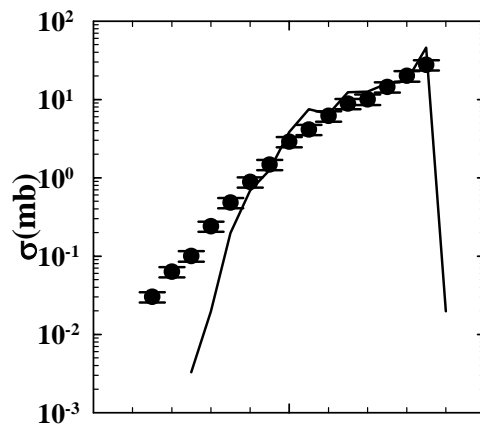
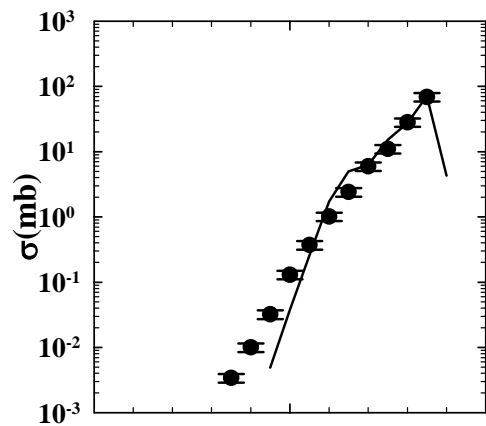


# P (1GeV)+<sup>56</sup>Fe

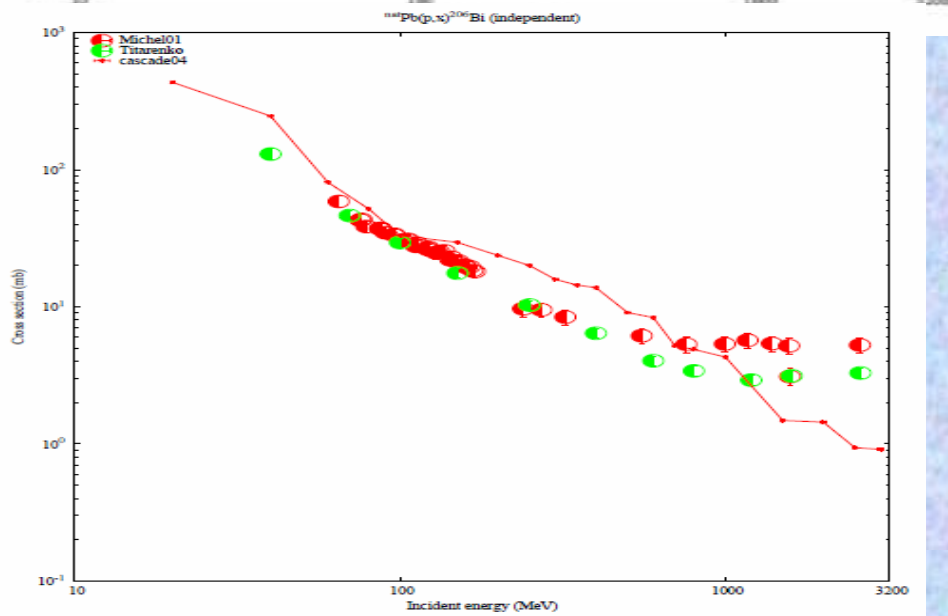
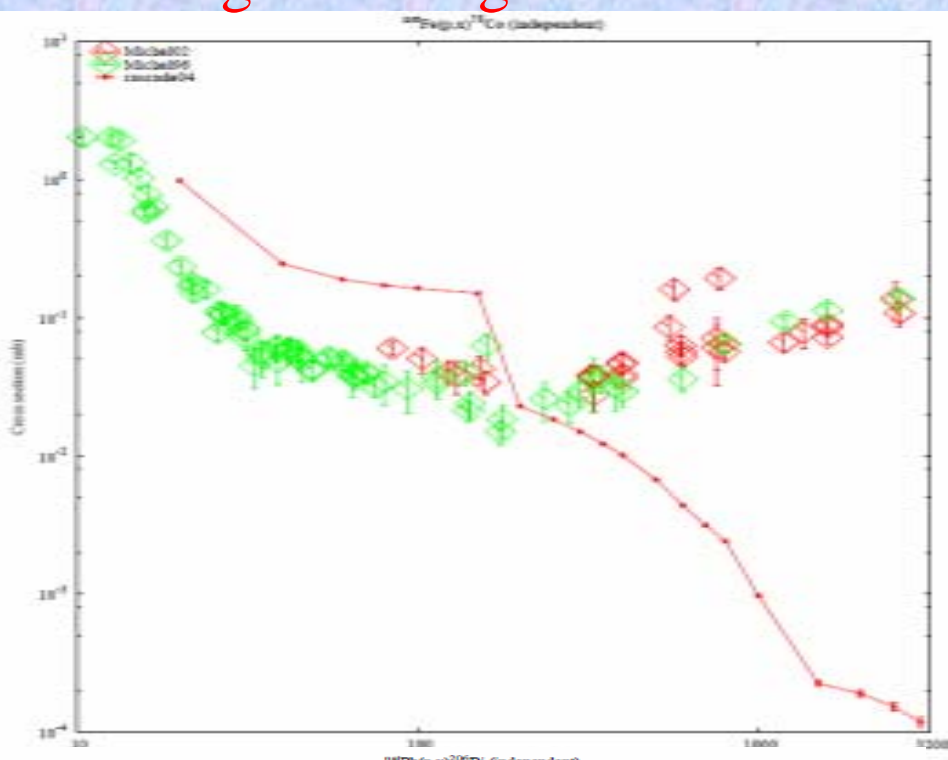


Mass number

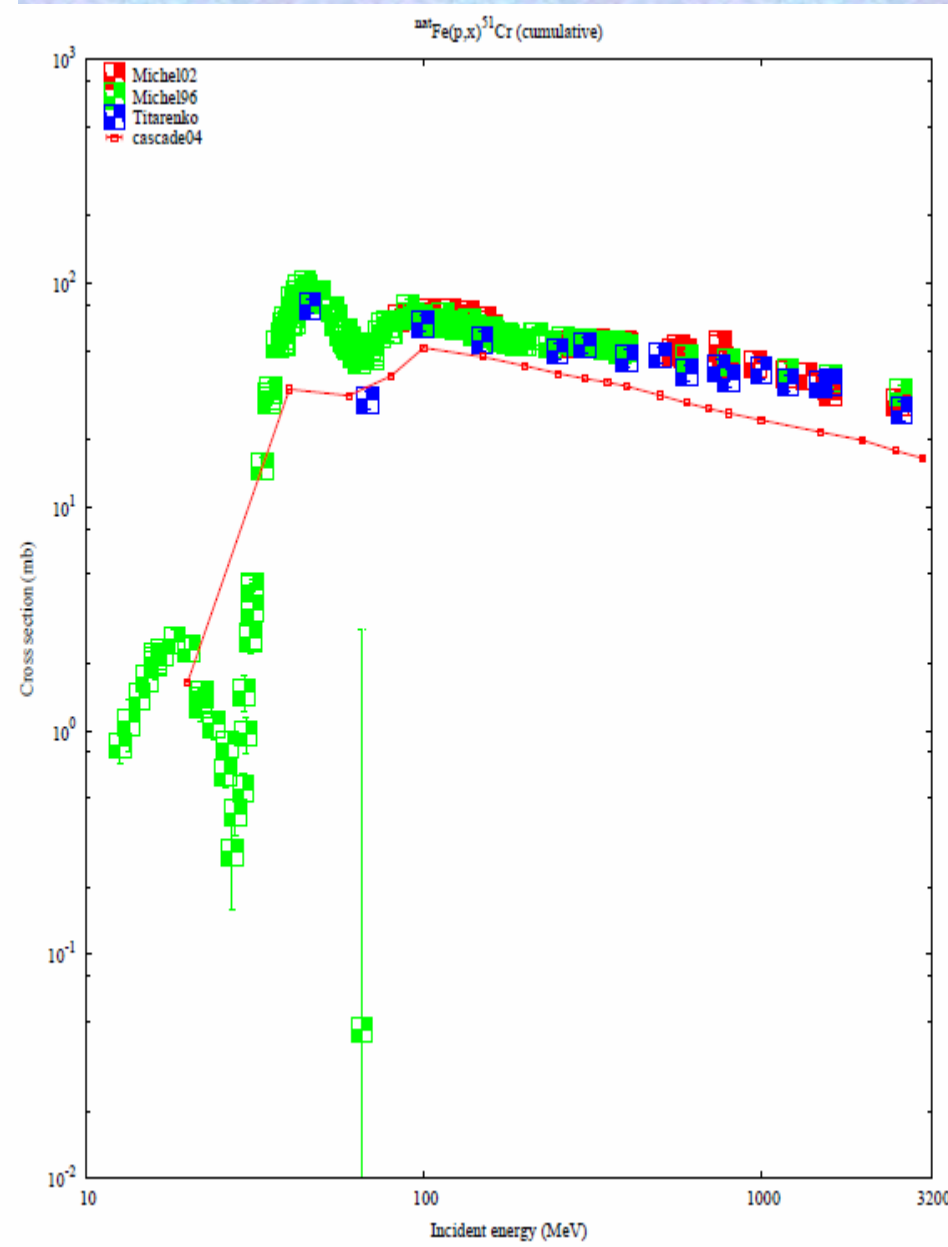
# P (1GeV)+<sup>238</sup>U



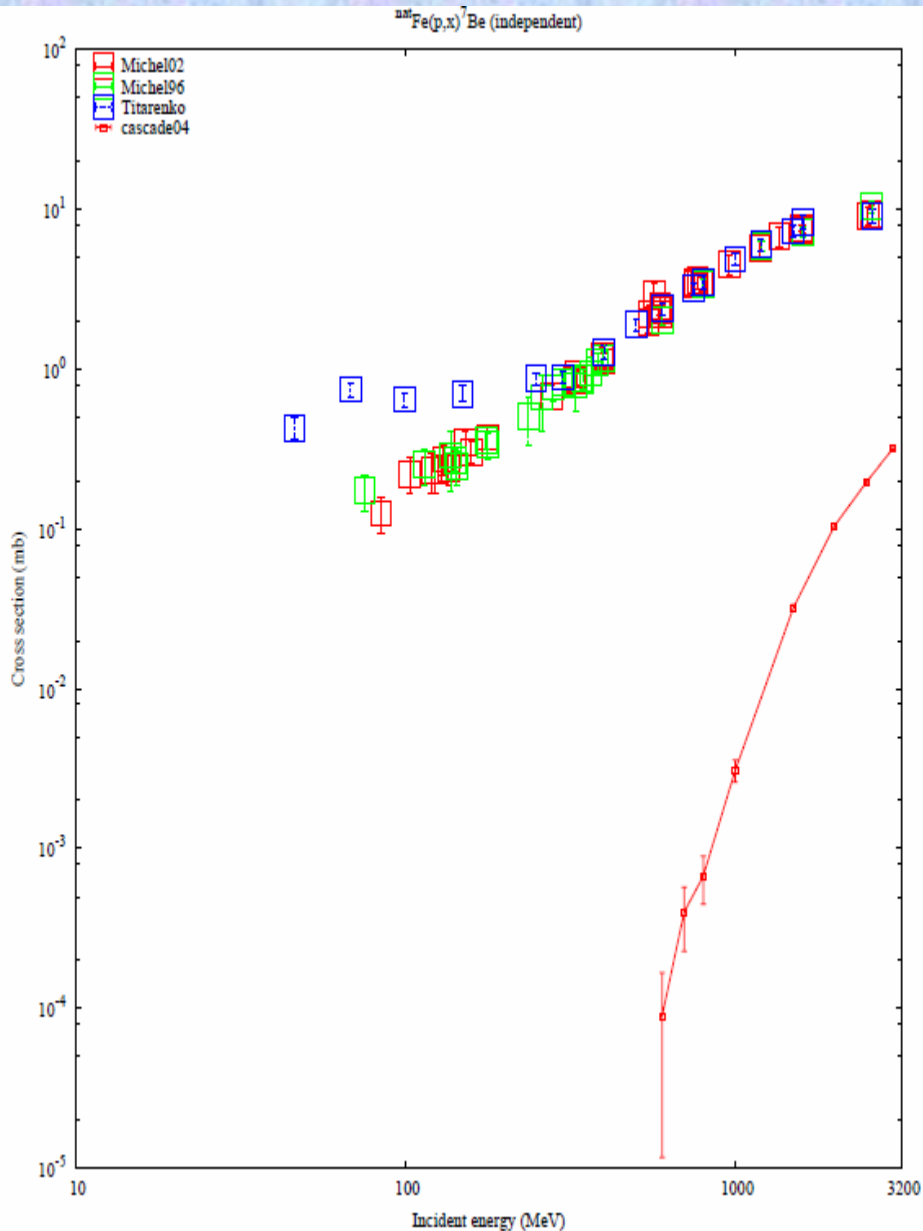
# Charge exchange



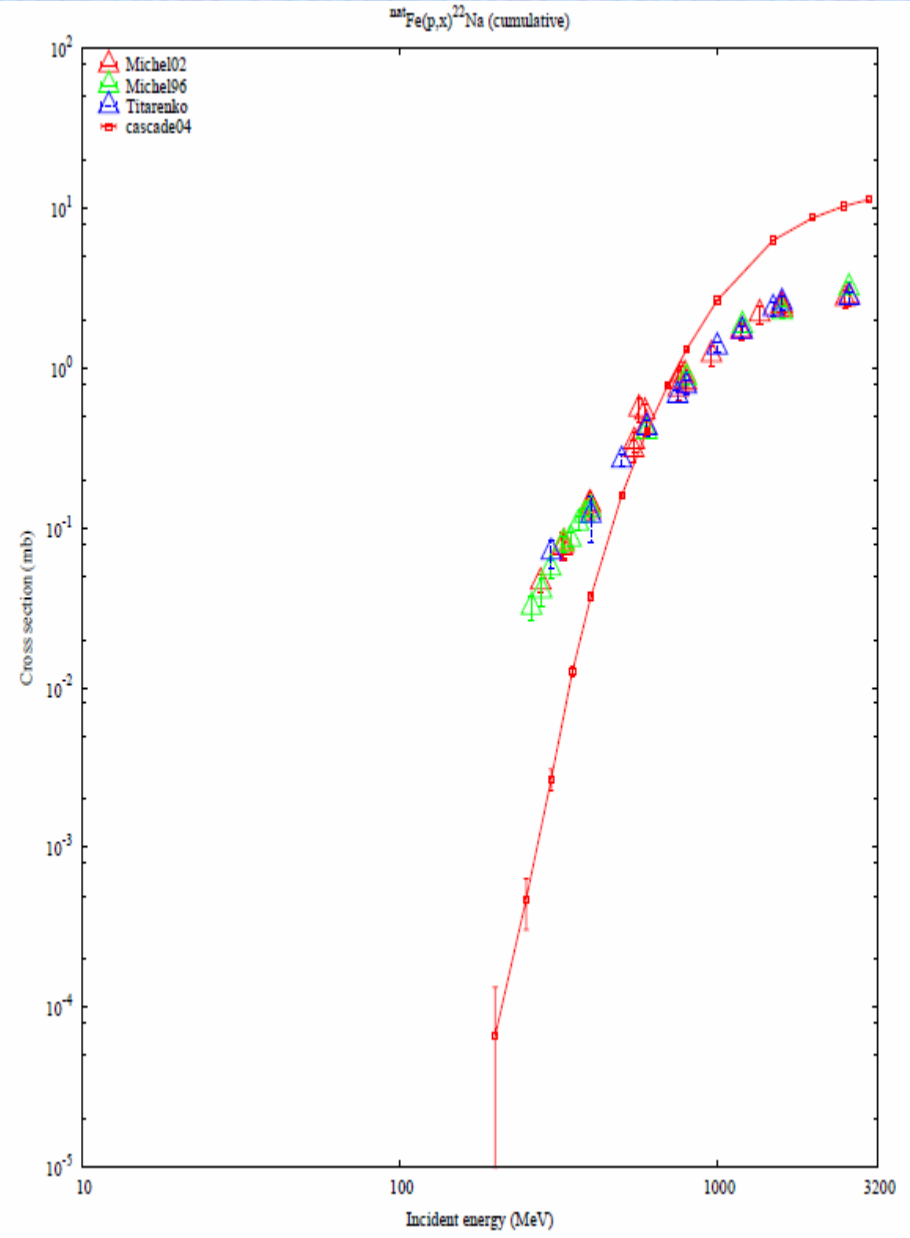
# Close to target

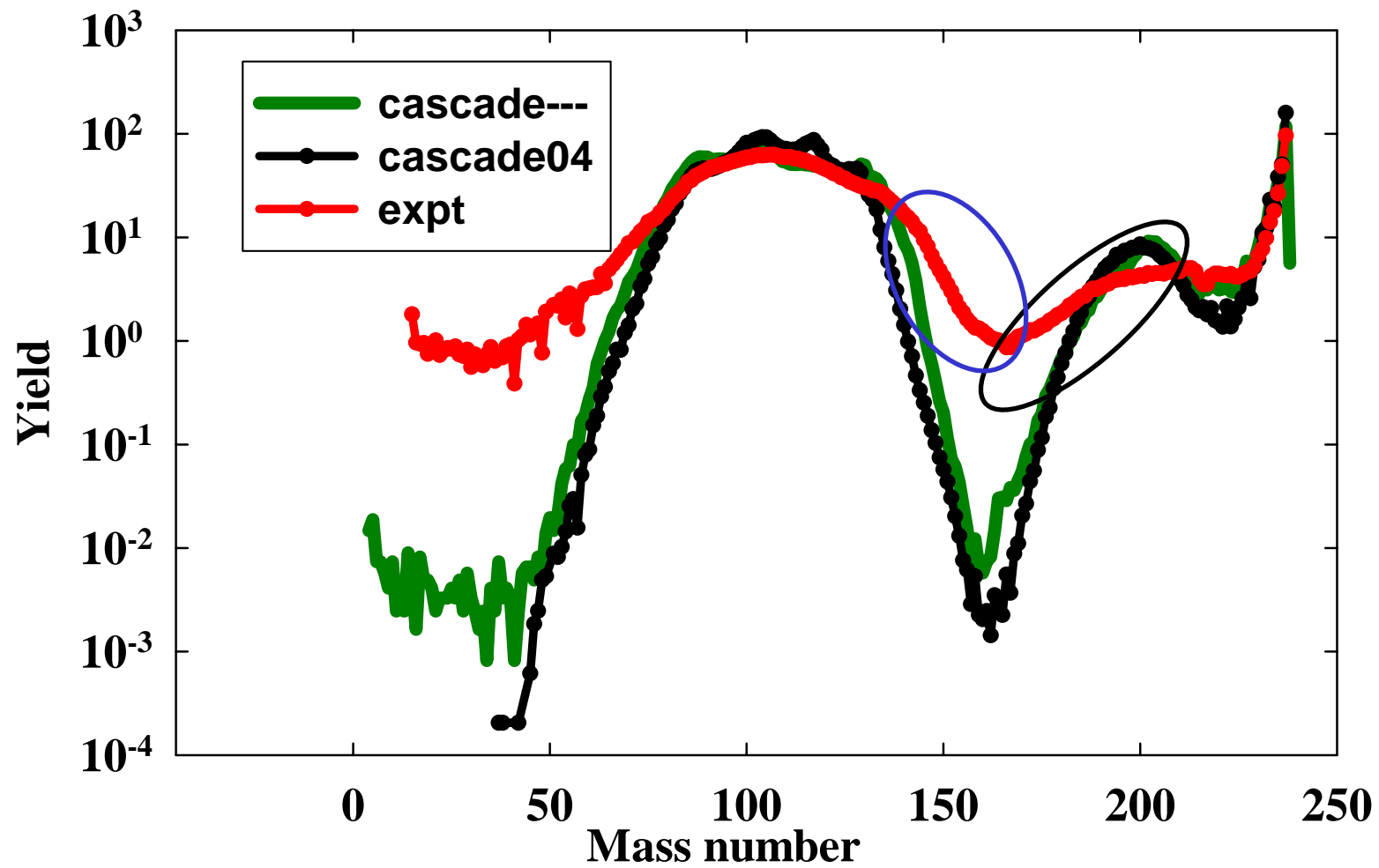


# Light fragments



# Far end of the target





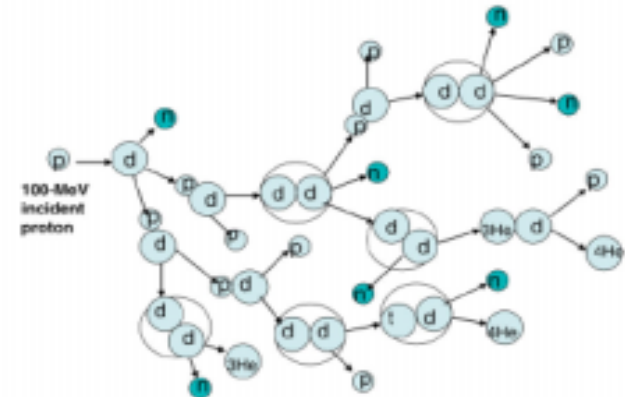


Fig. 1. A hypothetical cascade of breakup, fusion, and direct reactions initiated by a single proton. The cascade includes most of the reactions possible, although seldom will a proton induce the production of so many neutrons. To be competitive with spallation, on average two neutrons must be produced per 100-MeV proton. The reactions appear bunched together for illustrative purposes, but the proton momentum will carry most reaction particles strongly in the forward direction except for the energy-releasing fusion reactions.

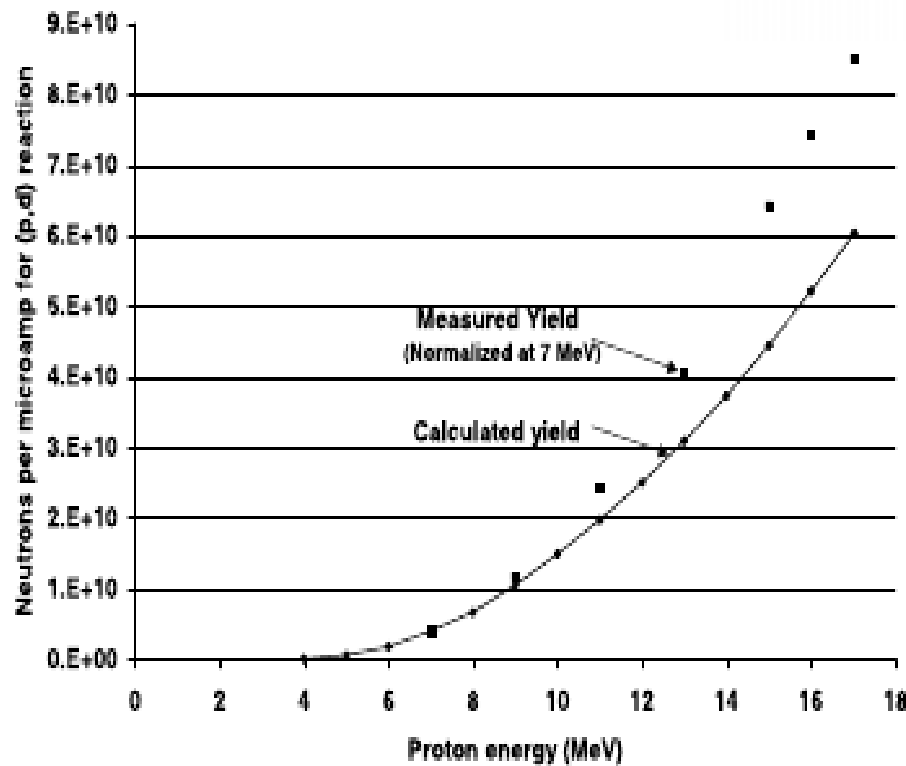


Fig. 3. Yield of neutrons from protons on a stopping-length deuterium target. The line through the lower set of points shows the expected yield based only on the  $p + d \rightarrow n + 2p$  reaction. The upper points in the 7- to 17-MeV range are the measured points. The measured yield is therefore increasing faster than from the calculated reaction alone, indicating the substantial enhancement already present from fusion and other reactions in the energy range of this experiment.



धन्यवाद

Thank you for your attention