

## Results obtained with nuclear models of Geant4 in IAEA Benchmark of Spallation

J. M. Quesada  
on behalf of the Geant4 Hadronic Group  
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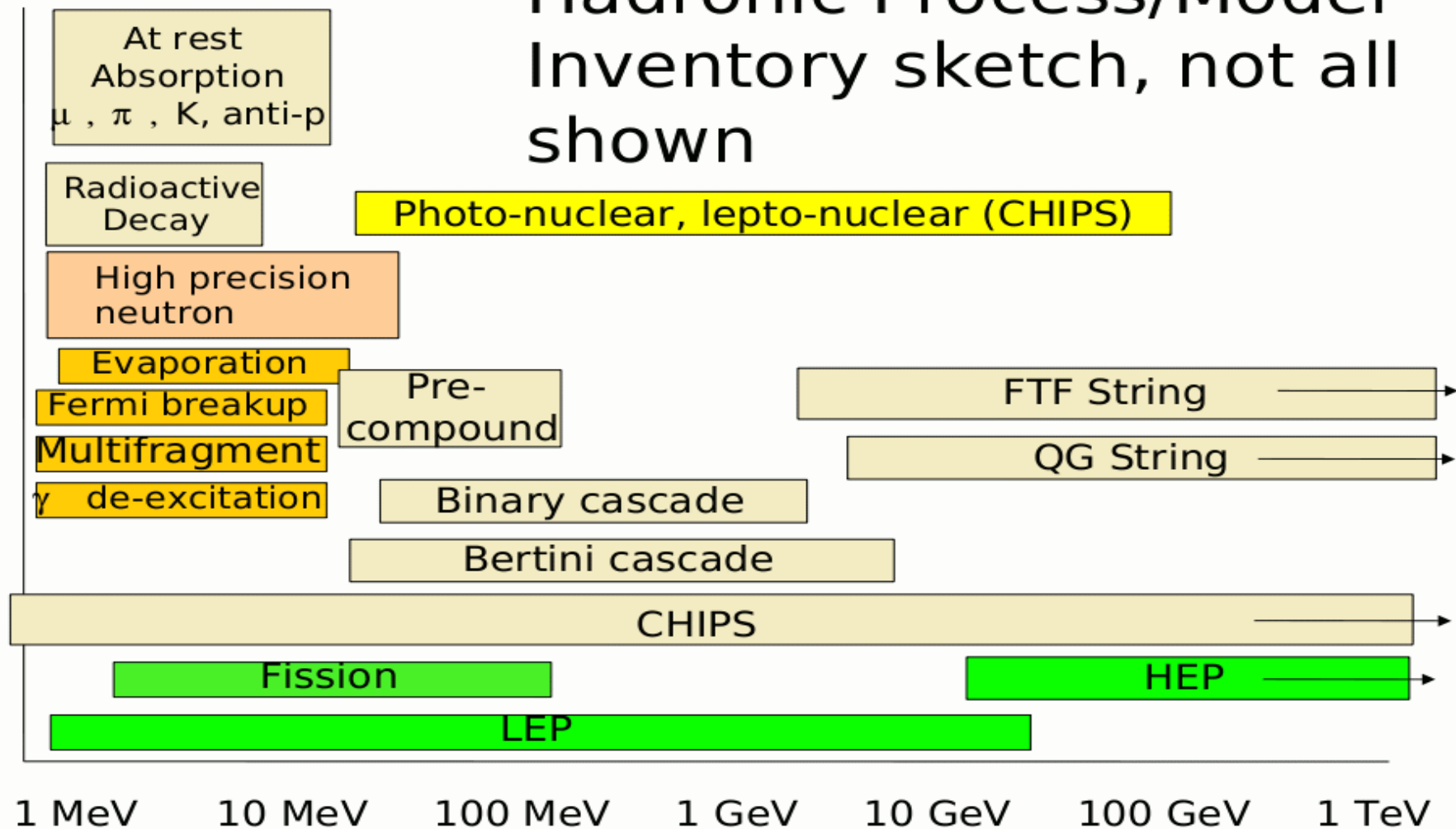
# General Introduction

- Geant4 is the C++, object-oriented successor to GEANT3
- Designed primarily with high energy physics in mind
  - but now used in medical and space applications as well
- It is a toolkit:
  - large degree of functionality and flexibility are provided
  - many different codes provided, including alternates covering the same regions of applicability
  - choice of which to use is up to user, but guidance provided by Geant4 developers
- All major physics processes covered:
  - electromagnetic, hadronic, decay, photo- and electro-nuclear

## Geant4 Hadronic Processes and Models

- Hadronic processes include
  - Elastic
  - Inelastic
  - Capture at rest
  - Neutron capture
  - Neutron-induced fission
  - Lepton-nuclear
  - Gamma-nuclear
- Each of the above processes is implemented by one or more:
  - models (which contain the physics algorithm)
  - cross sections (which determine mean free path, etc.)

## Hadronic Process/Model Inventory sketch, not all shown



- The large energy region considered in this benchmarking includes different interaction regimes.
- In order to predict the production cross sections, different reaction mechanisms must be considered
  - Cascade
  - Pre-equilibrium
  - Equilibrium de-excitation
- Bertini Cascade has its own pre-compound and de-excitation models
- Binary cascade uses the native Geant4 pre-compound and de-excitation models.

- **Binary:**
  - a time-dependent model which depends as little as possible on parameterization and therefore can be expected to be more predictive
  - is an *in house* development, including its own precompound and evaporation models.
  
- **Bertini:**
  - came from the INUCL code which was intended as an all-inclusive model.
  - It came with its own precompound and evaporation models. Neither of these are very different in origin from those in Binary, but the implementations are different.

- **CHIPS:**

- Quark-level event generator for the fragmentation of hadronic systems into hadrons.

- **QMD**

- Is a quantum extension of classical molecular-dynamics model.

- **INCL/ABLA :**

- C++ translation of INCL (v4.2) intranuclear cascade code
- C++ translation of ABLA (v3p) evaporation/fission code



- A re-engineered version of the INUCL code of N. Stepanov (ITEP)
- Employs many of the standard INC methods developed by Bertini (1968)
  - using free particle-particle collisions within cascade
  - step-like nuclear density
- Inelastic scattering of  $p$ ,  $n$ ,  $\pi$ ,  $K$ ,  $\Lambda$ ,  $\Sigma$ ,  $\Xi$
- Incident energies:  $0 < E < 10$  GeV
  - upper limit determined by lack of partial final state cross sections and the end of the cascade validity region
  - lower limit due to inclusion of internal nuclear de-excitation models

- **The Bertini model** is a classical **cascade**:
  - It is a solution to the Boltzmann equation on average
  - No scattering matrix calculated
- **Core code**:
  - Elementary particle collider: uses free cross sections
  - Up to and including 9-body final state partial cross sections for  $\pi^+p$ ,  $\pi^-p$ ,  $pp$ ,  $pn$  from the CERN compilations.
  - $K^+$ ,  $K^-$  partial cross sections.
  - $\pi^+n$ ,  $\pi^-n$ ,  $nn$  cross sections are obtained through isospin arguments
  - Partial cross sections have been reviewed and partially changed.
  - Generated secondaries:
    - pions, nucleons, kaons, hyperons.
    - No resonances
    - Deuterons, tritons,  $^3\text{He}$ , alphas (from evaporation phase only)
  - Cascade in nuclear medium
  - Final steps: pre-equilibrium and equilibrium decay of residual nucleus

- Nuclear entry point sampled over projected area of nucleus
- Incident particle is transported in density dependent nuclear medium
  - mean free path from total particle-particle cross sections
  - Nucleus modeled as 3 concentric, constant-density shells plus reflection/transmission shell boundaries.
  - nucleons have Fermi gas momentum distribution
  - Pauli exclusion invoked
- Projectile interacts with a single nucleon
  - hadron-nucleon interactions based on free cross sections and angular distributions
  - pions can be absorbed on quasi-deuterons

- Each secondary from initial interaction is propagated in nuclear potential until it interacts or leaves nucleus
  - can have reflection from density shell boundaries
  - Coulomb barrier added for all phases
- As cascade collisions occur, exciton states are built up, leading to equilibrated nucleus
  - selection rules for p-h state formation:  $\Delta p = 0, +/-1,$   
 $\Delta h = 0, +/-1, \quad \Delta n = 0, +/-2$
- Model uses its own exciton routine based on that of Griffin
  - Kalbach matrix elements used
  - level densities parametrized vs. Z and A

- Cascade ends and exciton model takes over when secondary KE drops below 20% of its original value or  $7 \times$  nuclear binding energy
- Nuclear evaporation follows for most nuclei
  - emission continues as long as excitation is large enough to remove a particle.
- For light, highly excited nuclei, Fermi breakup
- Fission included in phenomenological way for heavy elements

- H.P. Wellisch and G. Folger (CERN)
- Henning Weber (Frankfurt group)
- Based in part on Amelin's kinetic model
- Incident p, n
  - $0 < E < 3 \text{ GeV}$
  
- light ions
  - $0 < E < 3 \text{ GeV/A}$
  
- $\pi$ 
  - $0 < E < 1.5 \text{ GeV}$

- Hybrid between classical cascade and full QMD model
- Detailed model of nucleus
  - nucleons placed in space according to nuclear density
  - nucleon momentum according to Fermi gas model
- Nucleon momentum is taken into account when evaluating cross sections, i.e. collision probability
- Collective effect of nucleus on participant nucleons described by optical potential
  - Participant particle's equations of motion are integrated numerically.

- Nucleon-nucleon scattering (t-channel) resonance excitation cross-sections are derived from p-p scattering using isospin invariance, and the corresponding Clebsch-Gordan coefficients
  - elastic N-N scattering included
- Meson-nucleon inelastic (except true absorption) scattering modelled as s-channel resonance excitation. Breit-Wigner form used for cross section.
- Resonances may interact or decay
  - nominal PDG branching ratios used for resonance decay
  - masses sampled from Breit-Wigner form
- Developed in collaboration with Frankfurt group, broadly similar to UrQMD



- Calculate imaginary part of the R-matrix using free 2-body cross-sections from experimental data and parameterizations
- For resonance re-scattering, the solution of an in-medium BUU equation is used.
  - The Binary Cascade at present takes the following strong resonances into account:
    - The delta resonances with masses 1232, 1600, 1620, 1700, 1900, 1905, 1910, 1920, 1930, and 1950 MeV
    - Excited nucleons with masses 1440, 1520, 1535, 1650, 1675, 1680, 1700, 1710, 1720, 1900, 1990, 2090, 2190, 2220, and 2250 MeV

- Nucleon-nucleon elastic scattering angular distributions taken from :
  - Arndt phase shift analysis of experimental data up to 1.2 GeV
  - at higher energies PDG data are used
- Pauli blocking implemented in its classical form
  - final state nucleons occupy only states above Fermi momentum
- Coulomb barrier taken into account for charged hadrons

## Binary Cascade Modelling Sequence(4)

- If primary below 45 MeV, no cascade, just precompound
- Cascade stops when mean energy of all scattered particles is below  $0.2 \cdot A$ -dependent cut for the average kinetic energies of secondaries, which means :
  - 18 MeV for  $A < 31$
  - 14 MeV for  $A < 61$
  - 10 MeV for  $A < 121$
  - 9 MeV for  $A > 120$
- When cascade stops, the properties of the residual exciton system and nucleus are evaluated, and passed to pre-equilibrium de-excitation class
- Pre-equilibrium: exciton model

- Native pre-equilibrium de-excitation model in Geant4 is a version of standard exciton model.

### Key ingredients:

- Internal transition rates:
  - CEM (Cascade Exciton Model, Gudima et al). **Default**
  - Blann-Machner's parameterization.
- Emission rates:
  - Nucleon emission in standard exciton formulation.
  - Complex particle emission (d,t,<sup>3</sup>He, <sup>4</sup>He) from CEM.

- The transition rates (for  $\Delta n = -2, 0, +2$ ):

$$\lambda_{\Delta n}(p, h, E^*) = \frac{2\pi}{\hbar} |M_{\Delta n}|^2 \omega_{\Delta n}(p, h, E^*)$$

- The **total** transition rate:

$$\lambda_{total}(p, h, E^*) = \sum_{\Delta n=-2,0,+2} \lambda_{\Delta n}(p, h, E^*)$$

- The “j” particle (nucleon) emission probability distribution:

$$\lambda_c^j(p, h, E^*, \epsilon) = \frac{2s_j + 1}{\pi^2 \hbar^3} \mu_j \mathcal{R}_j(p, h) \frac{\omega(p-1, h, E^* - B_j - \epsilon)}{\omega(p, h, E^*)} \epsilon \sigma_{inv}(\epsilon)$$

- The “j” particle (cluster) emission probability distribution:

$$\lambda_c^j(p, h, E^*, \epsilon) = \frac{2s_j + 1}{\pi^2 \hbar^3} \mu_j \mathcal{R}_j(p, h) \gamma_j \frac{\omega(p-p_j, h, E^* - B_j - \epsilon)}{\omega(p, h, E^*)} \frac{\omega(p_j, 0, B_j + \epsilon)}{g_j} \epsilon \sigma_{inv}(\epsilon)$$

- The **total** emission rate:

$$\Gamma_c(p, h, E^*) = \sum_i \int_{V_i^c}^{E^* - B_j} \lambda_c^j(p, h, E^*, \epsilon) d\epsilon$$

**Inverse reaction cross sections play a mayor role in the calculation of (competing) emission probabilities.**

- Theory driven *old* parameterization (Dostrovski et al, 1959) (kept as option)

### New parameterization (after release 9.2)

- More realistic parameterization of reaction cross sections:
  - Chatterjee et al: Calculated with global optical model potentials, in turn fitted to reproduce available experimental data
  - Kalbach's retuning (PRECO code)
  - Wellisch's parameterization of proton reaction cross sections by direct fitting to experimental data
  - **Default option** combines the best combination of inverse cross sections (Wellisch's parameterization for protons and Kalbach's one for the rest)

- En principle, the transition from pre-equilibrium to equilibrium de-excitation should take place when the following condition is fulfilled:
 
$$\lambda_+(p, h, E) = \lambda_-(p, h, E)$$

- Which can be roughly estimated as:

$$n_{eq} = \sqrt{2gE^*}$$

- The practical need of less pre-equilibrium emission led to the introduction following probability :

$$P_{pe}(n/n_{eq}) = 1 - e^{-\frac{1}{2\sigma_{pre}^2} \left(\frac{n}{n_{eq}} - 1\right)^2}$$

for  $n < n_{eq}$  and equal to zero for  $n > n_{eq}$ , with  $\sigma_{pre} \approx 0.4$ .

After pre-equilibrium the properties of the residual nucleus are evaluated, and passed to the equilibrium de-excitation handler.

Four processes are considered:

- Fermi break-up , for  $Z < 9$ ,  $A < 17$  (Botvina *et al*)
- Statistical multifragmentation, for  $E^*/A > 3$  MeV (Botvina *et al*)

### Competitors:

- Fission (Bohr-Wheeler model + Amelin prescript.)
- Particle evaporation:
  - Evaporation model WE (Weisskopf-Ewing)
  - Generalized Evaporation Model GEM (Furihata).



## Evaporation models

- WE: evaporation of particles (n,p,d,t,<sup>3</sup>He,α ) from a completely degenerated Fermi gas (excited compound nucleus)

$$\Gamma_c^j(E^*, \epsilon) = \frac{2s_j + 1}{\pi^2 \hbar^3} \mu_j \frac{\exp[2\sqrt{a_0(E^* - S_n - \delta_0 - \epsilon)}]}{\exp[2\sqrt{a_0(E^* - \delta_0)}]} \epsilon \sigma_{inv}(\epsilon)$$

improved  
↓

- GEM: generalization for including heavier ejected fragments (Z>13 , A<29)
- Combination of *improved* WE for evaporation of n,p,d,t,<sup>3</sup>He,α and GEM for heavier fragments. **DEFAULT**
- Break down in explosive scenarios:
  - Very light systems → Fermi break-up
  - High excitation energies → Statistical Multifragmentation

- For very light systems even a relatively small  $E^*$  may be comparable with their total binding energy.
- Fermi model does the partitioning of excited light nuclei ( $Z < 9$ ,  $A < 17$ ) and the sampling of their kinetic energies, based on very basic physical grounds (phase space availability).
- It does not have any impact in present benchmark.
- It is relevant in reactions involving light compound nuclei (hadrontherapy)

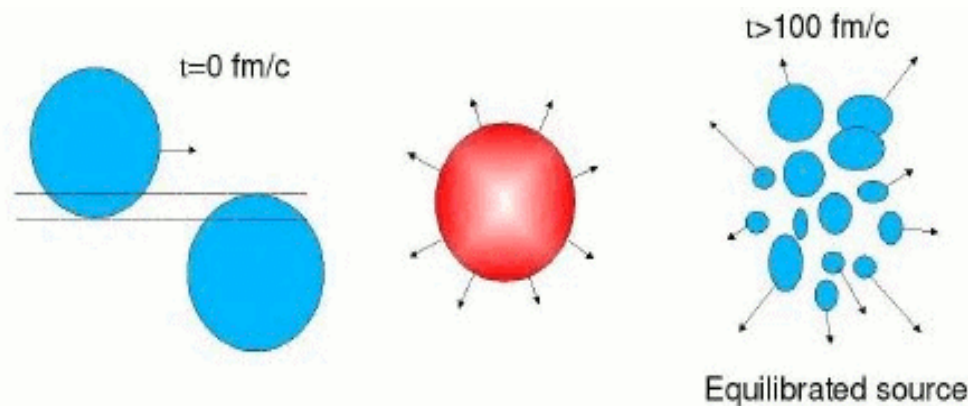
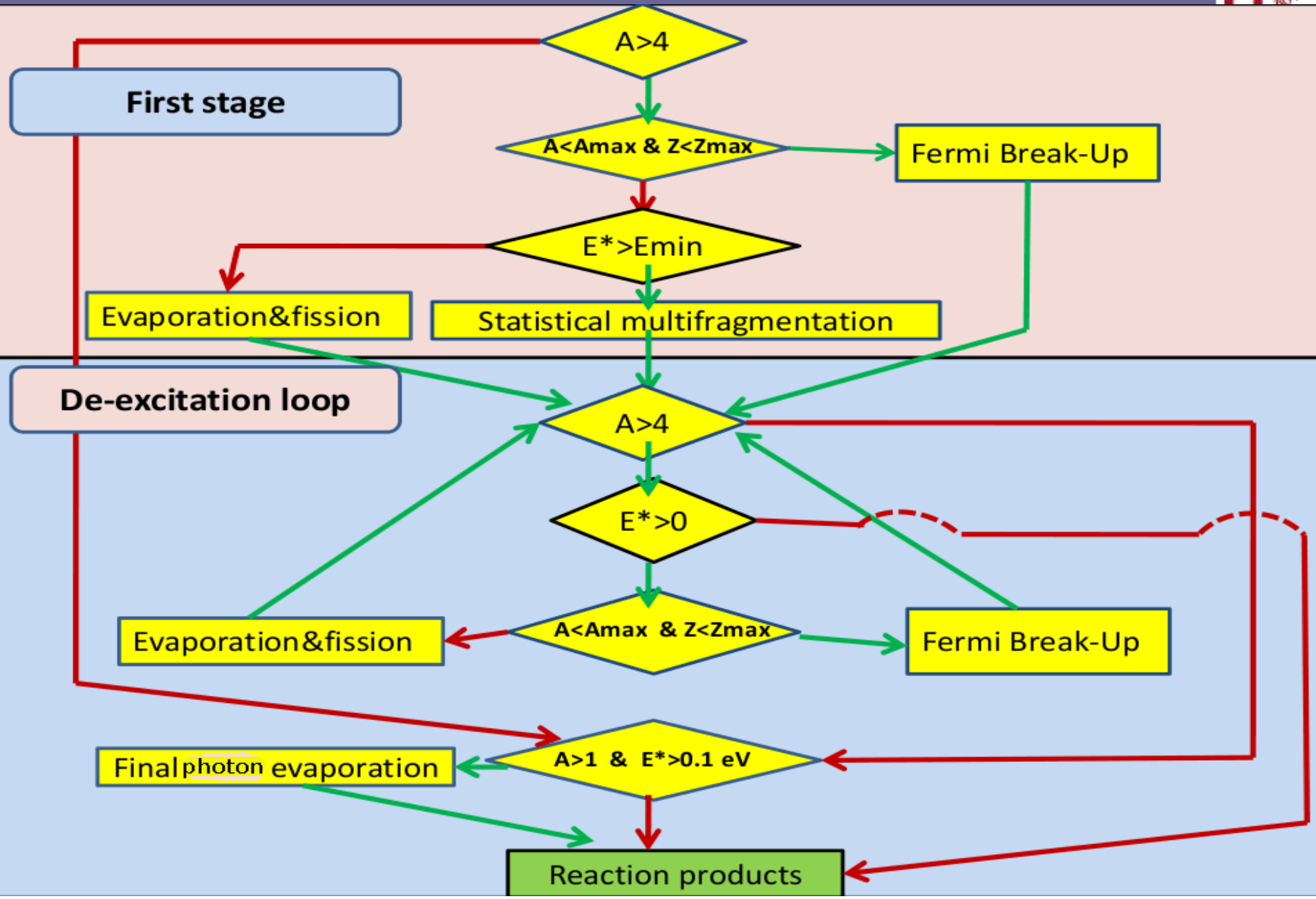


Figure: For  $E^*/A \geq 3$  MeV

- Implementation of Statistical Multifragmentation Model, developed since the seventies by Mishustin-Botvina.
- Igor Psenishnov (FIAS) has validated it against original FORTRAN code.



(related to geant4.9.2p01 official release results )

- No *ad hoc* tuning of level density parameter ratio  $a_{\text{fis}}/a_{\text{evap}}$ . (preliminary trials show that it is critical, as reported in previous works).
- No *soft transition* from pre-equilibrium (i.e. increment of equilibrium at the expenses of pre-equilibrium) .
- **Very important:** parameters tuned in a “model suite” shouldn’t be assumed to work in a different *environment*, i.e. with different *coupled* models.



Ad hoc tuning of parameters was clearly necessary in order to reproduce fission data. (Done in next release)

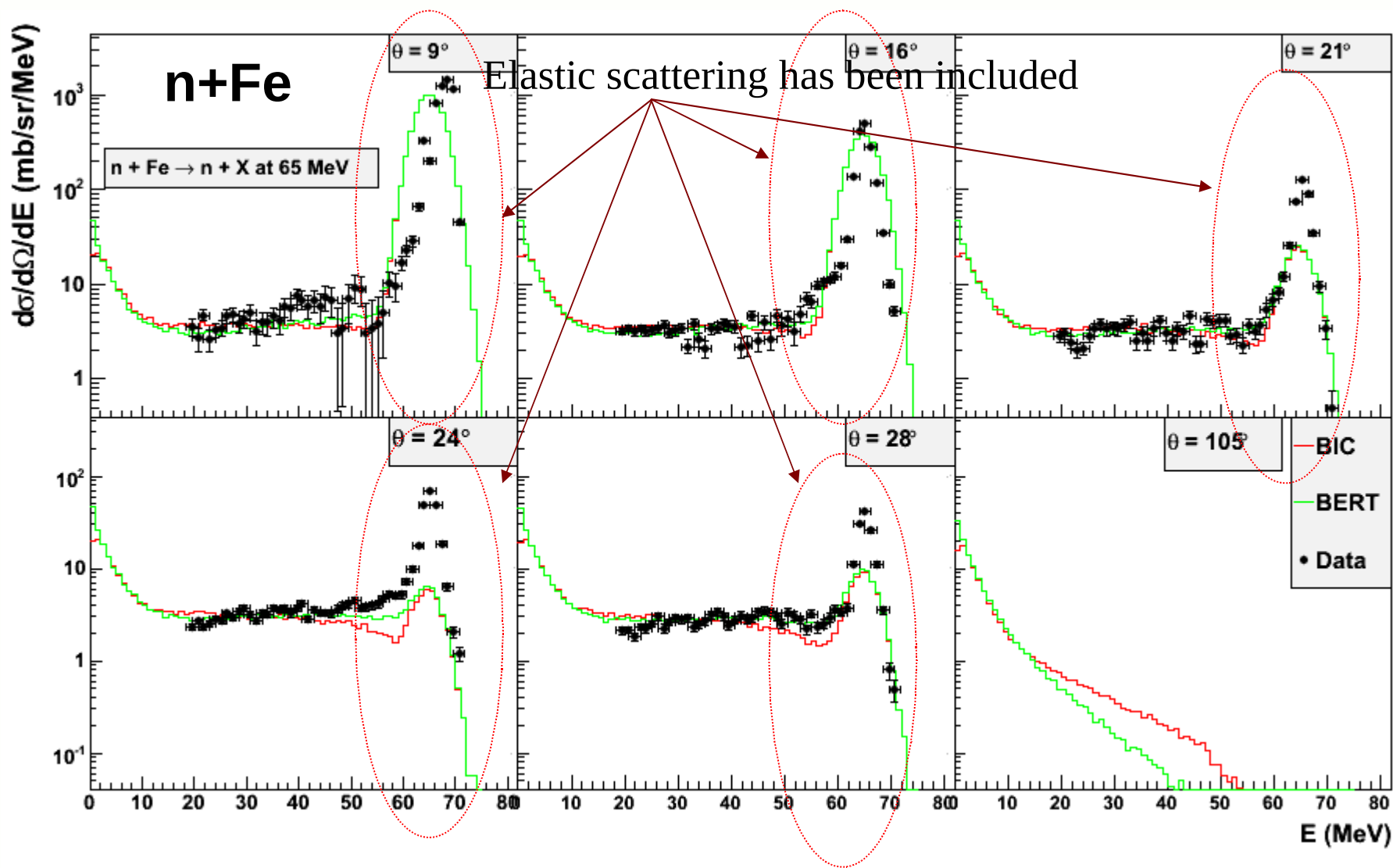
(included in geant4.9.3 official release )

- Transition probabilities at pre-equilibrium (exciton model) have been calculated according to CEM version of model
- Combined WE-GEM model has been implemented in de-excitation (allows description of IMF production)
- First retuning of parameters now undertaken using two example reactions
  - Tuning of level density parameter ratio  $a_{\text{fis}}/a_{\text{evap}}$ .
  - Tuning of the width of symmetric component of fission fragment distribution

## RESULTS

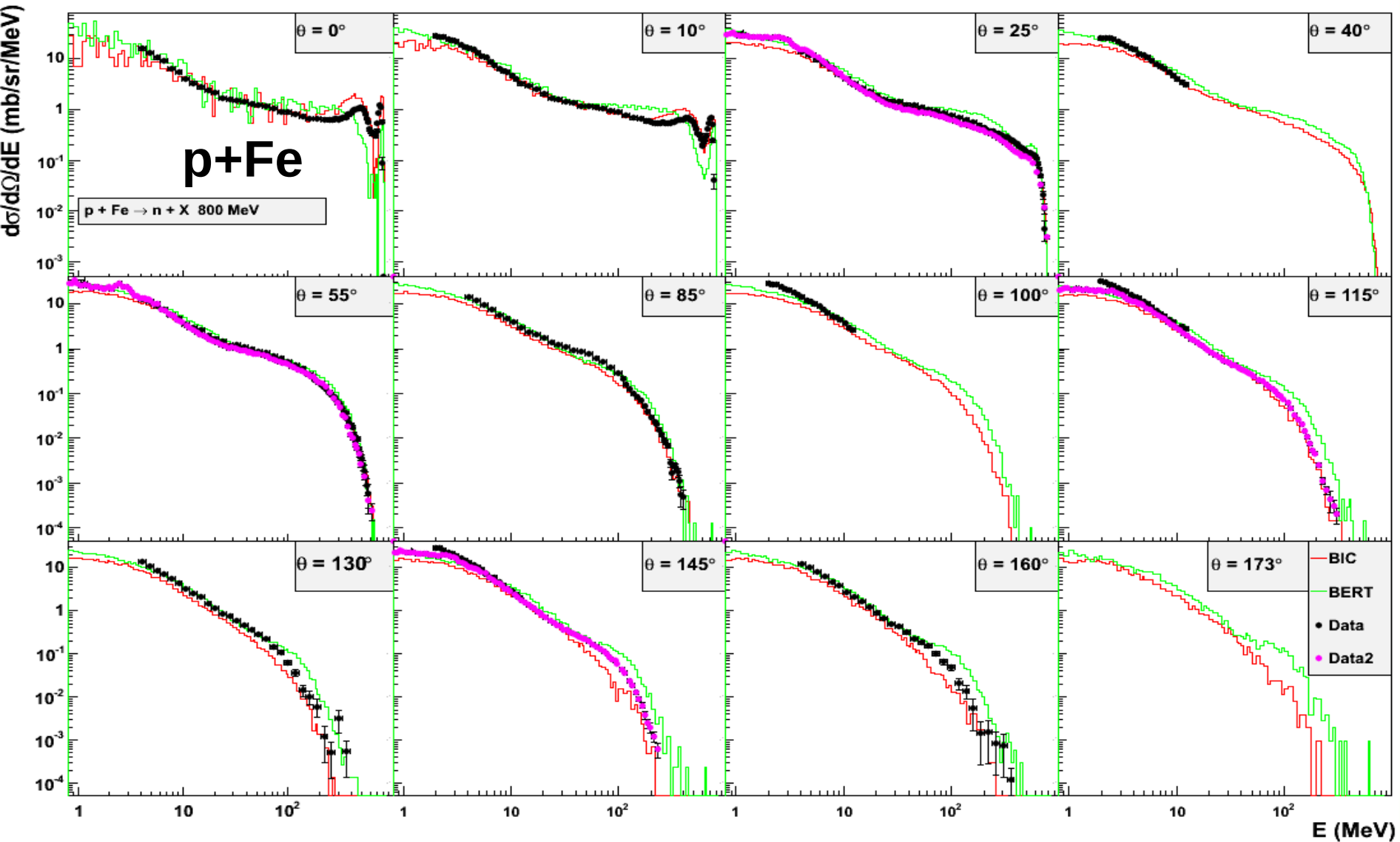
(Geant4 release 9.3)

## Neutron production at 63 MeV

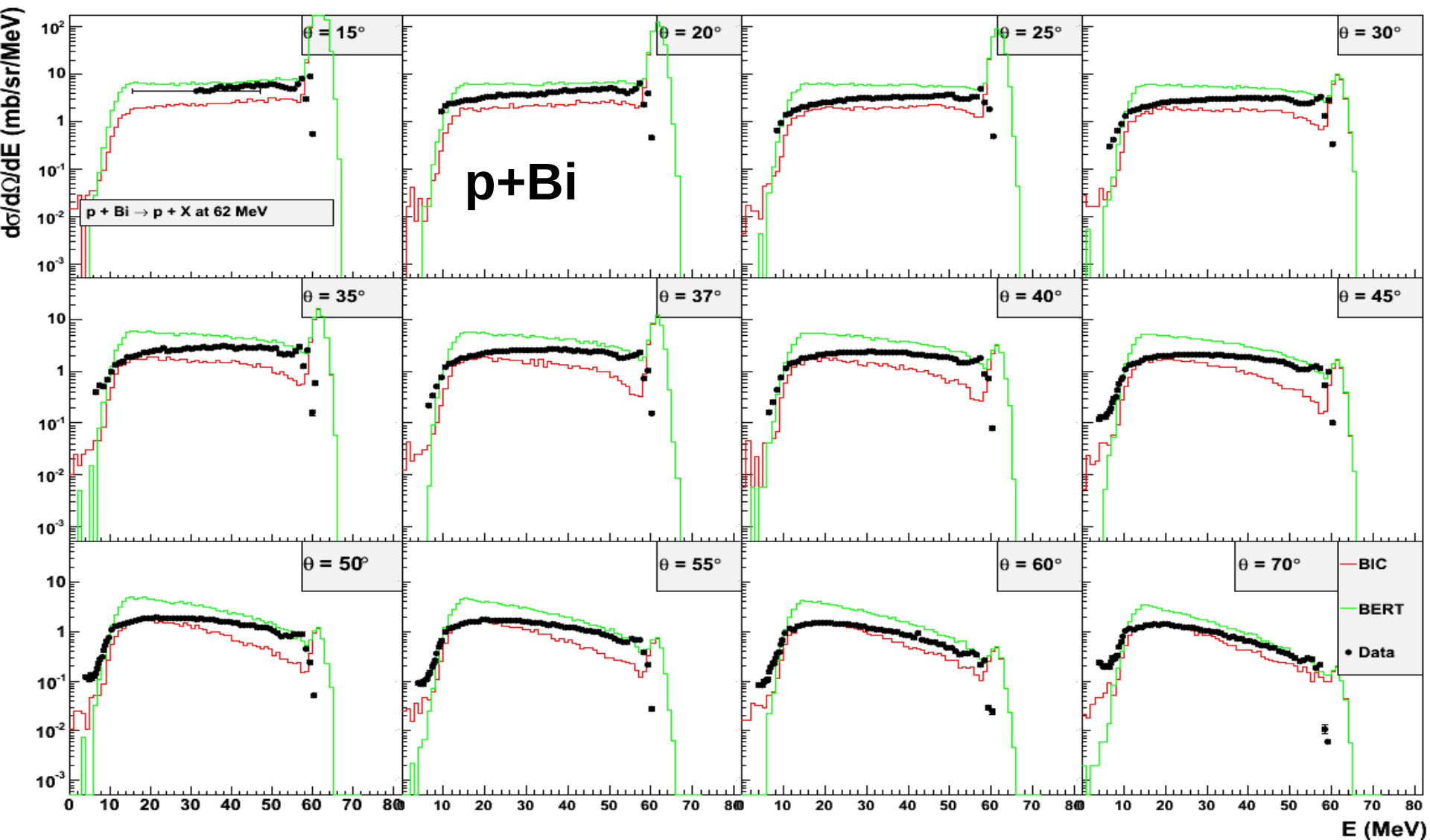




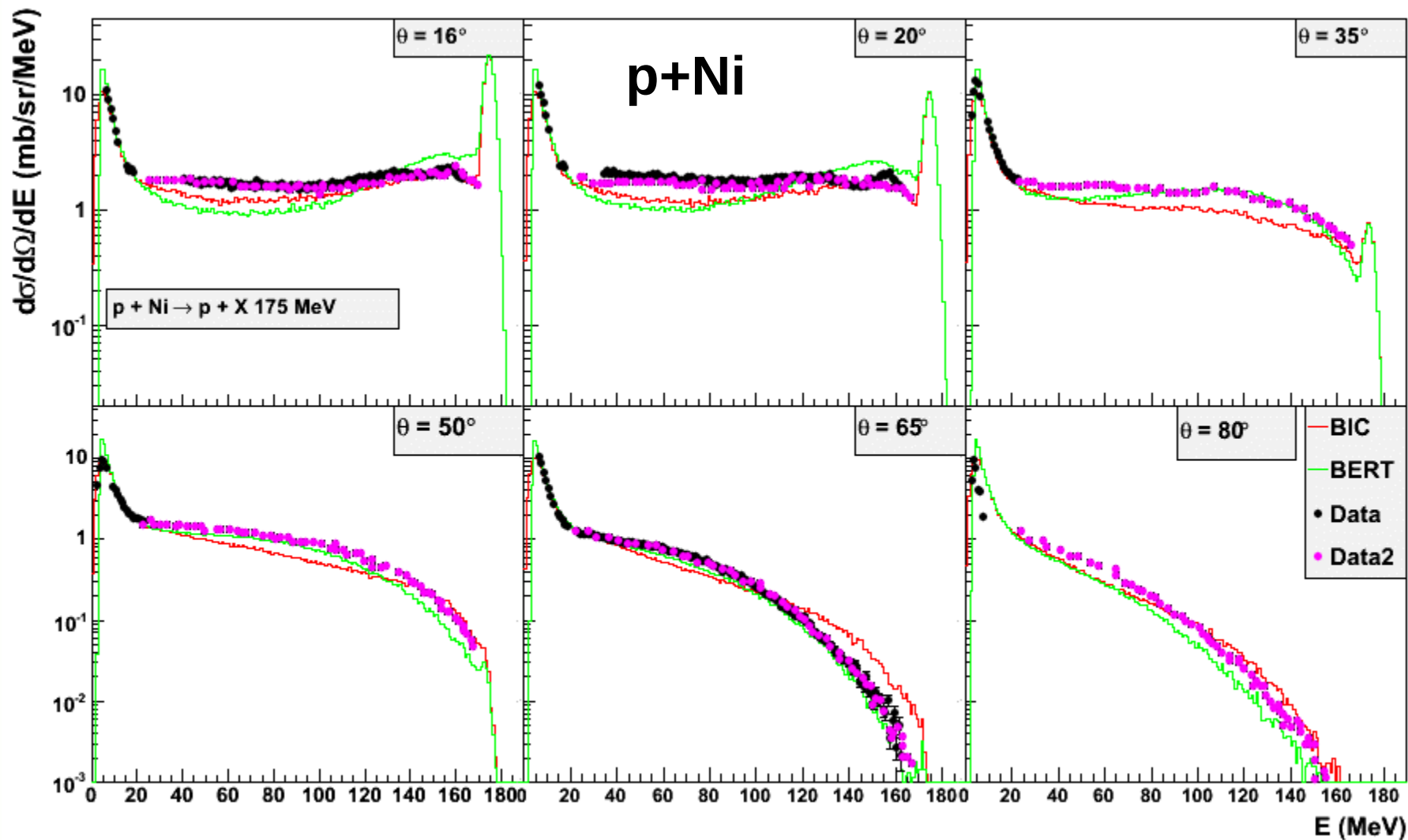
## Neutron production at 1200 MeV

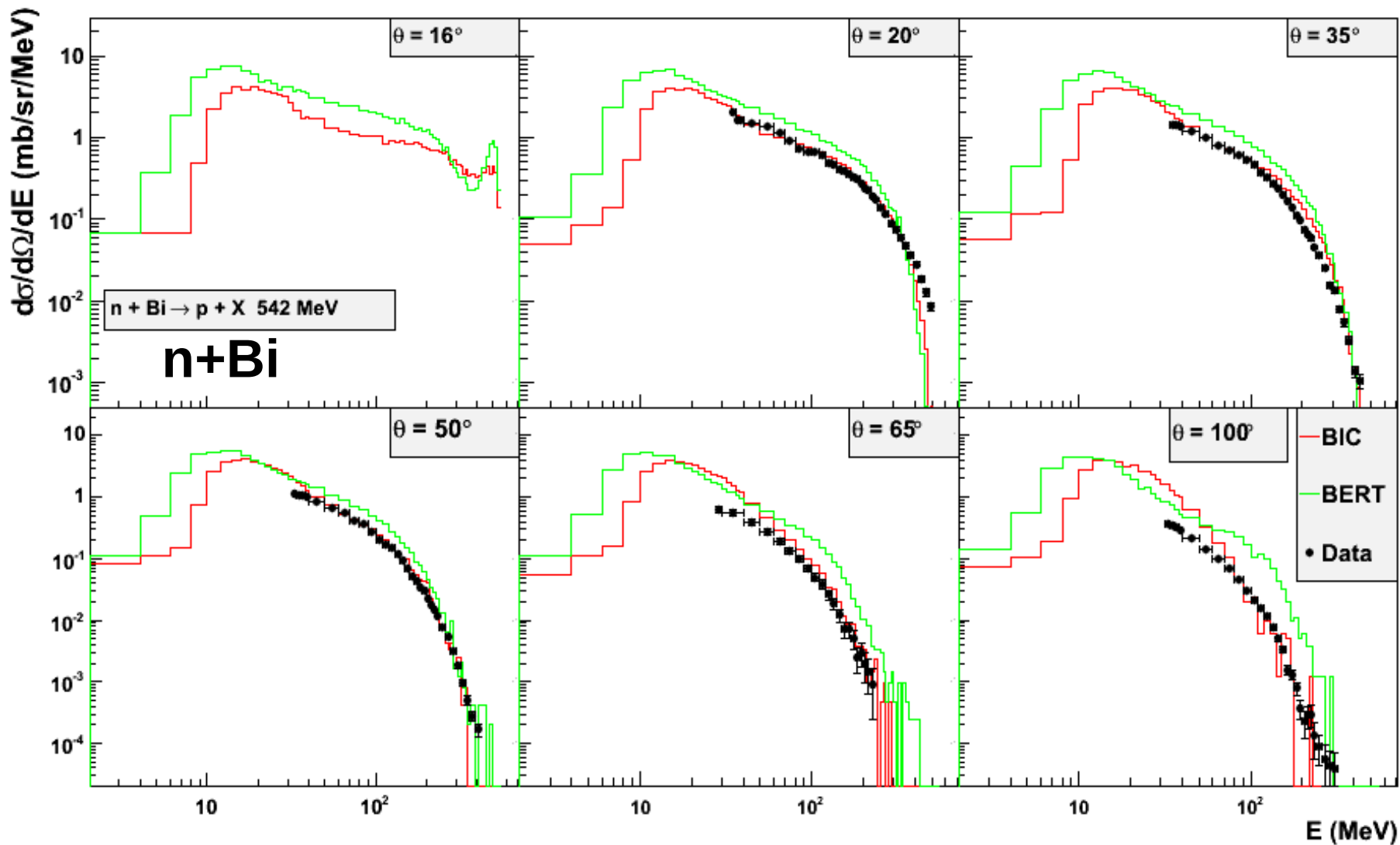


## Proton production at 62 MeV

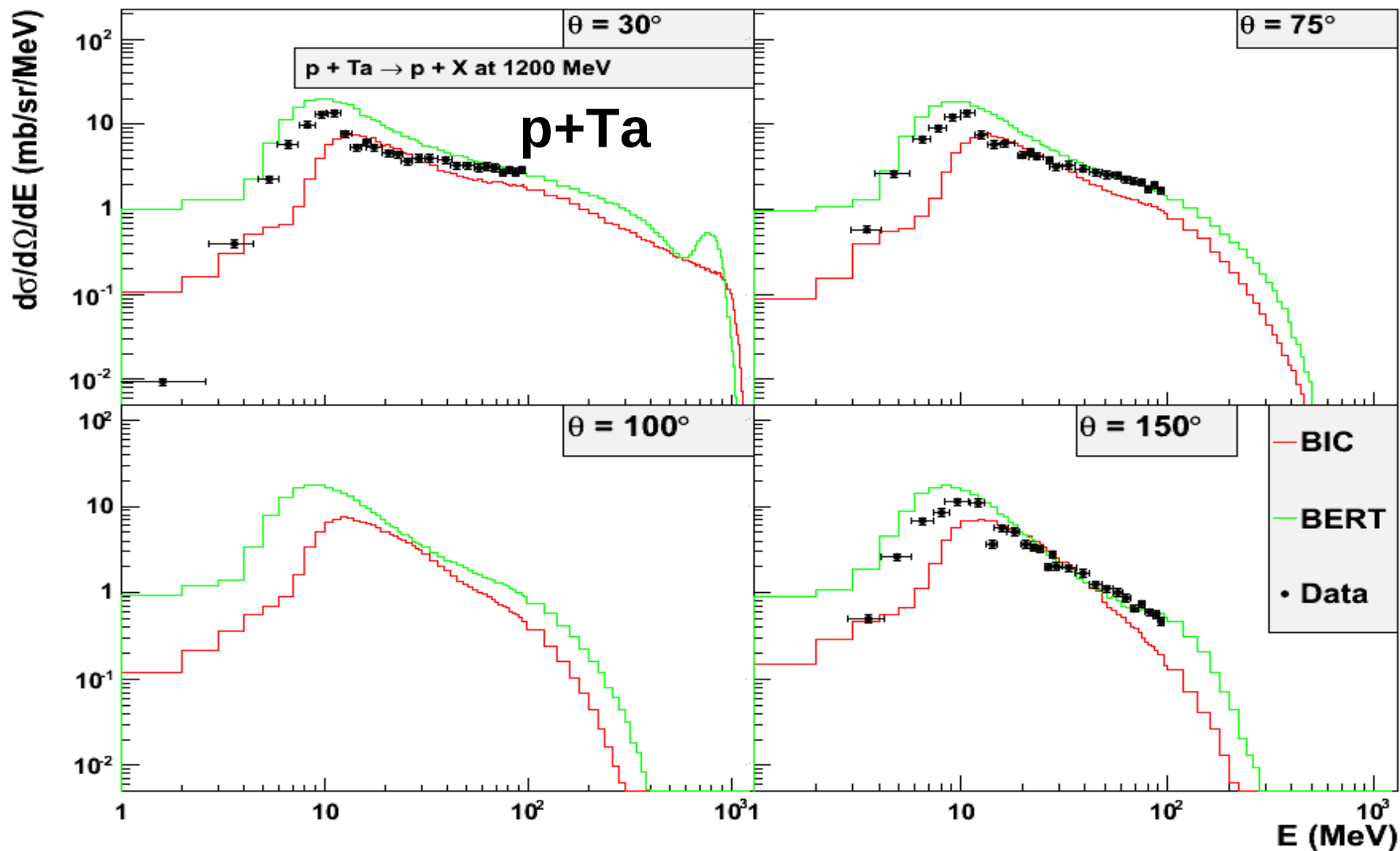


## Proton production at 175 MeV

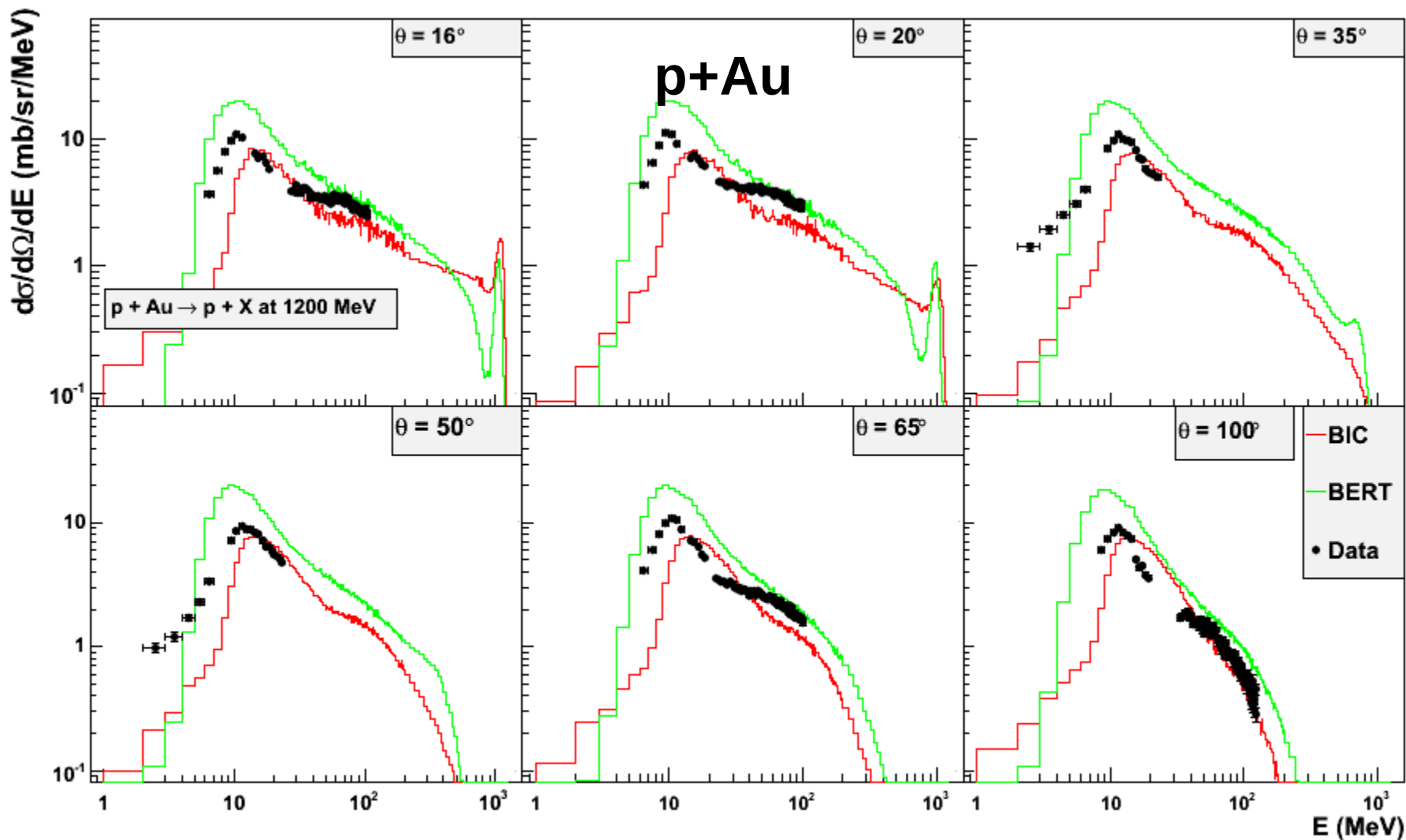


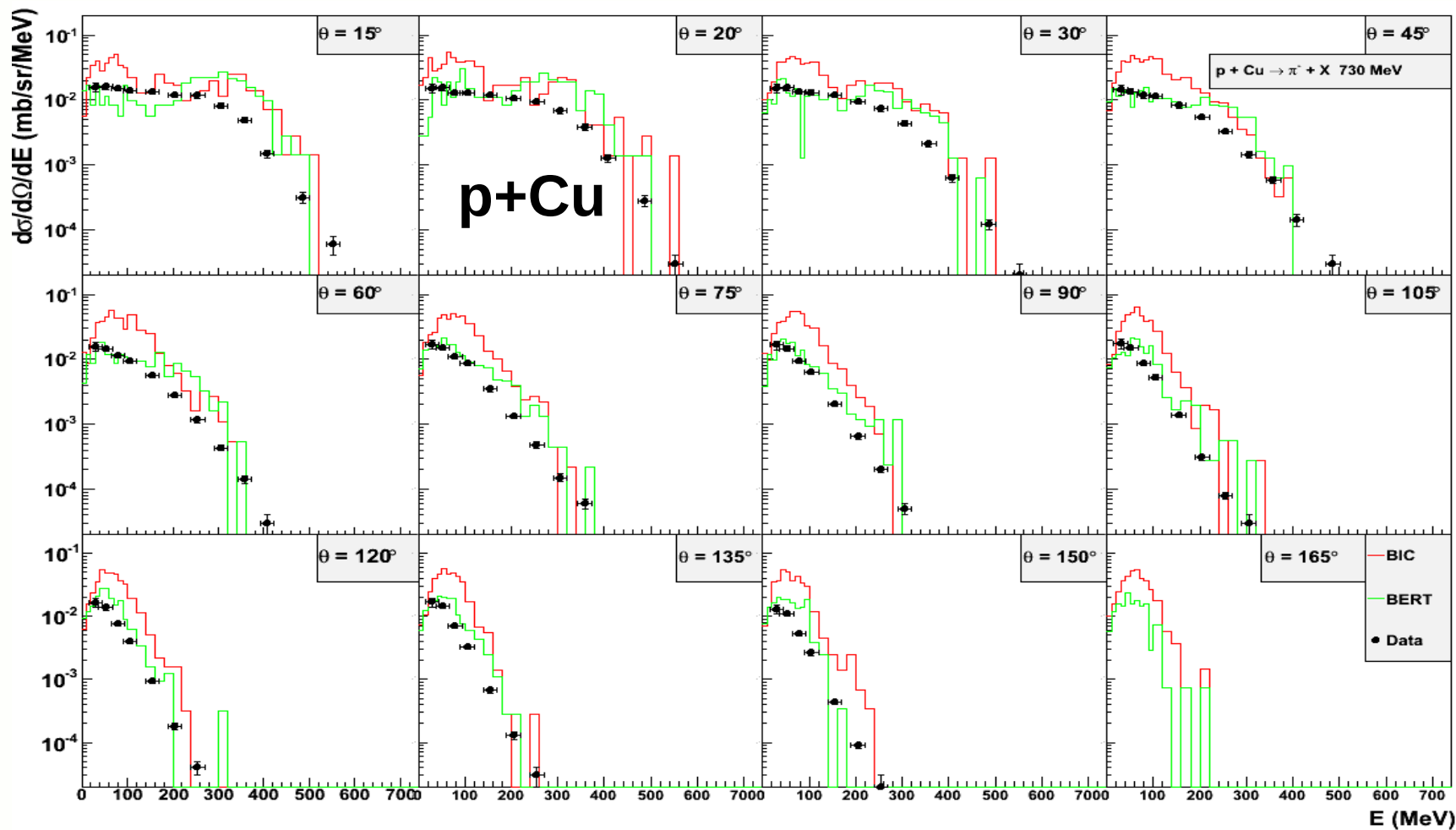


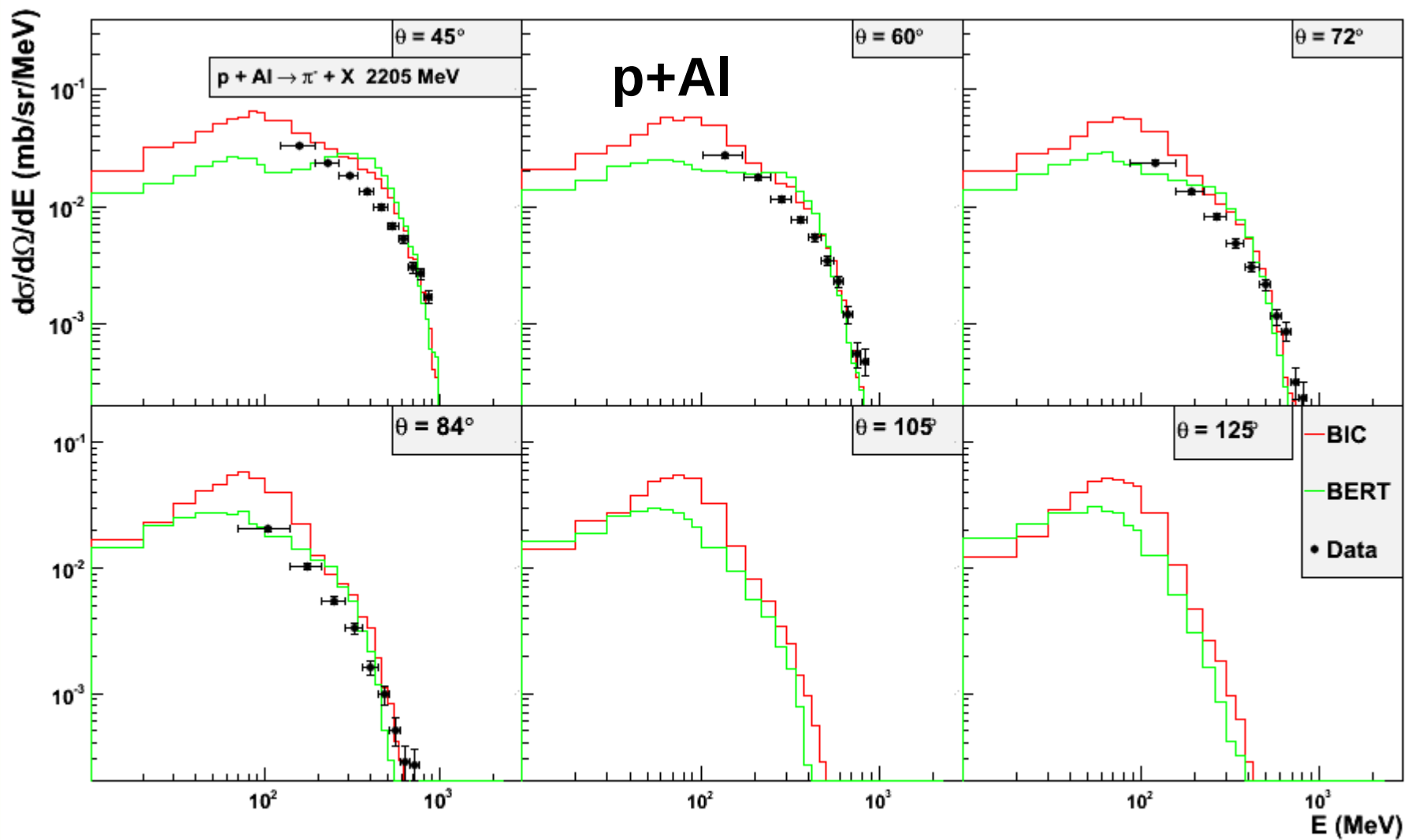
## Proton production at 1200 MeV



## Proton production at 1200 MeV

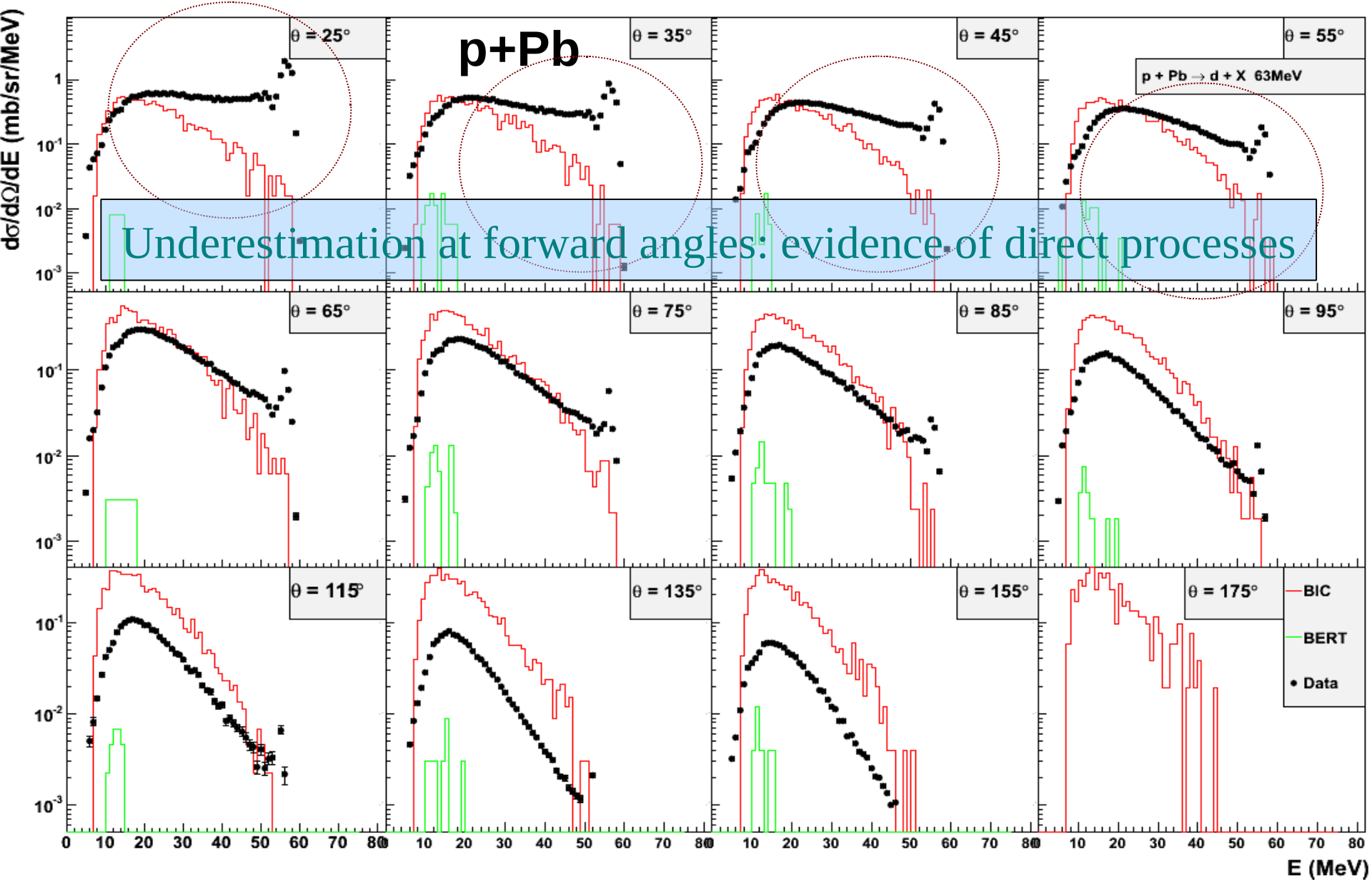


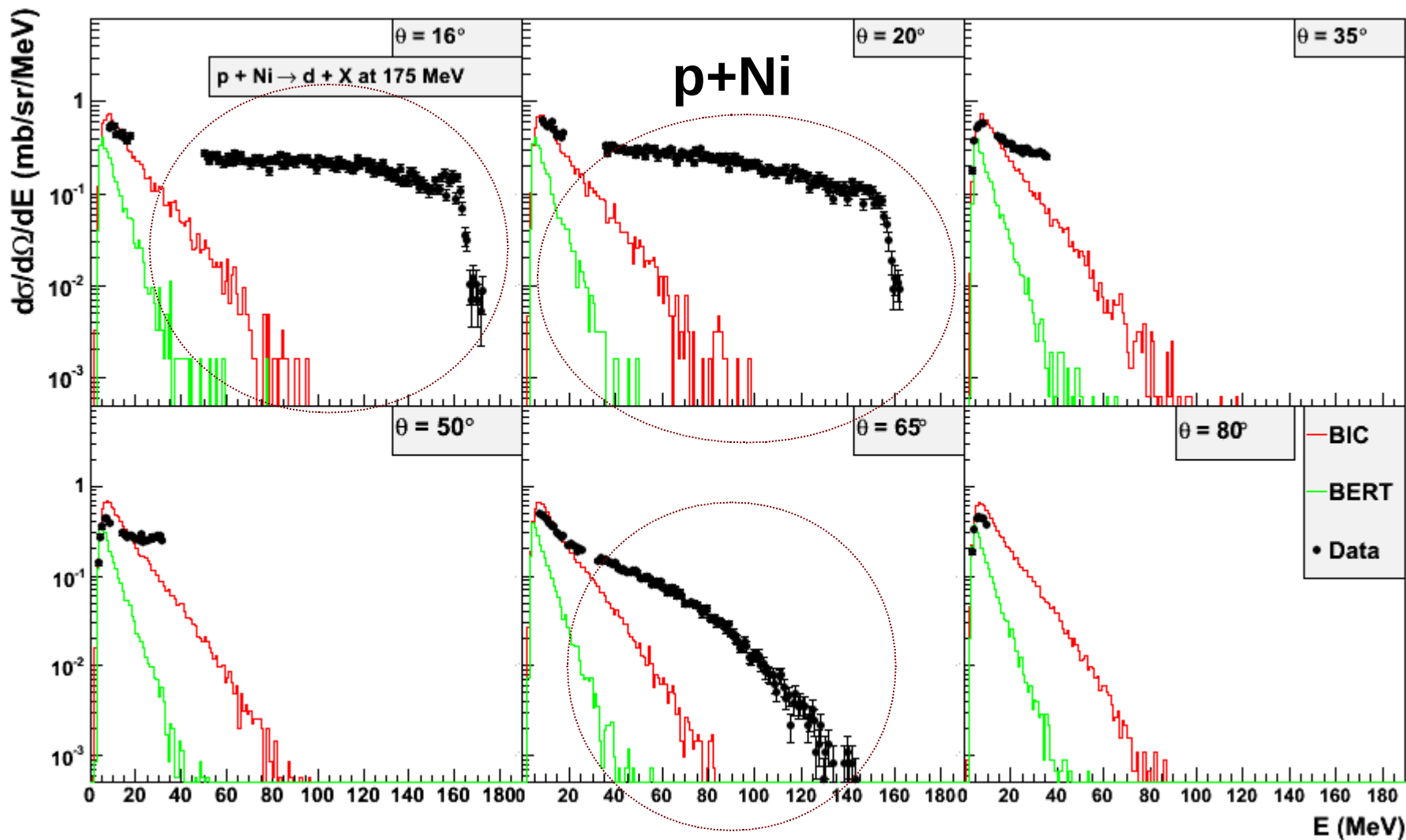


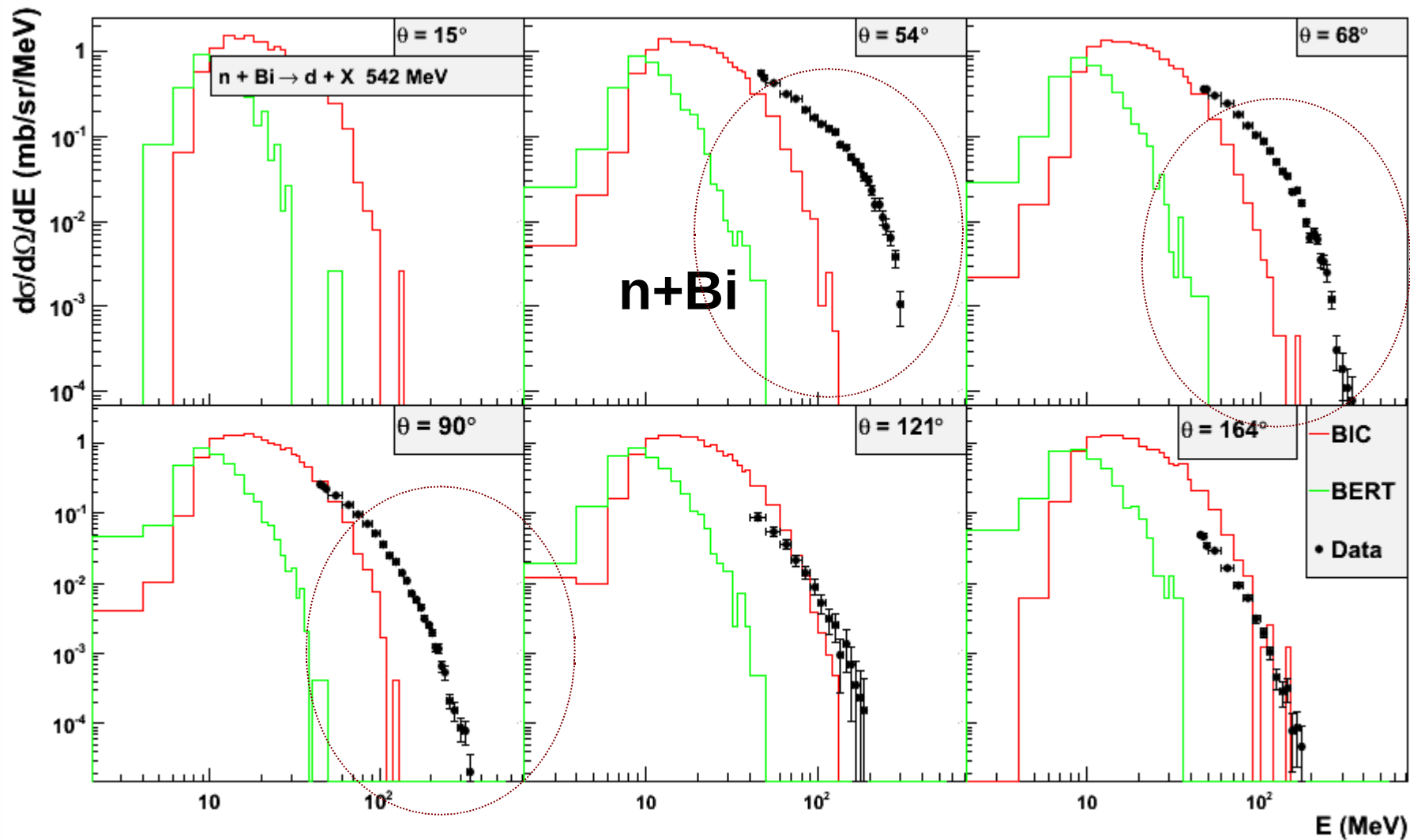


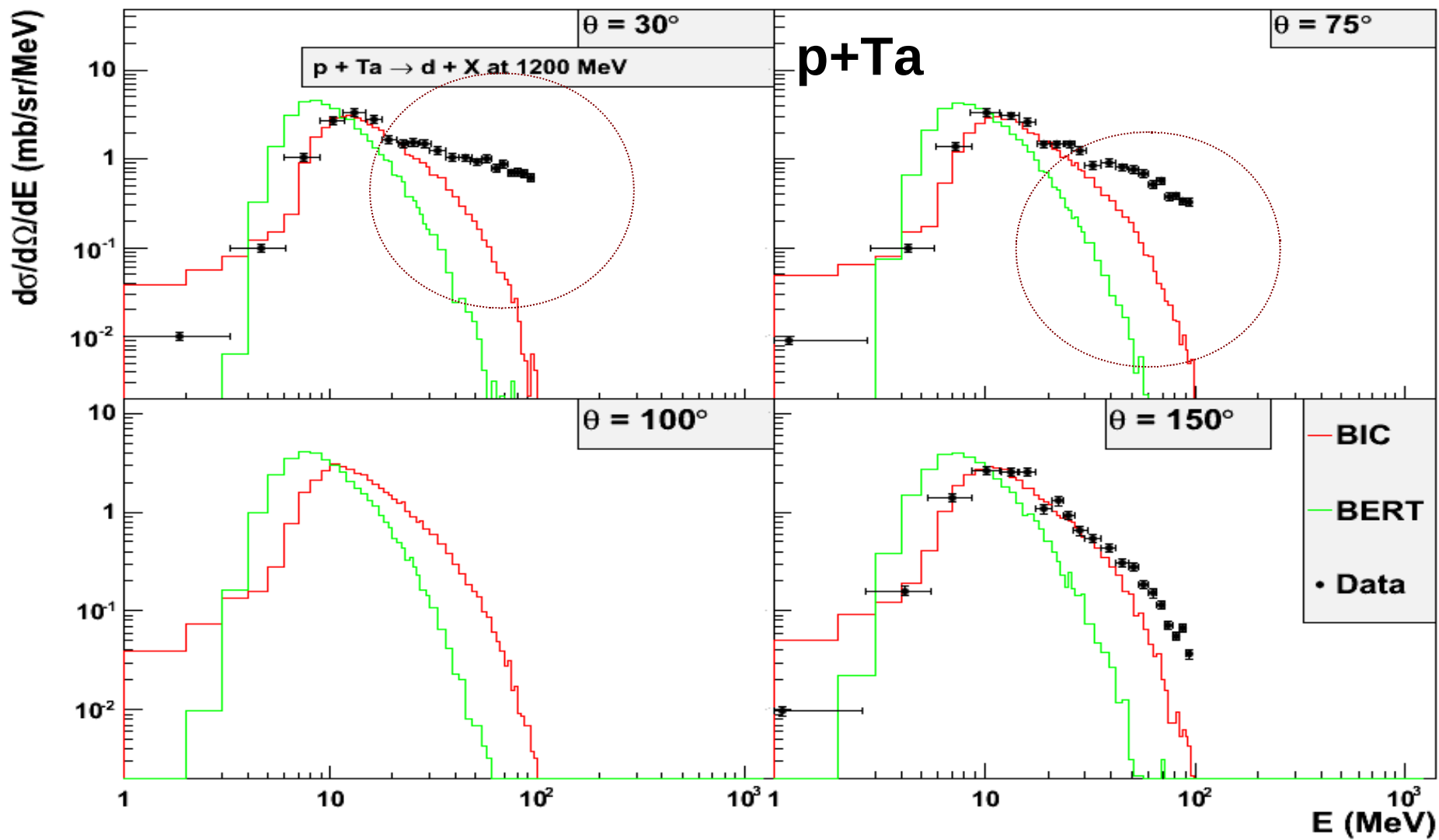


## Deuteron production at 63 MeV

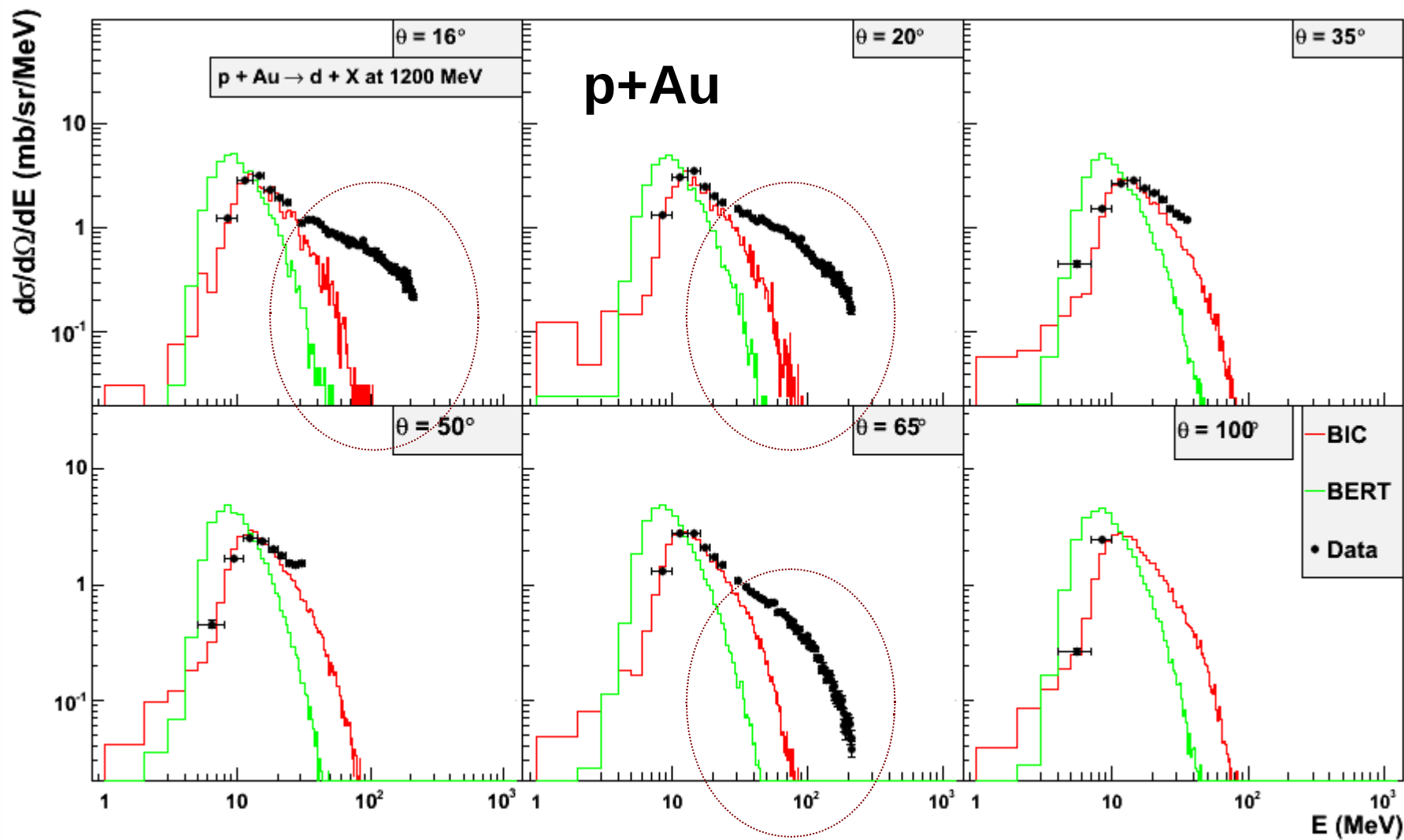




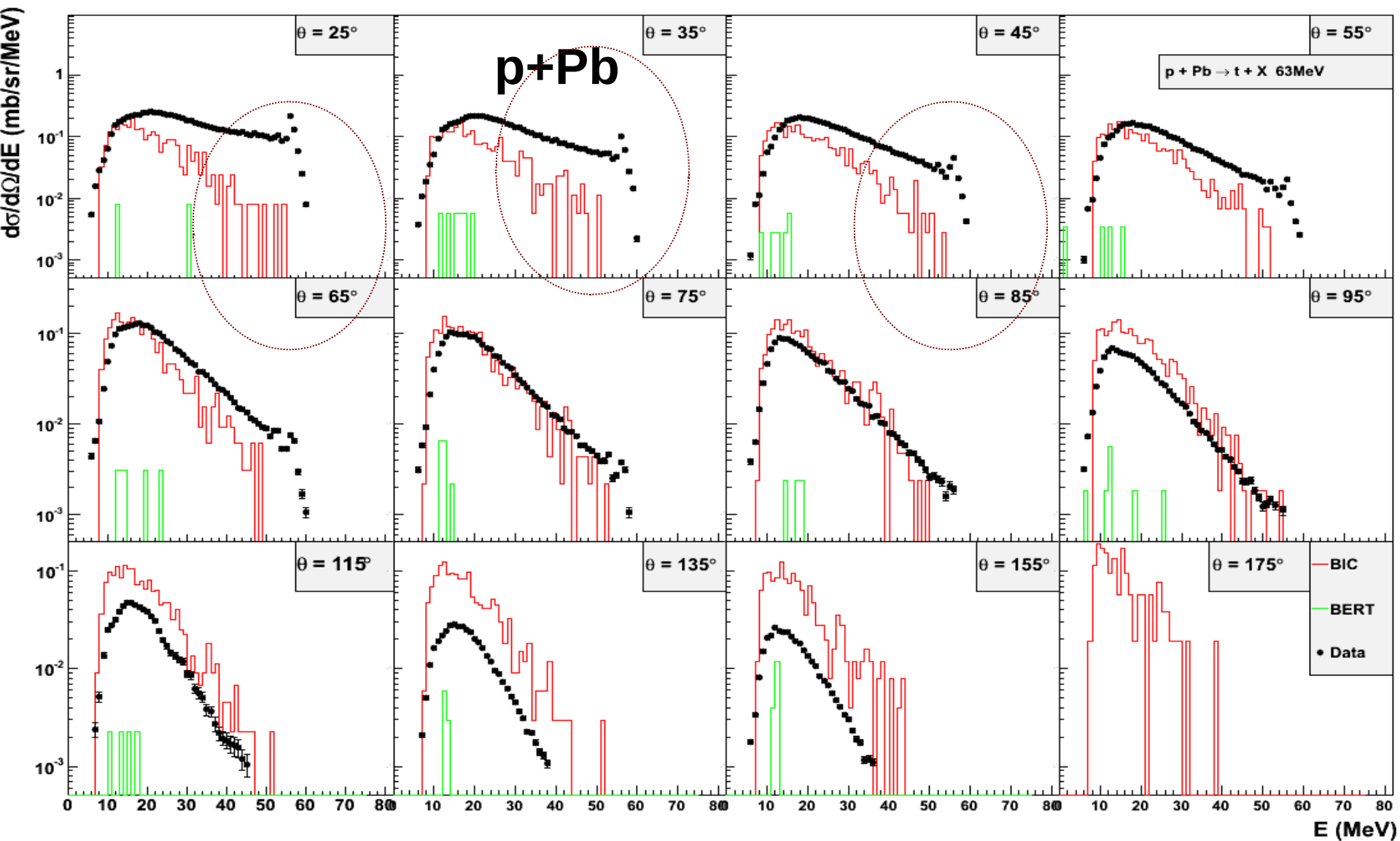




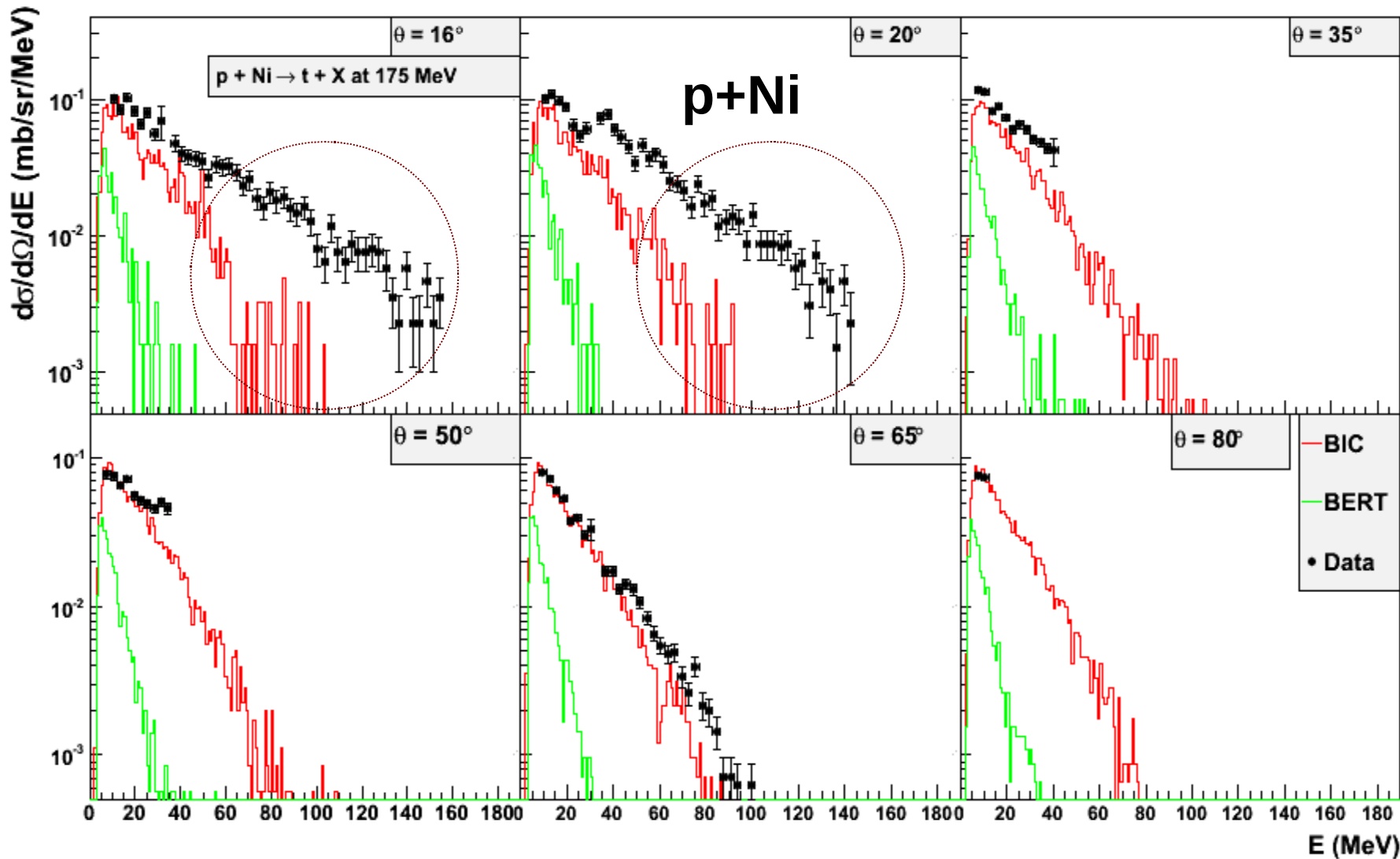
## Deuteron production at 1200 MeV

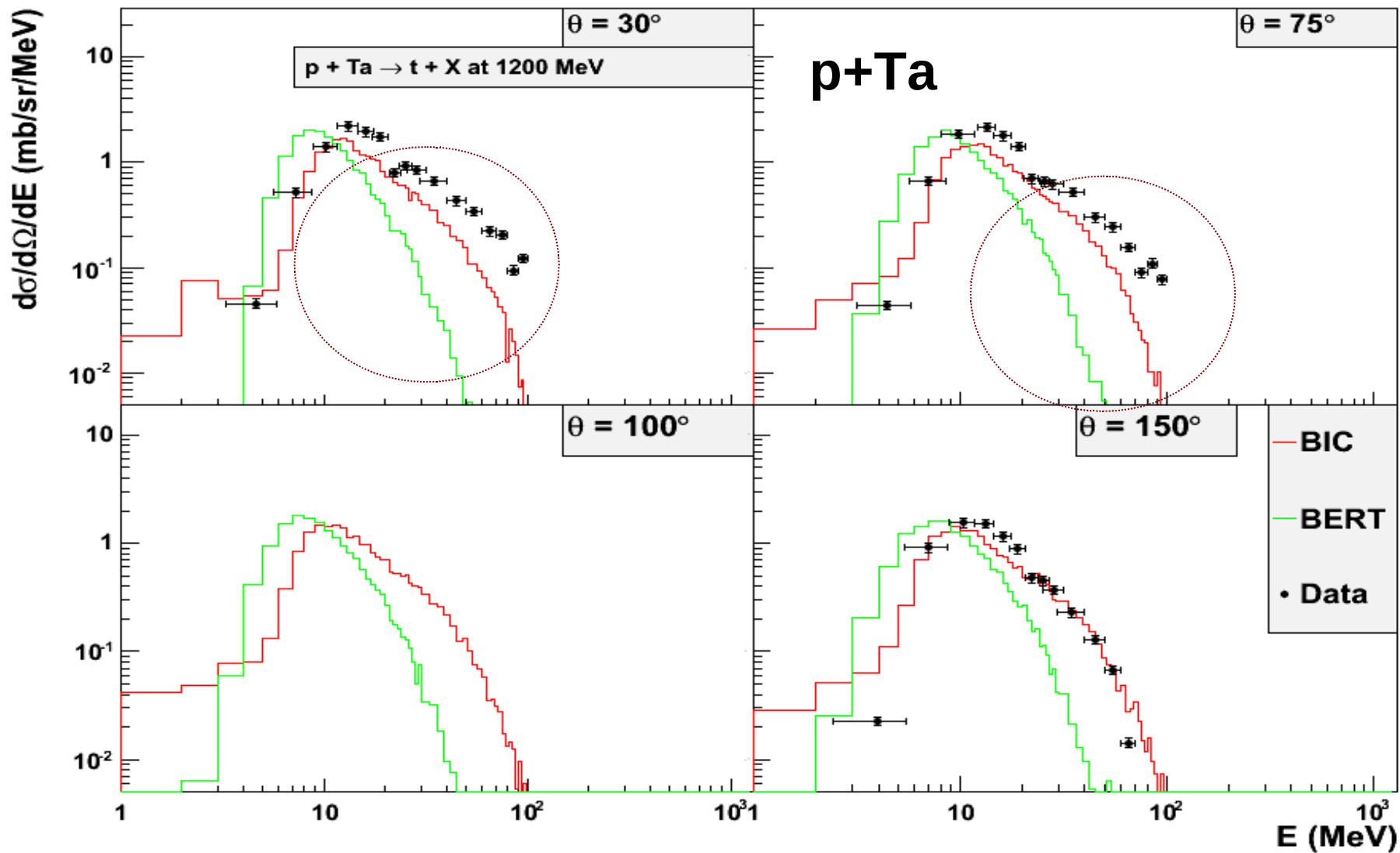


## Tritium production at 63 MeV

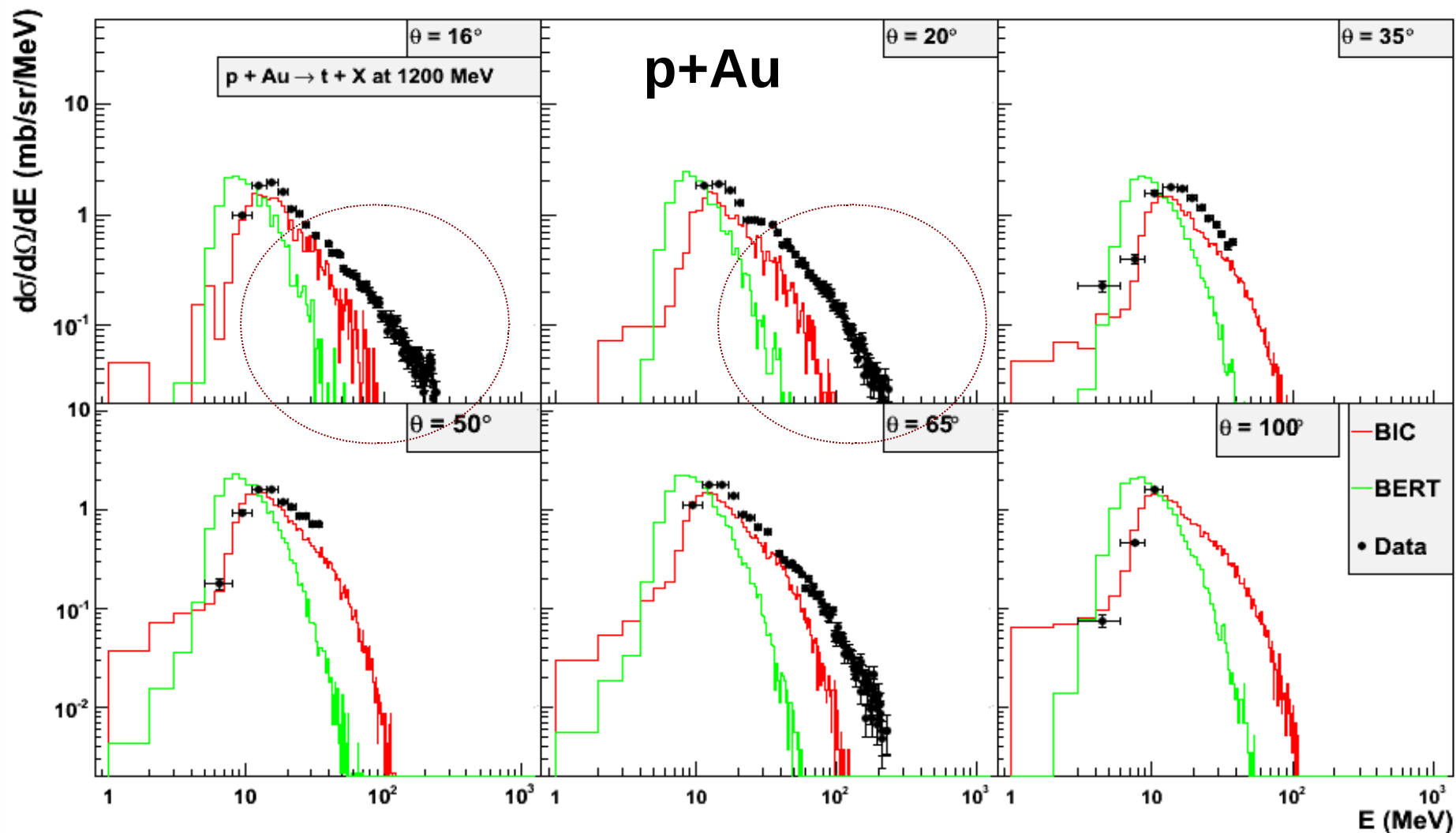


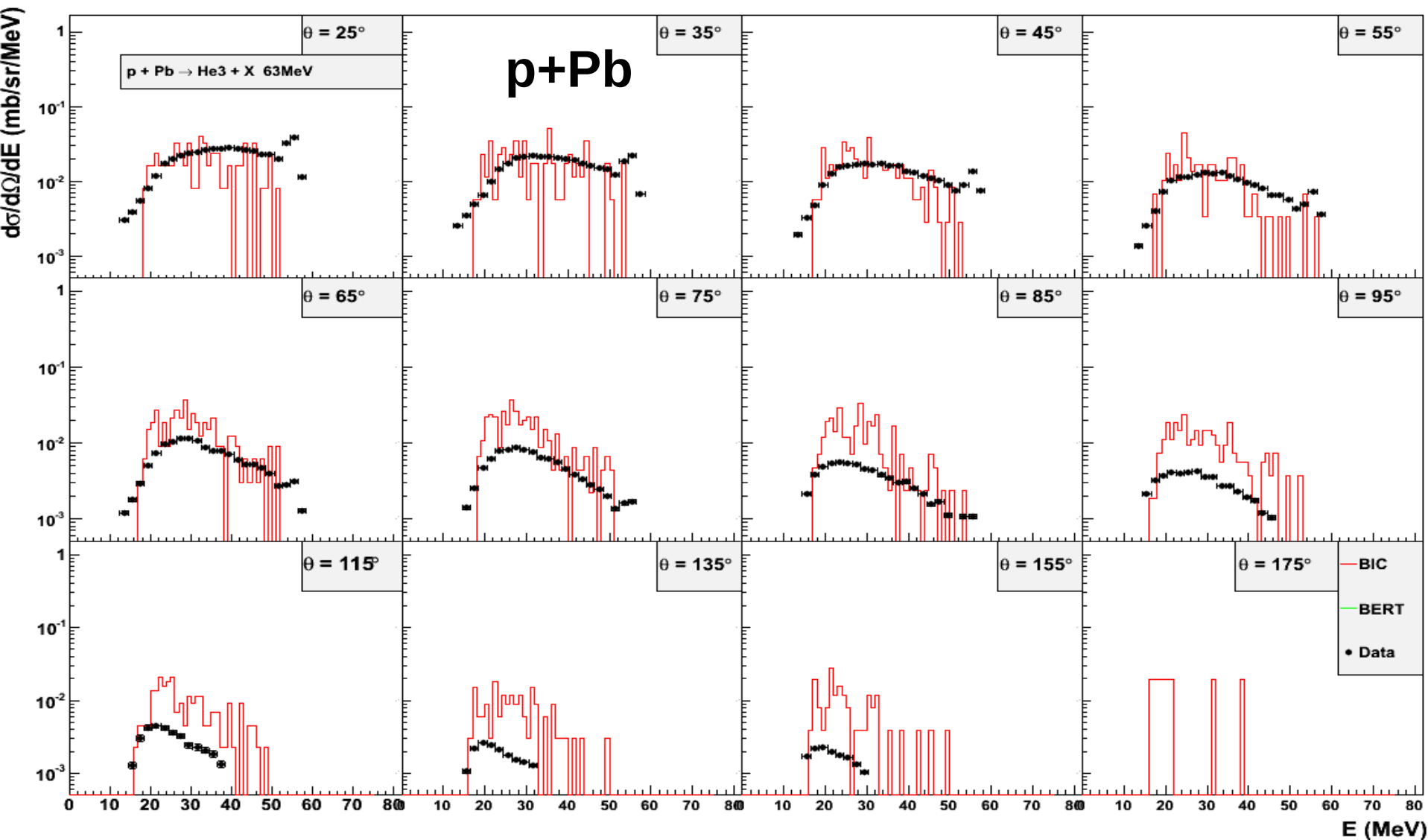
## Tritium production at 175 MeV

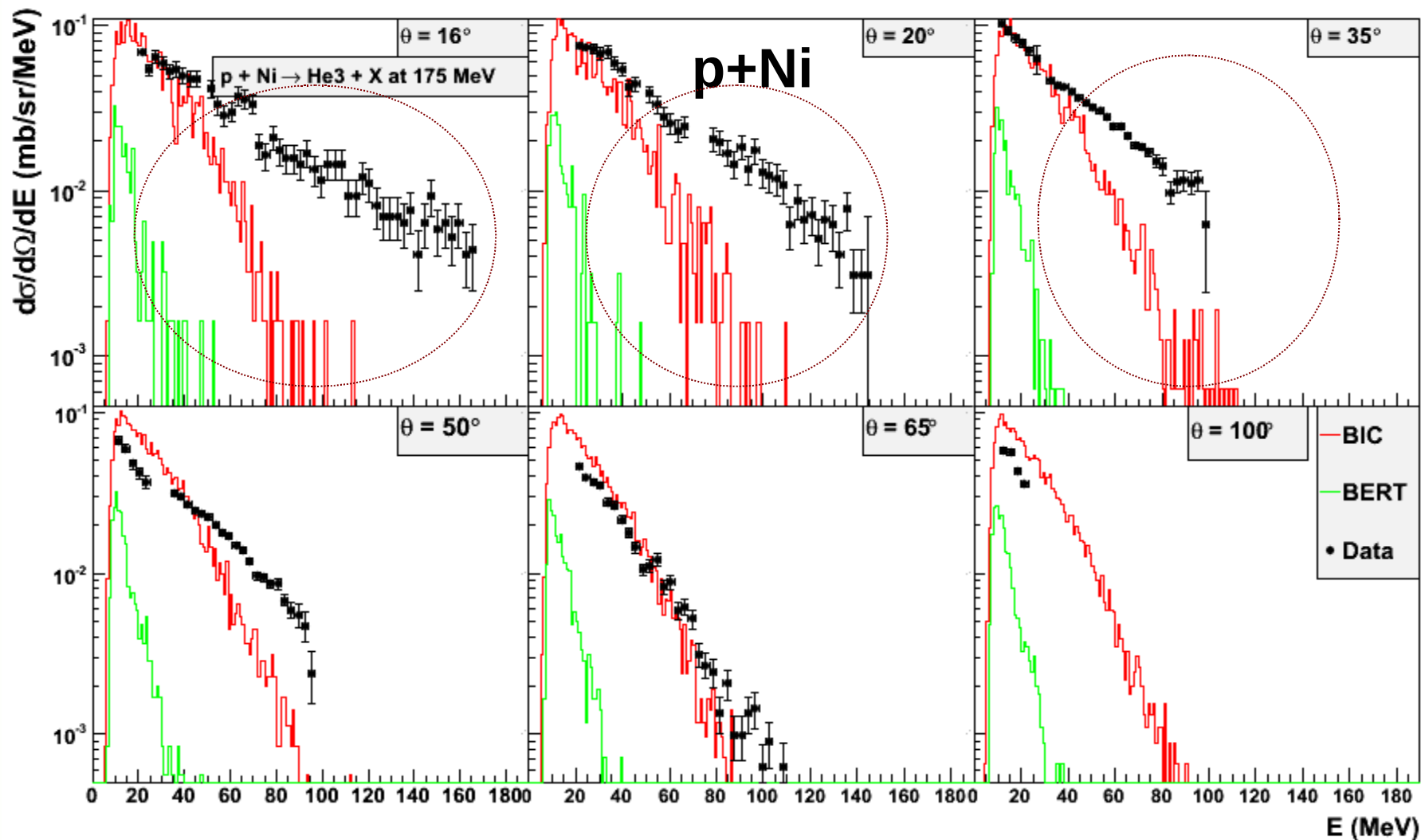


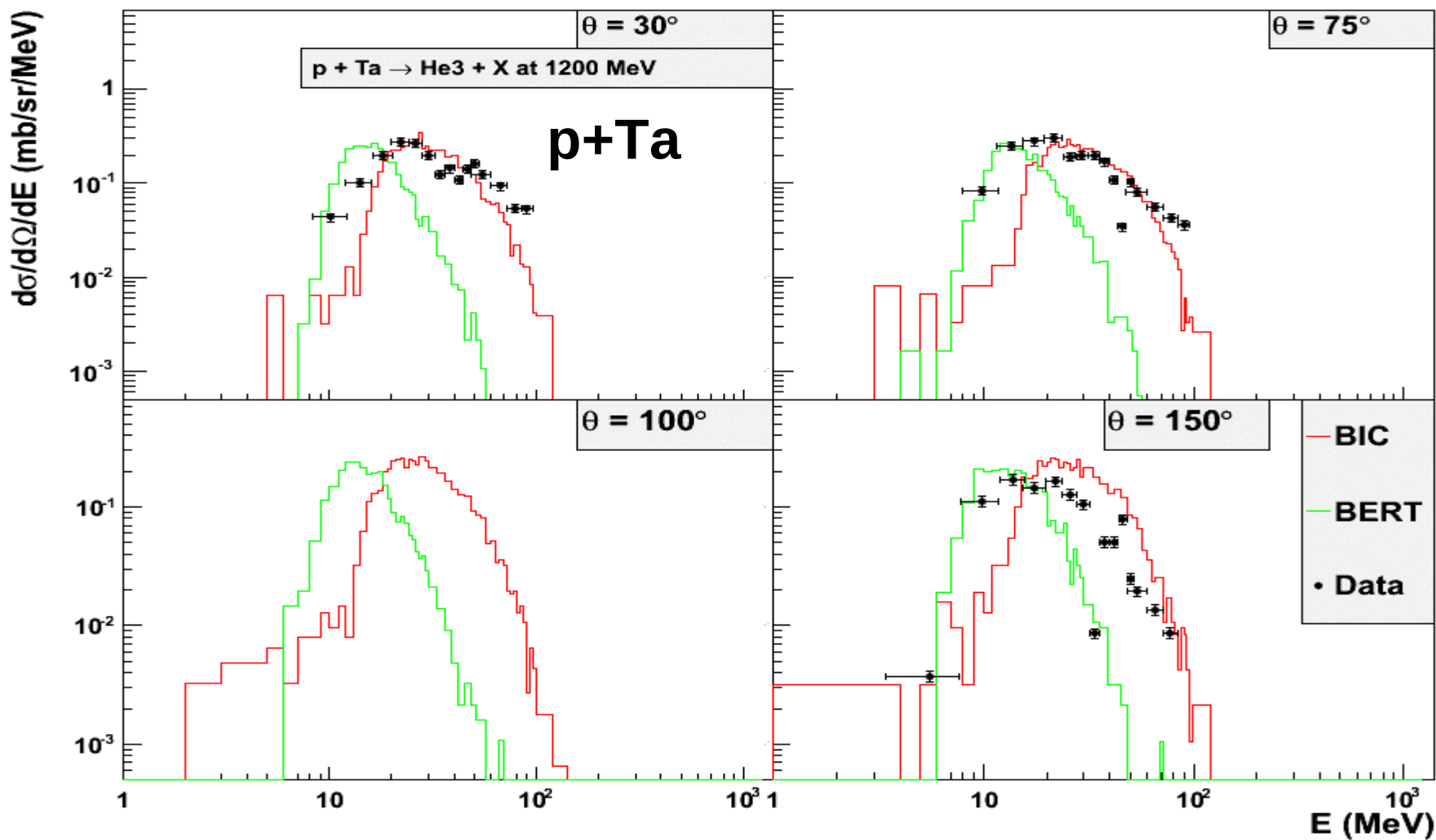


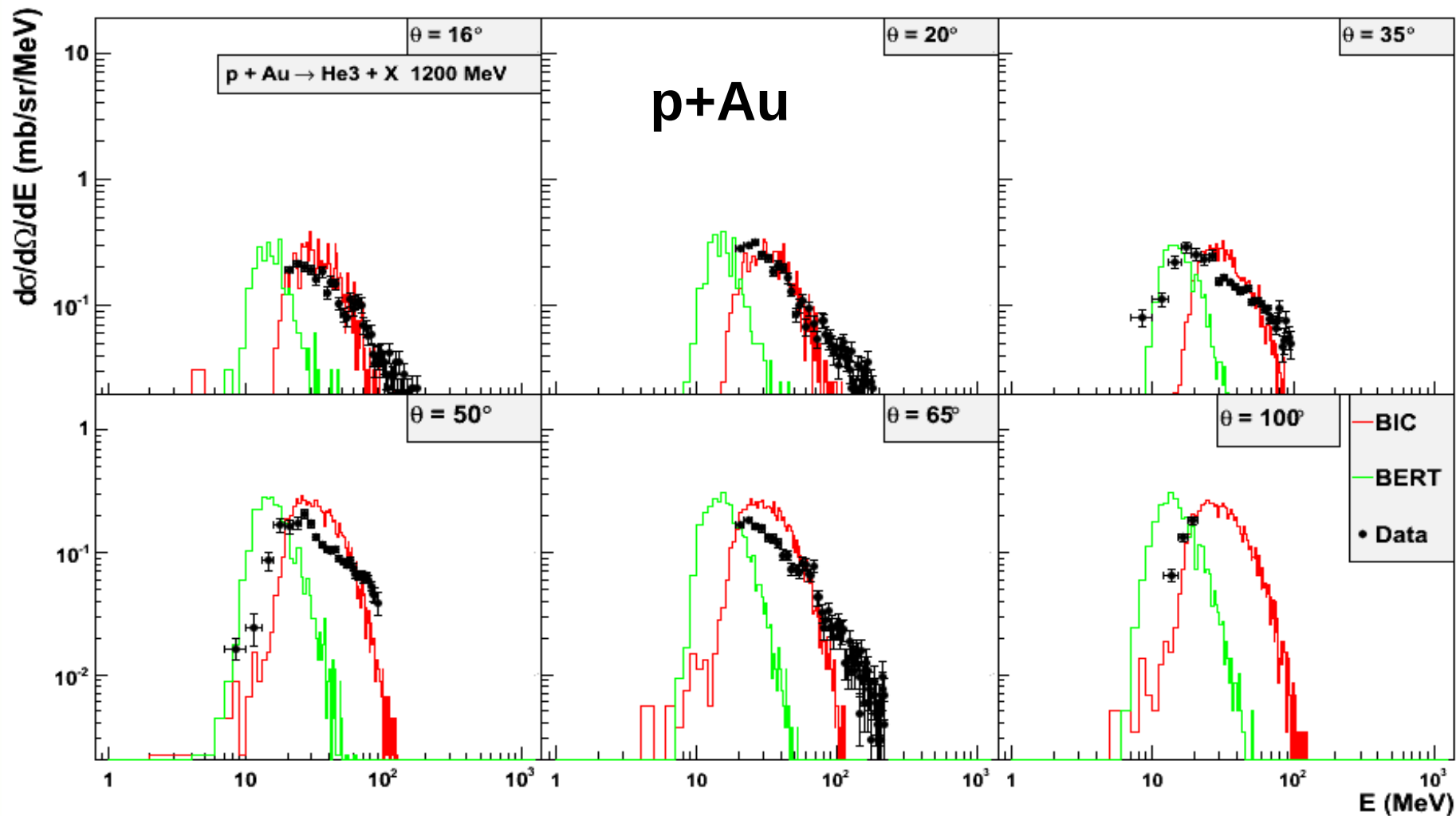


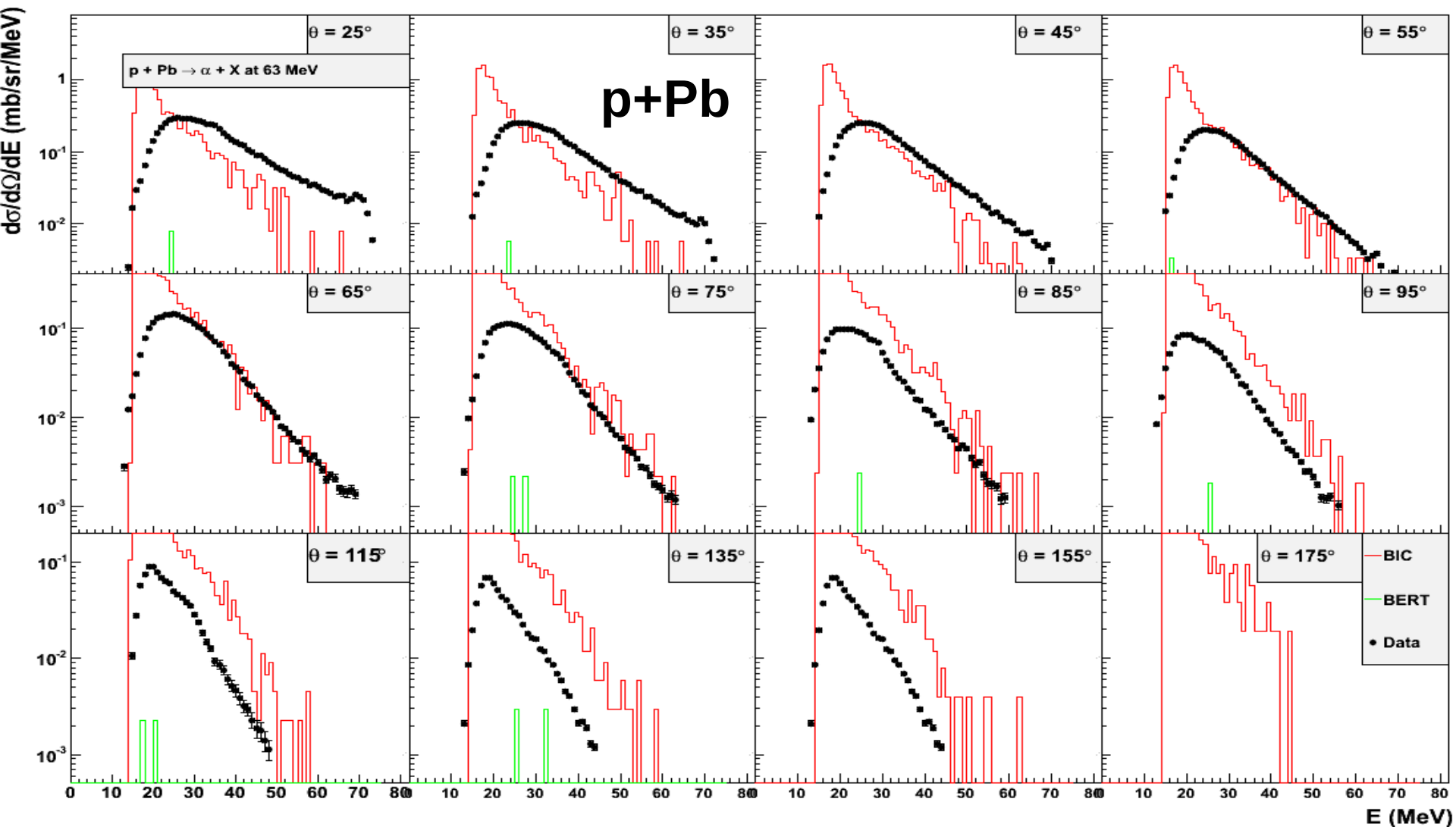


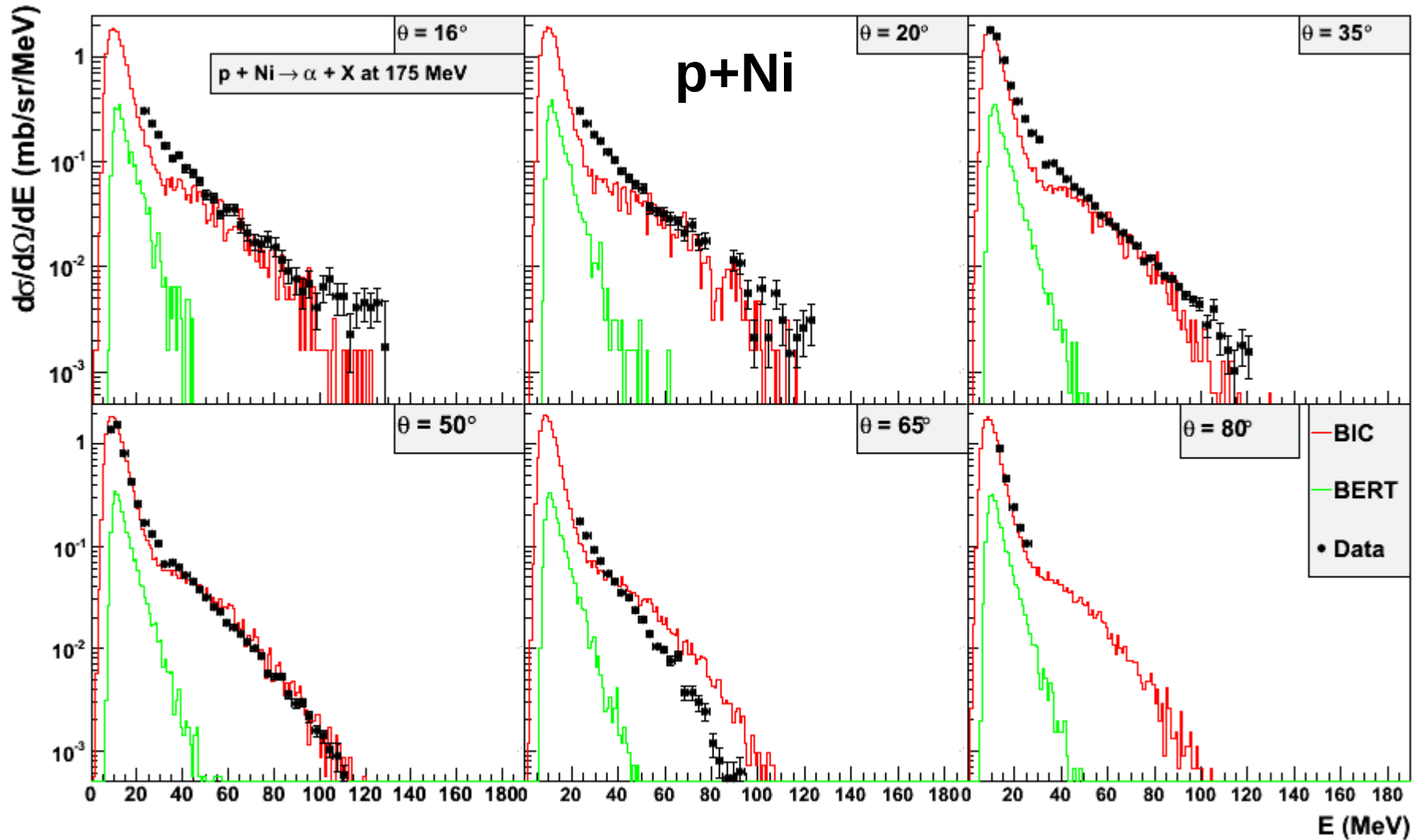




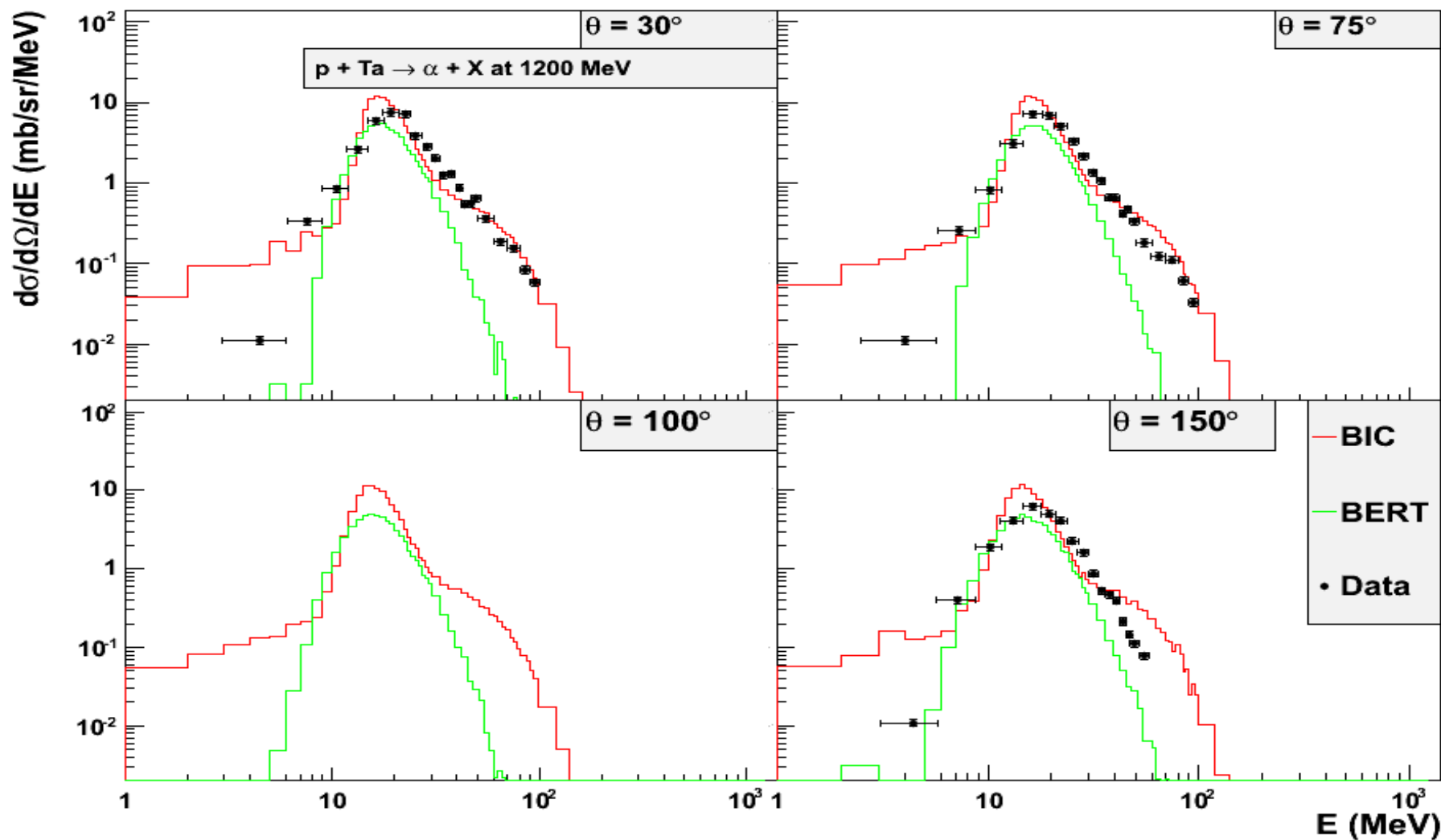




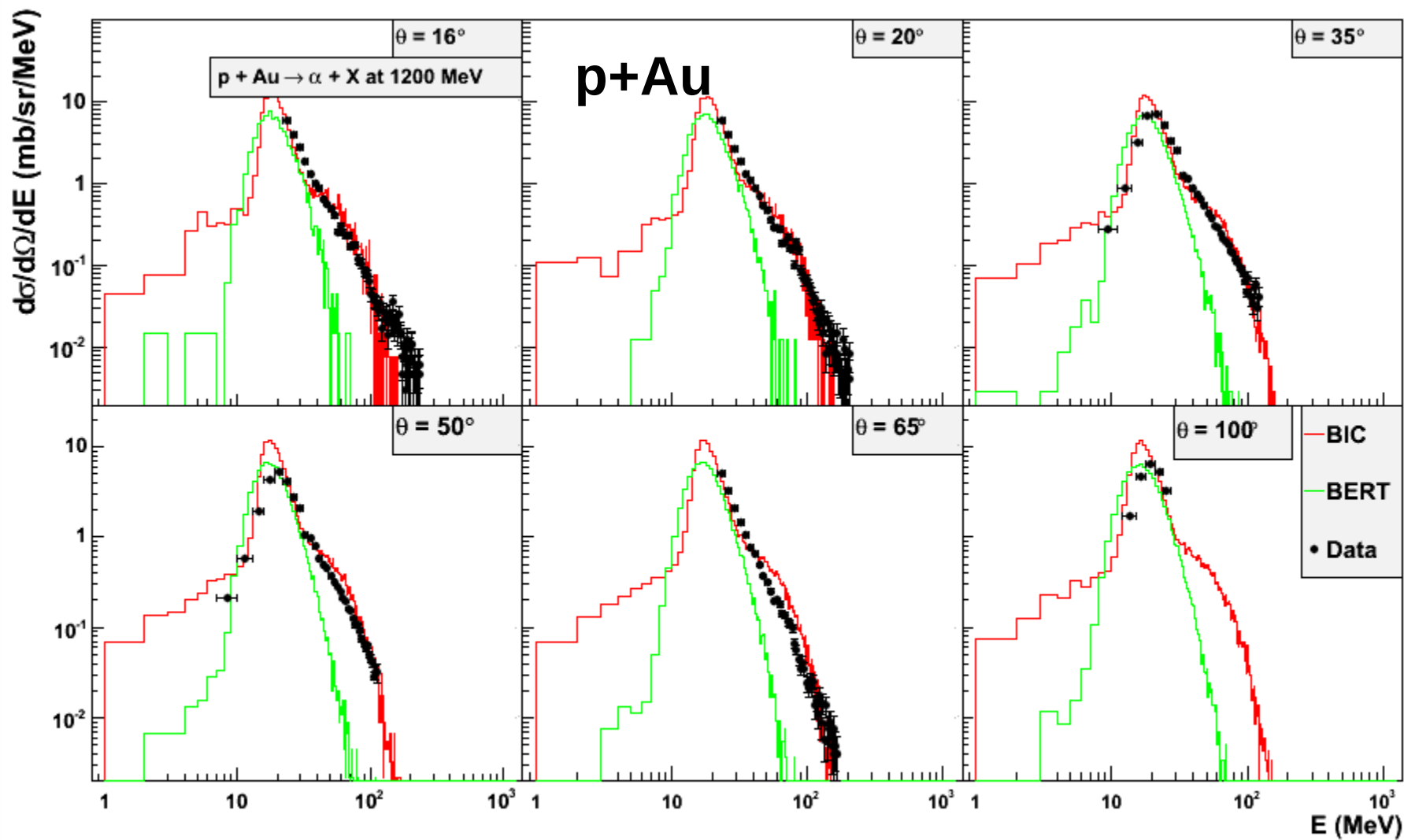




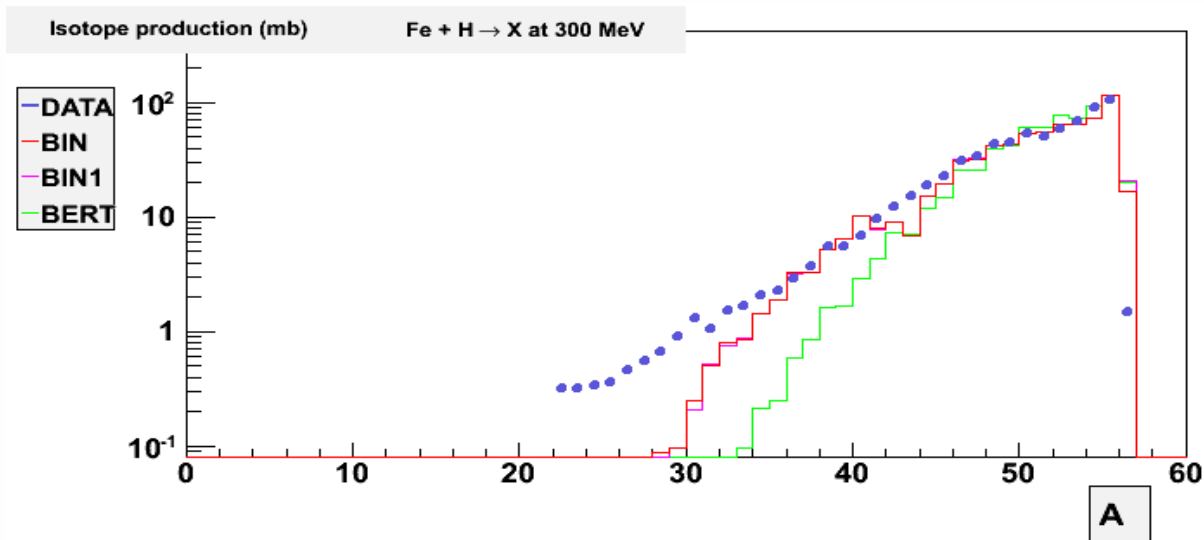
## Alpha production at 1200 MeV





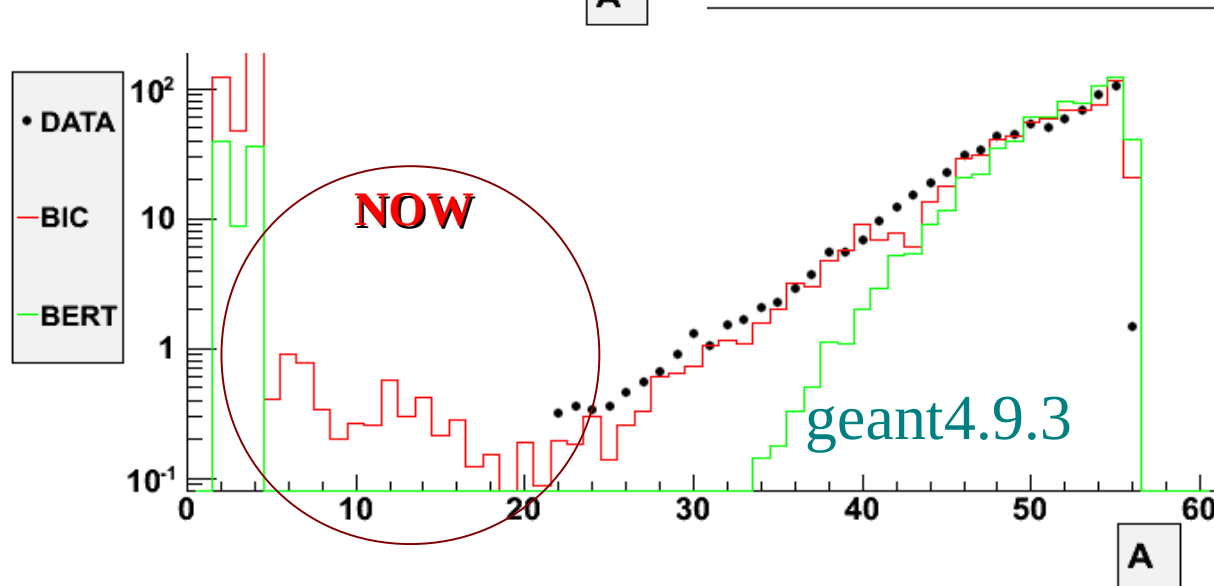


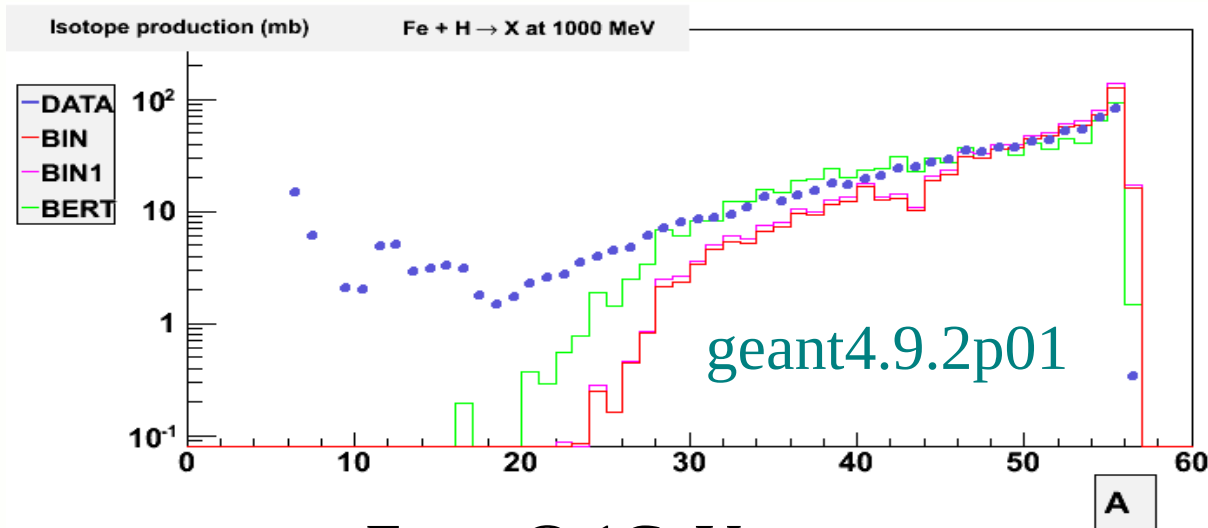
## Isotopic distribution at 0.3 GeV



Fe+p @ 0.3 GeV

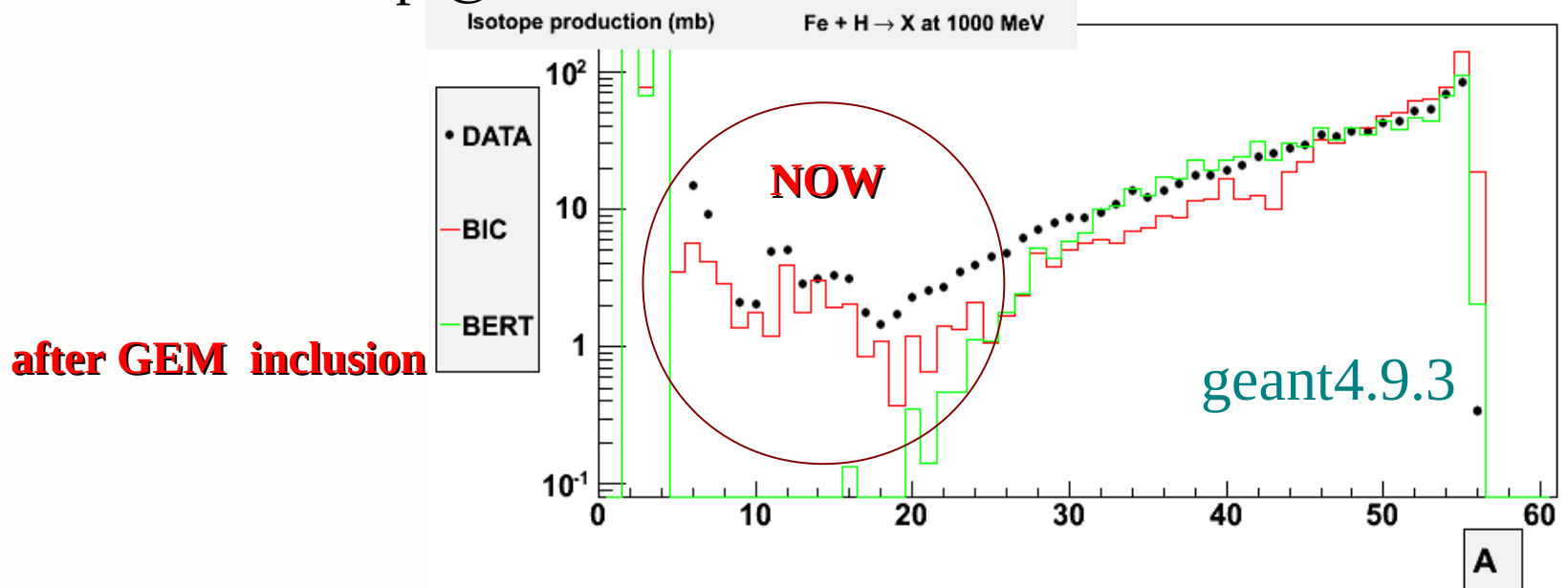
after GEM inclusion

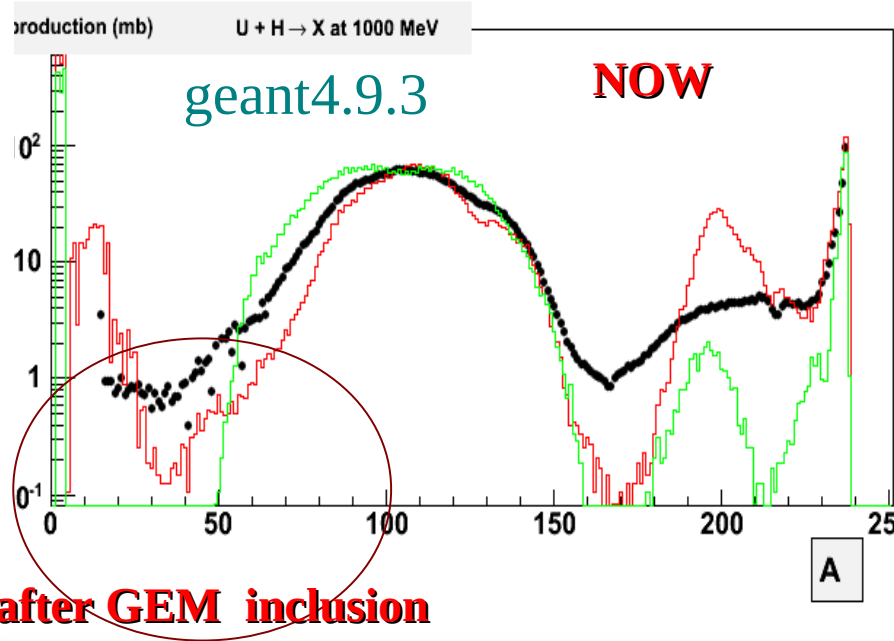
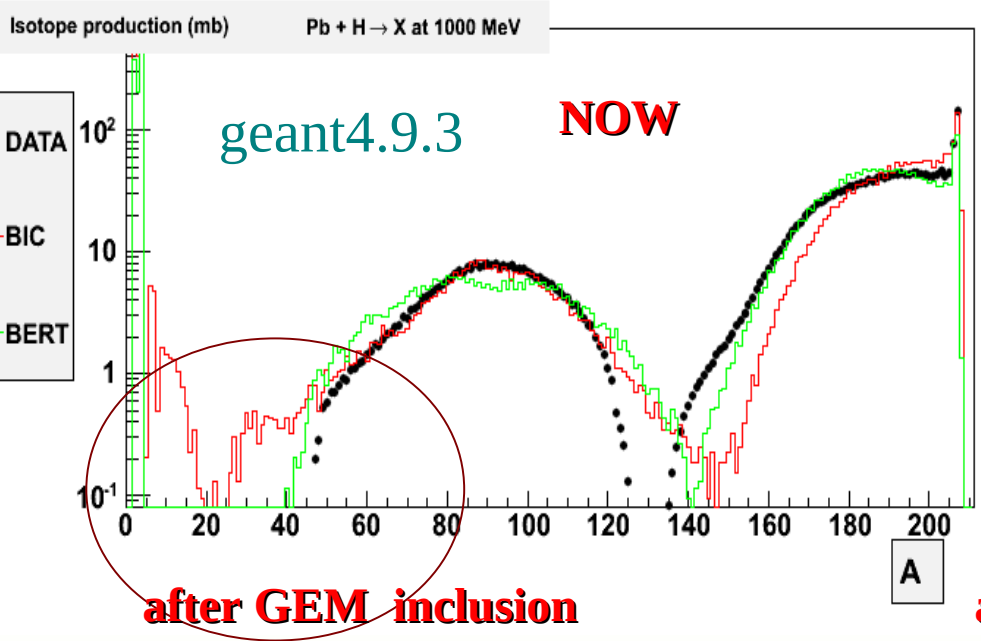
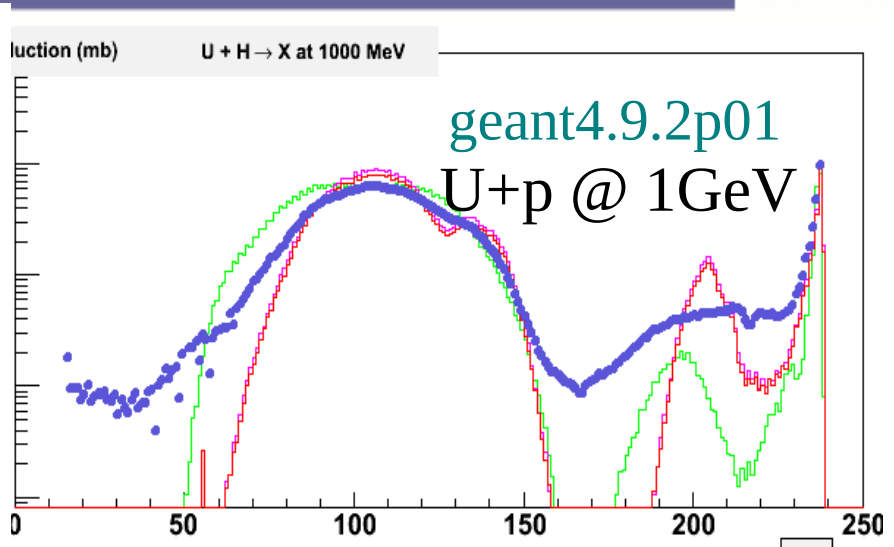
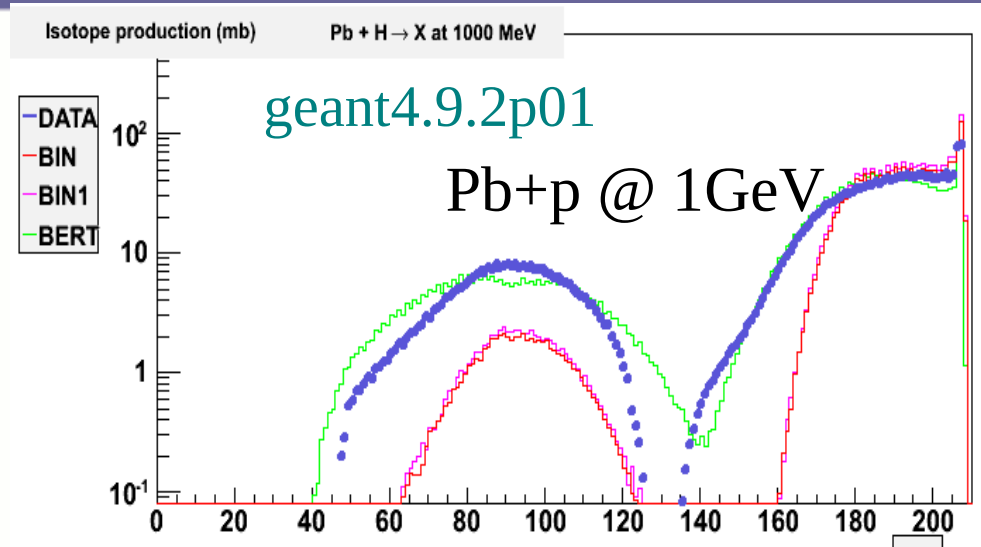


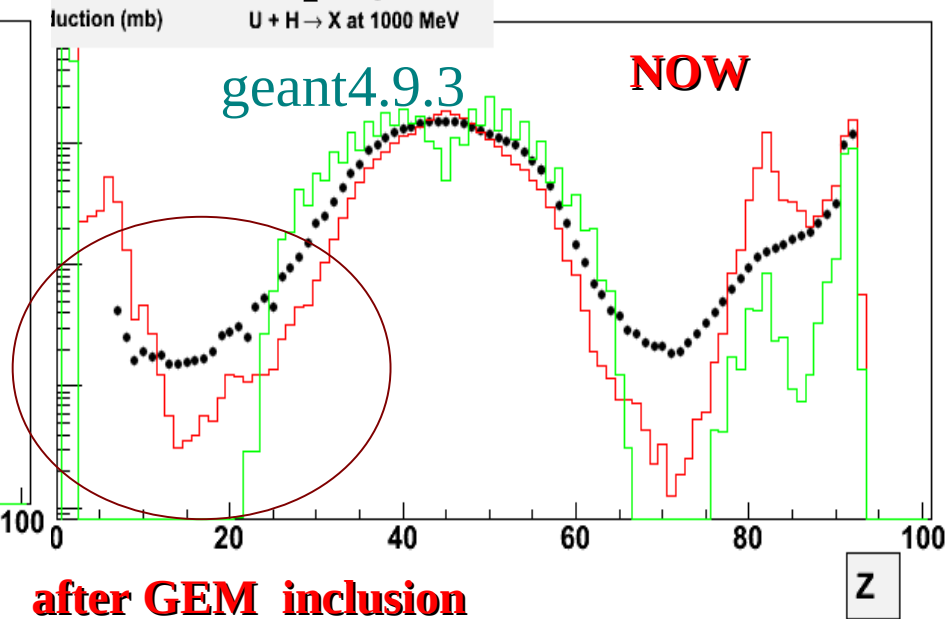
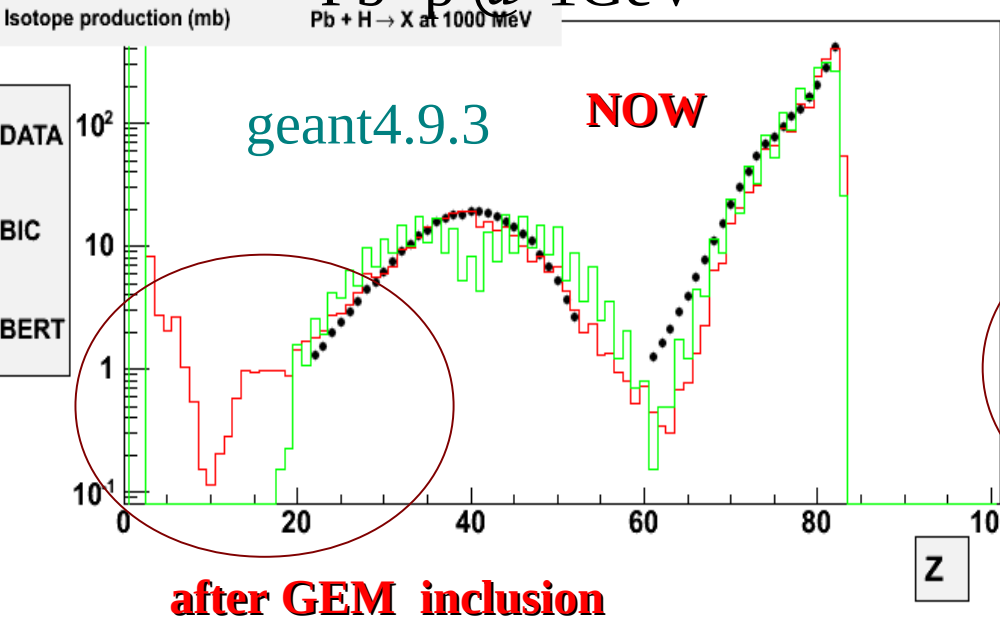
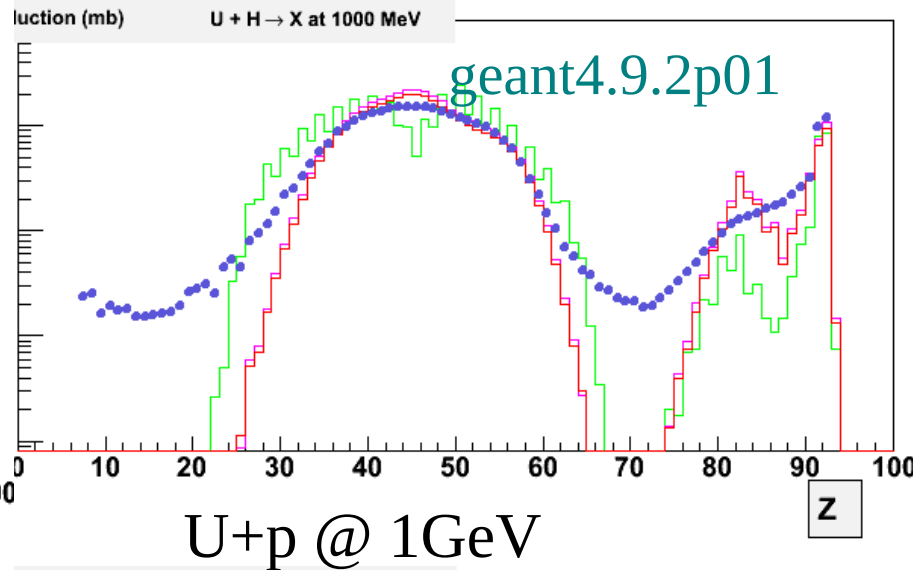
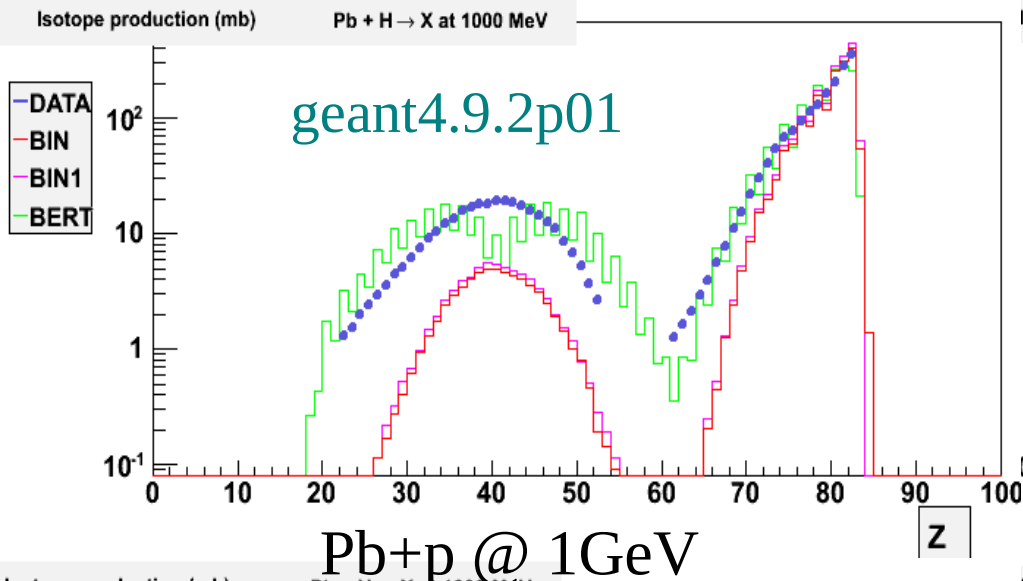


p+Fe

Fe+p @ 1GeV







## Work in progress (I)

(in development version, not yet in official release )

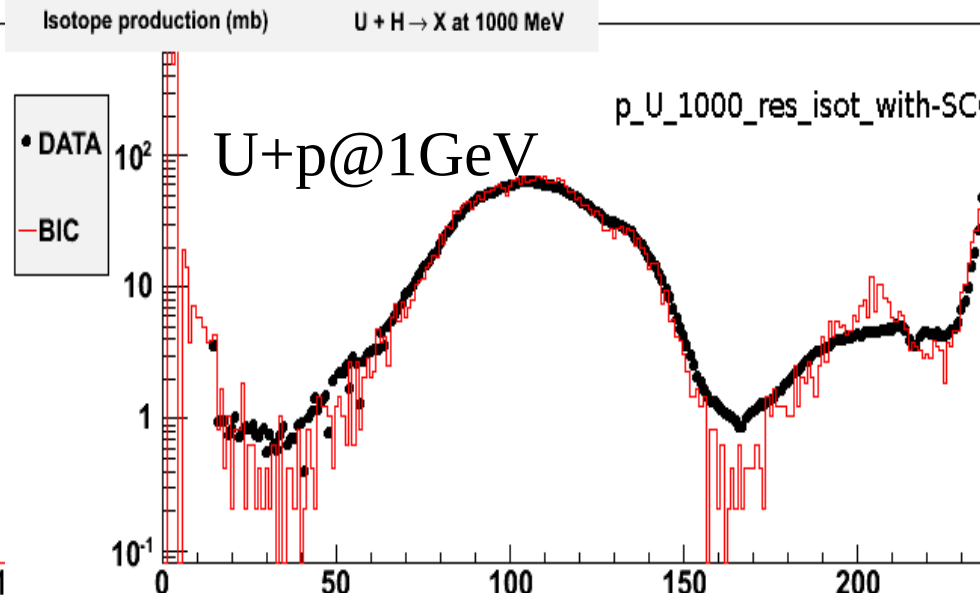
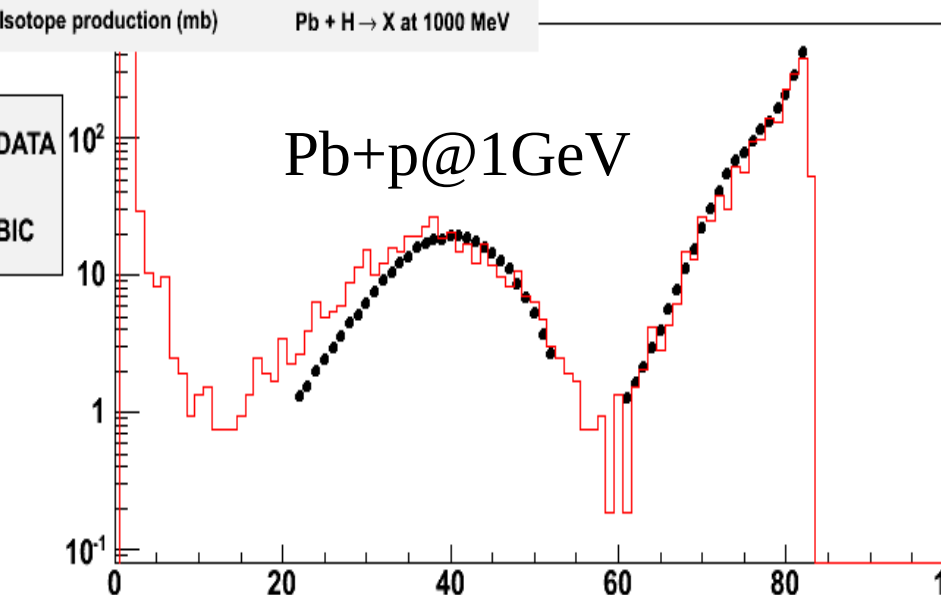
- There was still room for improvements, specially in the spallation products region
- *Soft cut-off* transition from pre-equilibrium cures this problem, **but**, in our case, it makes performance at fission worse.
- A new algorithm for  $n_{eq}$  has been implemented ( based on physical criterion of equilibration of transition probabilities, i.e. *exact* calculation for each fragment) , **instead** of

$$n_{eq} \sim \sqrt{2gU}$$

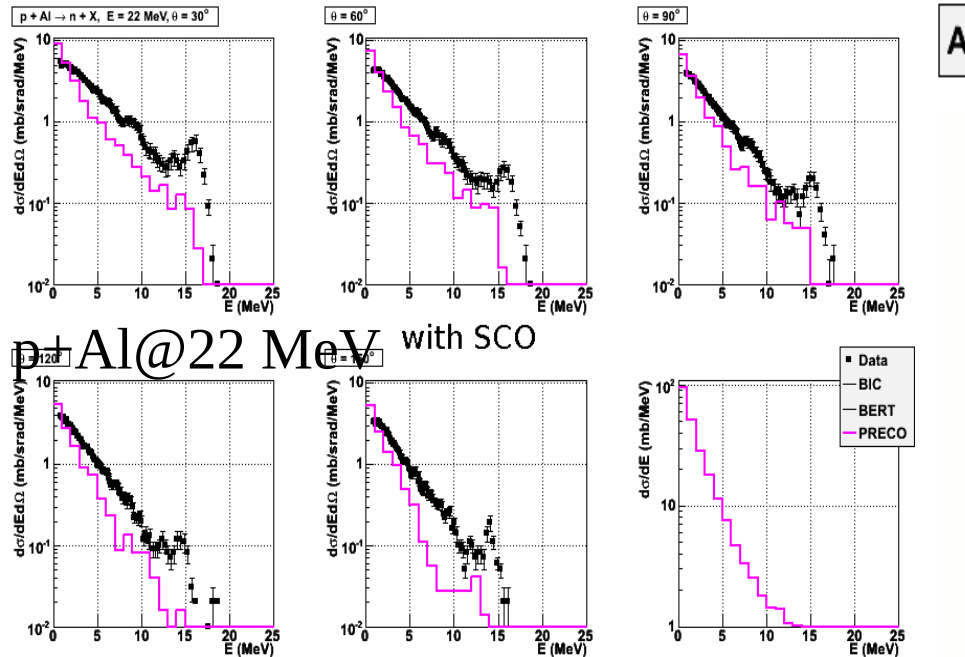
- The *diffusivity* ( $\sigma_{pre}$ ) of the transition has been drastically reduced (i.e. transition is now sharper)

(in development version, not yet in official release )

- No chance for a global set of parameters (optimal for any combination of models)
- Different sets of fission parameters were fitted for each choice (with/without *soft cut-off*).
- They are automatically selected



- “soft cut-off” is ON
- Fission parameters have been fitted
- The situation at pre-equilibrium is quite the same
- CPU time increase (factor ~ 1.5)





## Conclusions (1)

- Bertini agrees better with data for:
  - pion production (high and low energy)
- Binary agrees better with data for:
  - neutron production
  - low and medium energy protons
  - light ion production
  - isotopic distribution  
(price to pay: BIC is about 9 times slower than Bertini for Pb+H @ 1GeV)
- Cases where neither model is better overall
  - one model may be better for forward angles, the other for backward angles
- The fact that we cannot say that one model is clearly better than the other emphasizes the need for alternate models in same energy range

## Conclusions (2)

- This benchmark study demonstrated areas where improvement is needed. As a result:
  - recently made **improvements to pre-compound**
    - *soft* transition to equilibrium has been tuned
  - recently made **improvements to de-excitation**:
    - IMF evaporation (GEM model)
  - recently made **improvements to fission**
    - fission parameters have been tuned
    - improvements are still possible  
(angular momentum dependence of fission barrier?)
- plan to add **coalescence models for cascade stage**:
  - The underestimation in the fast cluster yield (especially deuterons) at small angles clearly points to the direct mechanism of complex particle production

Thanks for your attention